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# Formulation of and numerical studies with the Dutch Atmospheric Large-Eddy Simulation (DALES)

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## Abstract

The current version of the Dutch Atmospheric Large-Eddy Simulation (DALES) model is presented. DALES is a large-eddy simulation model designed for process studies of the atmospheric boundary layer, including convective and stable boundary layers as well as cloudy boundary layers. In addition, DALES can be used for studies of more specific cases, such as flow over sloping or heterogeneous terrain, and dispersion of inert and chemically active species. This paper contains an extensive description of the physical and numerical formulation of the code, and gives an overview of its applications and accomplishments in recent years.

## 1 Introduction

Modern atmospheric research relies on a cascade of observational and modeling tools. The most detailed type of numerical modeling available are Large-Eddy Simulations (LES). This type of modeling is widely used for atmospheric boundary layer (ABL) studies and provides, in combination with observations, the basis for many cloud and boundary layer parameterizations in models on the other side of the cascade, such as General Circulation Models.

The principle of LES is to resolve the turbulent scales larger than a certain filter width, and to parameterize the smaller, less energetic scales. This filter width is usually related to the grid size of the LES, and ranges typically between 1 m for stably stratified boundary layers, to 50 m for simulations of the cloud-topped ABL. In such a typical LES set up, up to 90% of the turbulent energy resides in the resolved scales. For applications of DALES like the ones presented in this paper, LES has the advantage over coarser models that it relies only weakly on parameterizations. In comparison with observational studies, LES has the advantage of providing a complete data set, in terms of time and space, and in terms of diagnosable variables. Especially the combined use of LES and observations is a popular methodology in process studies of the ABL.

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LES modeling of the ABL started in the late sixties (e.g., Lilly, 1967; Deardorff, 1972); cloudy boundary layers were first simulated by Sommeria (1976). From Nieuwstadt and Brost (1986) onward, several cycles of intercomparison studies compare state-of-the-art LES models with observational studies and with each other. The aim of these studies was not so much to determine which LES model performs best in which situation, but more to determine the strengths and weaknesses of LES. Two particularly active cycles are organized under the umbrella of the Global Energy and Water Cycle Experiment (GEWEX): the GEWEX Atmospheric Boundary Layers Study (GABLS), and the GEWEX Cloud System Study (GCSS) Boundary Layer Cloud Working Group. The GABLS focusses on the clear boundary layer, mainly on stable and transitional situations (Holtslag, 2006; Beare et al., 2006; Basu et al., 2008). The GCSS looks at different aspects of boundary layer clouds, mainly shallow cumulus and stratocumulus clouds (Bretherton et al., 1999a,b; Duynkerke et al., 1999, 2004; Brown et al., 2002; Siebesma et al., 2003; Stevens et al., 2001, 2005; Ackerman et al., 2009; van Zanten et al., 2010).

The Dutch Atmospheric Large-Eddy Simulation (DALES) has joined in virtually all of these intercomparisons. Besides convective, stable and cloud-topped boundary layers, DALES has also been used on a wide range of topics, such as studies of shear driven flow, heterogeneous surfaces, dispersion and turbulent reacting flows in the ABL, and of flow over sloped terrain. As such, DALES is one of the most all-round tested available LES models for studies of the ABL. In this paper, we aim to describe and validate DALES 3.2, the current version of DALES.

In the remainder of this paper, we first give a thorough description of the code in Sect. 2. In Sect. 3, an overview of studies conducted with DALES is given, both as a validation of the code as well as an overview of the capabilities of an LES like DALES. In Sect. 4, an outlook is given on future studies that are planned to be done with DALES, as well as an outlook on future improvements.

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## 2 Description of the code

### 2.1 Generalities

DALES is rooted in the LES model of Nieuwstadt and Brost (1986). Cuijpers and Duynkerke (1993) first used DALES for moist convection, and provide a general description of an older version of DALES. Large parts of the code have been changed ever since and contributions of many people over a number of years have resulted in the current version 3.2 of DALES. Currently, DALES is maintained by researchers from Delft University, the Royal Netherlands Meteorological Institute (KNMI), Wageningen University and the Max Planck Institute for Meteorology.

Notable changes in comparison with the version that has been described by Cuijpers and Duynkerke (1993), include: Different time integration and advection schemes, revised subfilter-scale, surface and radiation schemes, addition of a cloud-microphysical scheme, capabilities for chemical reactive scalar transport and for Lagrangian particle dispersion, for flow over heterogeneous and for flow over sloped terrain. These revisions in DALES result in faster simulations and higher stability, and in an easier and more extendable user interface. Due to the modular setup of the code, newly written code for specific applications of DALES can easily improve the code as a whole. This makes DALES suitable as a community model; besides the actively developing core users, the model is currently used in several other institutes across the world.

DALES 3.2 is released under the GPLv3 license. It is available at [dales.ablresearch.org](http://dales.ablresearch.org). Documentation is also available there. Although the code is completely free to use, to modify and to redistribute, it is regarded courtesy to share bug fixes and extensions that can be of general interest, and to keep in contact with the core developers. Given the experimental character of the code, it is also appreciated to discuss co-authorship in case of publications coming out of research conducted with DALES.

DALES is written in Fortran 95. The only dependency of DALES are on cmake for building, and on the Message Passing Interface (MPI). Some optional modules also require NetCDFv3. Code for Fourier transformations is provided as well, leaving

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DALES as portable as possible. To the best knowledge of the authors, DALES runs on all common combinations of platform architecture, compiler, and MPI implementation. Currently, an effort is being made to port DALES to NVIDIA graphical processors, using CUDA.

5 The prognostic variables of DALES are the three velocity components  $u_i$ , the liquid water potential temperature  $\theta_l$ , the total water specific humidity  $q_t$ , the rain water specific humidity  $q_r$ , the rain droplet number concentration  $N_r$ , and up to 100 passive or reactive scalars. Because of the one-and-a-half order scheme that parameterizes sub-filter scale dynamics, the subfilter-scale turbulent kinetic energy (SFS-TKE,  $e$ ) counts as an additional prognostic variable. To decrease simulation time, most prognostic variables can be switched off; only calculations of  $u_i$ ,  $e$ , and  $\theta_l$  are obligatory.

Given that ice is not currently implemented in the model, the total water specific humidity is defined as the sum of the water vapor specific humidity  $q_v$  and the cloud liquid water specific humidity  $q_c$ :

$$15 \quad q_t = q_v + q_c \quad (1)$$

Note that this definition excludes the rain water specific humidity  $q_r$  from  $q_t$ . Any conversion between rain water on the one hand, and cloud water or water vapor on the other hand, will therefore enter the equations for  $q_t$  and for  $\theta_l$  as an addition source term. As a definition of  $\theta_l$ , we use the close approximation explained by Emanuel (1994):

$$20 \quad \theta_l \approx \theta - \frac{L}{c_{pd}\Pi} q_c \quad (2)$$

with  $L=2.5 \times 10^6 \text{ J kg}^{-1}$  the latent heat release of vaporization,  $c_{pd}=1004 \text{ J kg}^{-1} \text{ K}^{-1}$  the heat capacity of dry air, and  $\Pi$  the exner function:

$$\Pi = \left( \frac{p}{p_0} \right)^{\frac{R_d}{c_{pd}}}, \quad (3)$$

25 in which  $R_d=287.0 \text{ J kg}^{-1} \text{ K}^{-1}$  is the gas constant for dry air and  $p_0=10^5 \text{ Pa}$  is a reference pressure.

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In the absence of precipitation and other explicit source terms,  $\theta_l$  and  $q_t$  are conserved variables. The virtual potential temperature  $\theta_v$  is in good approximation defined with:

$$\theta_v \approx \left( \theta_l + \frac{L}{c_{pd}\Pi} q_c \right) \left( 1 - \left( 1 - \frac{R_v}{R_d} \right) q_t - \frac{R_v}{R_d} q_c \right), \quad (4)$$

with  $R_v = 461.5 \text{ J kg}^{-1} \text{ K}^{-1}$ , the gas constants for water vapor. The most important thermodynamical constants that are used throughout this paper are summarized in Table 1.

DALES assumes the Boussinesq approximation, and is run on an Arakawa C-grid (see Fig. 2). The pressure, the SFS-TKE, and the scalars are defined at cell center, the three velocity components are defined at the West side, the South side, and the bottom side of the grid cell, respectively.

Hereafter, quantities that are averaged over the LES-filter are denoted with a tilde  $\tilde{\cdot}$ , time averages with a overbar  $\overline{\cdot}$ , and averages over the two horizontal directions of the domain with angular brackets  $\langle \cdot \rangle$  (slab average). The prognostic scalars can often be treated simultaneously as the generic scalar field  $\varphi \in \{\theta_l, q_t, q_r, N_r, s_n\}$ . Primes denote the subfilter-scale fluctuations with respect to the filtered mean. To remain consistent with notational conventions as used in literature and also in the source code of DALES, some symbols can have different meaning between different subsections. In such cases, the immediate context should always make it clear what each symbol stands for in a particular section. Vertical velocities and fluxes are in general positive when directed upward; only the radiative and sedimentation fluxes point downward, following conventions.

In the following sections, different components of the model are described one by one. Sections 2.2–2.6 describe the physical and numerical components that are necessary to conduct a minimal experiment with DALES. After that, Sects. 2.7–2.11 describe various forcings and source terms that extend the core of DALES for use in more specific applications. Finally, Sect. 2.12 describes the most relevant statistical routines in DALES.

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## 2.2 The governing equations

Within the Boussinesq approximation the equations of motion, after application of the LES filter, are given by

$$\frac{\partial \tilde{u}_i}{\partial x_i} = 0, \quad (5)$$

$$\frac{\partial \tilde{u}_i}{\partial t} = -\frac{\partial \tilde{u}_i \tilde{u}_j}{\partial x_j} - \frac{\partial \pi}{\partial x_i} + \frac{g}{\theta_0} \tilde{\theta}_v \delta_{i3} + \mathcal{F}_i - \frac{\partial \tau_{ij}}{\partial x_j}, \quad (6)$$

$$\frac{\partial \tilde{\varphi}}{\partial t} = -\frac{\partial \tilde{u}_j \tilde{\varphi}}{\partial x_j} - \frac{\partial R_{u_j, \varphi}}{\partial x_j} + \mathcal{S}_\varphi, \quad (7)$$

where the tildes denote the filtered mean variables. Molecular transport terms have been neglected. The  $z$ -direction ( $x_3$ ) is taken to be vertical.  $\theta_0$  is the reference state potential temperature taken to be equal to the value at the surface, and  $\mathcal{F}_i$  represents other forcings, including large scale forcings and the Coriolis forcing

$$\mathcal{F}_i^{\text{cor}} = -2\epsilon_{ijk} \Omega_j \tilde{u}_k, \quad (8)$$

where  $\Omega$  is the earth's angular velocity. Source terms for scalar  $\varphi$  are denoted by  $\mathcal{S}_\varphi$ , and may include of microphysical ( $\mathcal{S}^{\text{mcr}}$ ), radiative ( $\mathcal{S}^{\text{rad}}$ ), chemical ( $\mathcal{S}^{\text{chem}}$ ), large-scale ( $\mathcal{S}^{\text{ls}}$ ), and relaxation ( $\mathcal{S}^{\text{rel}}$ ) terms. The subfilter-scale (SFS), or residual, scalar fluxes are denoted by  $R_{u_j, \varphi} \equiv \widetilde{u_j \varphi} - \tilde{u}_j \tilde{\varphi}$ , i.e., the contribution to the resolved motion from all scales below the LES filter width. The anisotropic SFS-stress tensor is defined by

$$\tau_{ij} \equiv \widetilde{u_i u_j} - \tilde{u}_i \tilde{u}_j - \frac{2}{3} \delta_{ij} e, \quad (9)$$

where  $e = \frac{1}{2}(\widetilde{u_i u_i} - \tilde{u}_i \tilde{u}_i)$  is the subfilter-scale turbulent kinetic energy (SFS-TKE).

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Formally, the trace of the SFS-stress has been included in the modified pressure

$$\pi = \frac{\tilde{p}}{\rho} + \frac{2}{3}e. \quad (10)$$

To enforce continuity (Eq. 5), a Poisson equation for  $\pi$  is solved

$$\frac{\partial^2 \pi}{\partial x_i^2} = \frac{\partial}{\partial x_i} \left( -\frac{\partial \tilde{u}_i \tilde{u}_j}{\partial x_j} + \frac{g}{\theta_0} \tilde{\theta}_v \delta_{i3} + \mathcal{F}_i - \frac{\partial \tau_{ij}}{\partial x_j} \right). \quad (11)$$

- 5 Since computations are performed in a double periodic domain, the Poisson equation is solved by applying a Fast Fourier Transform in the homogenous directions followed by solving a linear system in the  $z$ -direction. A schematic overview of how the different processes affect the different variables is given in Fig. 1.

### 2.3 Subfilter-scale model

- 10 In DALES, the SFS stress tensor is modeled as an eddy diffusivity:

$$\tau_{ij} = -K_m \left( \frac{\partial \tilde{u}_i}{\partial x_j} + \frac{\partial \tilde{u}_j}{\partial x_i} \right), \quad (12)$$

$$R_{u_j, \varphi} = -K_h \frac{\partial \tilde{\varphi}}{\partial x_j}, \quad (13)$$

with the eddy diffusivity coefficients  $K_m$  and  $K_h$ . In DALES, the eddy diffusivity coefficients are modeled in two ways: Either as a function of the SFS-TKE  $e$  (Deardorff, 1973) (which is the default), or using Smagorinsky closure (Smagorinsky, 1963).

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### 2.3.1 SFS-TKE model

The prognostic equation for  $e$ :

$$\frac{\partial e}{\partial t} = -\frac{\partial \tilde{u}_j e}{\partial x_j} - \tau_{ij} \frac{\partial \tilde{u}_i}{\partial x_j} + \frac{g}{\theta_0} R_{w,\theta_v} - \frac{\partial R_{u_j,e}}{\partial x_j} - \frac{1}{\rho_0} \frac{\partial R_{u_j,p}}{\partial x_j} - \varepsilon, \quad (14)$$

with  $\varepsilon$  the dissipation rate. To close Eq. (14), we need to parameterize all the right-hand-side terms but the first one. The SFS-TKE production by shear (the second term) is closed with Eq. (12). Following Deardorff (1980), we use for the third term, the production due to buoyancy:

$$\frac{g}{\theta_0} R_{w,\theta_v} = \frac{g}{\theta_0} \left( A R_{w,\theta_1} + B R_{w,q_t} \right), \quad (15)$$

with coefficients  $A$  and  $B$  depending on the local thermodynamic state (dry or wet):

$$\left. \begin{aligned} A &= A_d = 1 + \frac{R_v}{R_d} \tilde{q}_t \\ B &= B_d = \left( \frac{R_v}{R_d} - 1 \right) \theta_0 \end{aligned} \right\} \text{if } q_c = 0 \quad (16)$$

$$\left. \begin{aligned} A &= A_w = \frac{(1 - \tilde{q}_t + \tilde{q}_s \frac{R_v}{R_d} (1 + \frac{L}{R_v T}))}{1 + \frac{L^2 q_s}{c_{pd} R_v T^2}} \\ B &= B_w = A_w \frac{L}{c_{pd}} - \theta_0 \end{aligned} \right\} \text{if } q_c > 0, \quad (17)$$

where  $q_s$  is the saturation specific humidity at the given temperature. At a cloud interface, it is a matter of choice whether to use the dry or the wet coefficients in calculation of the buoyancy production. Especially in situations where the properties of the cloud deck are around the buoyancy reversal criterion, this choice proves to be critical (Randall, 1980; Bretherton et al., 2004; de Roode, 2007). To determine whether a parcel that is a mixture of saturated and unsaturated air is saturated itself, we calculate the amount of unsaturated air that is needed to evaporate all the liquid water in a mixed

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air parcel. In particular the critical mixing ratio  $\chi_*$  defines the ratio of cloudy to total air mass for a just saturated mixed air parcel (Stevens, 2002):

$$\chi_* = \frac{\left(A_d \frac{L}{c_{pd}} - \frac{R_v}{R_d} \tilde{\theta}_l\right) q_c}{(A_d - A_w) \Delta \tilde{\theta}_l + (B_d - B_w) \Delta \tilde{q}_t}, \quad (18)$$

where  $\Delta \tilde{\theta}_l = \tilde{\theta}_l(z + \Delta z) - \tilde{\theta}_l(z - \Delta z)$  and  $\Delta \tilde{q}_t = \tilde{q}_t(z + \Delta z) - \tilde{q}_t(z - \Delta z)$  are the differences over the cloud interface. If  $\chi_*$  is smaller than 0.5, the dry coefficients are used to determine the SFS buoyancy production. If  $\chi_* > 0.5$ , the wet coefficients are used.

The fourth and fifth term in Eq. (14) are together parameterized as

$$-\frac{\partial}{\partial x_j} \left( R_{u_j, e} + \frac{1}{\rho_0} R_{u_j, p} \right) = \frac{\partial}{\partial x_j} \left( 2K_m \frac{\partial e}{\partial x_j} \right). \quad (19)$$

Under the assumption of 3-D homogenous isotropic stationary turbulence, and for a sharp spectral cutoff filter, a relation between the dissipation rate  $\varepsilon$  and the SGS-TKE  $e$  can be found by integration of the energy spectrum  $E(k) = \alpha \varepsilon^{2/3} k^{-5/3}$  from a filter wavenumber  $k_f$ , that lies within the inertial subrange, to infinity. This leads to

$$\varepsilon = e^{3/2} k_f \left( \frac{3}{2} \alpha \right)^{-3/2}, \quad (20)$$

with  $\alpha = 1.5$  denoting the Kolmogorov constant. Assuming a Kolmogorov turbulent spectrum between  $\kappa = 0$  and  $\kappa = k_f$ , the production of SFS-TKE due to shear is equal to:

$$\begin{aligned} \mathcal{P} &= 2K_m \int_0^{k_f} k^2 E(k) dk \\ &= \frac{3}{2} K_m \alpha \varepsilon^{2/3} k_f^{4/3}. \end{aligned} \quad (21)$$

The eddy diffusivity for momentum can be found by equating locally the shear-production of SFS-TKE to the dissipation:

$$\mathcal{P} = \varepsilon.$$

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This yields for  $K_m$

$$K_m = \frac{e^{1/2}}{k_f} \left( \frac{3}{2} \alpha \right)^{-3/2} = c_m \lambda e^{1/2},$$

$$\text{with } c_m = \frac{c_f}{2\pi} \left( \frac{3}{2} \alpha \right)^{-3/2} \quad (22)$$

where we defined the filter width  $c_f \lambda = \frac{2\pi}{k_f}$ . The eddy diffusivity for heat is modeled

5 similarly as  $K_h = c_h \lambda e^{1/2}$ , and for the dissipation  $\varepsilon$  we can write:

$$\varepsilon = \frac{c_\varepsilon}{\lambda} e^{3/2}, \text{ with } c_\varepsilon = \frac{2\pi}{c_f} \left( \frac{3}{2} \alpha \right)^{-3/2} \quad (23)$$

In unstable flow,  $c_f \lambda$  can be taken proportional to the grid size:

$$\lambda = \Delta = (\Delta x \Delta y \Delta z)^{1/3}, \quad (24)$$

$$c_f = 2.5, \quad (25)$$

10 see Cuijpers (1990). However, this no longer holds for stable situations, i.e., when  $\frac{\partial \tilde{\theta}_v}{\partial z} > 0$ . In that case,  $\lambda$  is taken to be

$$\lambda = \min \left( \Delta, c_N \frac{e^{1/2}}{N} \right), \quad (26)$$

with  $N^2 = \frac{g}{\theta_0} \frac{\partial \tilde{\theta}_v}{\partial z}$  denoting the Brunt-Väisälä frequency, and  $c_N = 0.76$ . A stability correction is also applied on  $c_h$  and  $c_\varepsilon$ :

$$15 \quad c_h = \left( c_{h,1} + c_{h,2} \frac{\lambda}{\Delta} \right) c_m, \quad (27)$$

$$c_\varepsilon = c_{\varepsilon,1} + c_{\varepsilon,2} \frac{\lambda}{\Delta}. \quad (28)$$

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Now all parameters of the subfilter-scale parameterization of DALES are defined; they are summarized in Table 2.

Substituting the closure relations and parameters into Eq. (14) obtains the prognostic equation for  $e^{1/2}$  as implemented in DALES:

$$\begin{aligned}
 5 \quad \frac{\partial e^{1/2}}{\partial t} = & -\tilde{u}_j \frac{\partial e^{1/2}}{\partial x_j} + \frac{1}{2e^{1/2}} \left[ K_m \left( \frac{\partial \tilde{u}_j}{\partial x_i} + \frac{\partial \tilde{u}_i}{\partial x_j} \right) \frac{\partial \tilde{u}_i}{\partial x_j} - K_h \frac{g}{\theta_0} \frac{\partial (A\tilde{\theta}_1 + B\tilde{q}_t)}{\partial z} \right] \\
 & + \frac{\partial}{\partial x_j} \left( 2K_m \frac{\partial e^{1/2}}{\partial x_j} \right) - \frac{c_\varepsilon e}{2\lambda}, \quad (29)
 \end{aligned}$$

which closes the system.

### 2.3.2 Smagorinsky SFS modeling

10 The way Smagorinsky is implemented in DALES, which is including a stability correction, can be derived from Eq. (29) by neglecting the transport terms and time dependency. After some algebra, we arrive at:

$$K_m = c_s^2 \lambda^2 \left( \frac{1}{2} S^2 \right)^{\frac{1}{2}} \left( 1 - \frac{\text{Ri}_g}{\text{Pr}} \right)^{\frac{1}{2}}, \quad (30)$$

with:

$$c_s = c_m^{\frac{3}{4}} c_e^{-\frac{1}{4}} = 0.22 \quad (31)$$

$$15 \quad S^2 = \left( \frac{\partial \tilde{u}_j}{\partial x_i} + \frac{\partial \tilde{u}_i}{\partial x_j} \right)^2$$

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$$Ri_g = \frac{g}{\theta_0} \frac{\partial(A\tilde{\theta}_1 + B\tilde{q}_1)}{\partial z} \frac{1}{\frac{1}{2}S^2}$$

the Smagorinsky constant, rate of strain and the gradient Richardson number, respectively.  $Pr = K_m/K_h = 0.33$  is the Prandtl number, equivalent to  $\frac{c_m}{c_h}$  in the previous section.

## 2.4 Boundary conditions: the surface model

5 The LES model contains a set of options to solve the interaction of the turbulent flow with its bottom boundary. As DALES has a no-slip boundary at the bottom, but does not resolve the flow to the scale of the molecular friction, it requires a model to parameterize the turbulent drag and the exchange of scalars between the surface and the atmosphere. The surface fluxes enter the domain at subfilter-scale, since by definition  
 10 the resolved fluctuations in the vertical velocity at the surface are equal to zero.

We followed the common way of parameterizing turbulent fluxes in atmospheric models by applying the transfer laws as given by Louis (1979):

$$u_{*0}^2 = \sqrt{\overline{u'w'}_0^2 + \overline{v'w'}_0^2} = C_m U^2 \quad (32)$$

$$\overline{w'\varphi'} = -C_\varphi U (\tilde{\varphi} - \varphi_0) \quad (33)$$

15 where  $u_{*0}$  is the friction velocity,  $u$  and  $v$  are the two horizontal components of the total horizontal wind  $U = \sqrt{\overline{u^2} + \overline{v^2}}$ ,  $\varphi_0$  is the surface value of scalar  $\varphi$ , and  $C_m$  and  $C_\varphi$  the drag coefficients for momentum and scalars, respectively.

In DALES we assume that the first model level is in the atmospheric surface layer. Therefore, we can apply the scaling arguments of Businger et al. (1971); Yaglom  
 20 (1977), who used Monin-Obukhov similarity theory to show that in this layer the following relations hold:

$$u_{*0} = \frac{\kappa U}{\ln\left(\frac{z_1}{z_{0m}}\right) - \Psi_m\left(\frac{z_1}{L}\right) + \Psi_m\left(\frac{z_{0m}}{L}\right)} \quad (34)$$

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$$\varphi_* = -\frac{\widetilde{w' \varphi'}}{u_{*0}} = \frac{\kappa (\widetilde{\varphi} - \varphi_0)}{\ln\left(\frac{z_1}{z_{0\varphi}}\right) - \Psi_\varphi\left(\frac{z_1}{L}\right) + \Psi_\varphi\left(\frac{z_{0h}}{L}\right)} \quad (35)$$

in which  $z_{0m}$  and  $z_{0\varphi}$  are the roughness lengths for momentum and an arbitrary scalar,  $z_1$  is the height of the first model level and  $U(z_1)$  and  $\widetilde{\varphi}(z_1)$  the horizontal wind and the value of the scalar at this height.  $\Psi_m$  and  $\Psi_\varphi$  are the integrated flux gradient relationships and

$$L = -\frac{u_*^3}{\kappa \frac{g}{\theta_0} \widetilde{w' \theta'_{v0}}} \quad (36)$$

is the Obukhov length, with  $\kappa=0.4$  the von Karman constant. DALES uses the functions for  $\Psi_m$  and  $\Psi_\varphi$  as provided by Beljaars (1991).

In DALES  $u_{*0}$ ,  $L$  and, optionally, the characteristic scalar scales  $\varphi_*$  and moisture scale  $q_*$  are determined iteratively, as there is a circular dependence between these variables. Per time step a slab averaged value for these variables is derived based on horizontally averaged wind and scalar gradients. Now, we can calculate mean drag coefficients  $C_m$  and  $C_\varphi$  by:

$$C_m = \frac{u_{*0}^2}{\langle U^2 \rangle} \quad (37)$$

$$C_\varphi = \frac{u_{*0} \varphi_*}{\langle U \rangle (\langle \widetilde{\varphi} \rangle - \varphi_0)} \quad (38)$$

Although all locations in the horizontal use the same drag coefficient, we calculate local fluxes by using the local values of the wind and scalars. The subfilter-scale momentum fluxes are calculated by decomposing Eq. (32) along the two components of the horizontal wind vector, whereas Eq. (33) gives the scalar flux. This results in

$$\widetilde{u' w'_0} = -C_m U \widetilde{u} \quad (39)$$

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$$\widetilde{v'w'}_0 = -C_m U \widetilde{v} \quad (40)$$

$$\widetilde{w'\varphi'} = -C_\varphi U (\widetilde{\varphi} - \varphi_0) \quad (41)$$

The generic formulation for the scalar flux Eq. (41) does not hold for evaporation over a surface where moisture is not freely available, such as a vegetated land surface or a bare soil. Here, we define the aerodynamic resistance  $r_a$  as  $(C_q U)^{-1}$  and introduce the stomatal resistance  $r_s$  in the equation to get:

$$\widetilde{w'q'} = -\frac{1}{r_a + r_s} (\widetilde{q} - q_s) \quad (42)$$

DALES has four options to calculate the fluxes at the bottom boundary:

1. *Parameterized surface scalar and momentum fluxes, parameterized surface values.* Here, a Land Surface Model (LSM, see Sect. 2.4.1) calculates the surface temperature and the stomatal resistance which enters in the evaporation equation based on the vegetation type that is assigned to the grid cell. The variables  $u_{*0}$ ,  $L$  and  $\varphi_*$  are solved iteratively to get the drag coefficients. This the option that is associated to a fully interactive land surface. Combined with the radiation model, this options allows for the simulation of full diurnal cycles, in which the surface temperature is a free variable.
2. *Parameterized surface scalar and momentum fluxes, prescribed surface values.* In this option  $u_{*0}$ ,  $L$  and  $\varphi_*$  are solved iteratively to get the drag coefficients. This option is often associated with studies of the marine boundary layer, in which the surface temperature is nearly constant during the time scale of an LES simulation. It is also applied in the simulation of stable boundary layers. For simulations over land, a fixed stomatal resistance  $r_s$  can be prescribed.
3. *Prescribed surface scalar fluxes, prescribed  $u_{*0}$ .* In this option no iterations are necessary and the scalar surface values  $\varphi_0$  are calculated diagnostically. This is an option that is commonly used in simulations of free convection where  $u_{*0}=0$ .

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4. *Prescribed surface scalar fluxes, parameterized  $u_{*0}$ .* Here  $u_{*0}$  and  $L$  are resolved by iteration, whereas  $\varphi_*$  is diagnostically calculated as a function of the prescribed scalar fluxes and the calculated  $u_{*0}$ . This is the preferred option for daytime convection over land in situations containing a large-scale pressure gradient.

5 Prescribed fluxes or surface values may depend on time; linear interpolation is then performed between the given “anchor” points.

In addition to the previous description which treated homogeneous surfaces, DALES is also able to simulate heterogeneously forced ABLs. Under such conditions, only the prescribed scalar fluxes boundary conditions are available. Each grid cell has then its own value for the scalar flux, whereas the momentum flux is dynamically computed.

### 2.4.1 Land surface model

DALES has the option to use a land surface model (LSM). The LSM consists of two parts, namely a skin layer at which the surface energy balance is solved for each separate grid cell and a four layer soil scheme which calculates the soil temperature profile for each grid cell. In the skin layer the following surface energy balance equation is solved:

$$C_{\text{sk}} \frac{dT_s}{dt} = Q_* - \rho c_p \widetilde{w' \theta'} - \rho L_v \widetilde{w' q'} - G \quad (43)$$

in which  $C_{\text{sk}}$  is the heat capacity per unit of area of the skin layer (see Duynkerke, 1999),  $T_s$  is the surface temperature,  $Q_*$  is the net radiation and  $G$  is the ground heat flux. If the LSM is used, the stomatal resistance  $r_s$  in Eq. (42) is calculated using the Jarvis-Stewart parameterization (Jarvis, 1976).

The ground heat flux is parameterized as:

$$G = \Lambda (T_s - T_{\text{soil}_1}) \quad (44)$$

in which  $\Lambda$  is a bulk conductance for the stagnant air in the skin layer (Duynkerke, 1999) and  $T_{\text{soil}_1}$  is the temperature of the highest soil layer.

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The soil consists of four layers in which the heat transport is solved using a simple diffusion equation in which both the conductivity and the heat capacity are a function of the properties of the material of which the soil is built and of the moisture content of the soil. The temperature at the bottom of the lowest soil layer is prescribed.

## 5 2.5 Boundary conditions: the sides and top

In comparison with the bottom boundary, the boundary conditions at the top of the domain and in the horizontal directions are relatively straightforward. In the horizontal directions, periodic boundary conditions are applied. At the top of the domain, we take:

$$\frac{\partial \tilde{u}}{\partial z} = \frac{\partial \tilde{v}}{\partial z} = 0; \quad \tilde{w} = 0; \quad \frac{\partial \tilde{\varphi}}{\partial z} = \text{cst.} \quad (45)$$

10 Horizontal fluctuations at the top of the domain (for instance gravity waves) are damped out by a sponge layer through an additional forcing/source term:

$$\mathcal{F}_i^{\text{sp}}(z) = \frac{1}{t^{\text{sp}}} (\langle \tilde{u}_i \rangle - \tilde{u}_i), \quad (46)$$

$$\mathcal{S}_\varphi^{\text{sp}}(z) = \frac{1}{t^{\text{sp}}} (\langle \varphi \rangle - \tilde{\varphi}), \quad (47)$$

15 with  $t^{\text{sp}}$  a relaxation time scale that goes from  $t_0^{\text{sp}} = 1/(2.75 \times 10^{-3}) \text{ s} \approx 6 \text{ min}$  at the top of the domain to infinity at the bottom of the sponge layer.

## 2.6 Numerical scheme

A Cartesian grid is used, with optional grid stretching in the  $\hat{z}$  direction. For clarity, an equidistant grid is assumed in the discussion of the advection scheme. The grid is staggered in space, on an Arakawa C-grid; the pressure, the SFS-TKE and the scalars are defined at  $\mathbf{x} + \frac{1}{2}(\Delta x, \Delta y, \Delta z)$ , the  $\tilde{u}$  is defined at  $\mathbf{x} + \frac{1}{2}(0, \Delta y, \Delta z)$ , and similar for  $\tilde{v}$  and  $\tilde{w}$ . The level of cell center is called the full level (denoted with an “f”); the level

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where  $w$  is located is called the half level (an “h”). The (variable) vertical grid spacing  $\Delta z$  is defined centered around the belonging level (see Fig. 2).

To be able to use multiple processes, thus decreasing the wall clock time of experiments, DALES 3.2 has been parallelized by dividing the domain in separate stripes in the  $y$ -direction. Tests show that this method is computationally efficient as long as the amount of processes is smaller than a quarter of the number of gridpoints in the  $y$ -direction. In the near future, we plan to also divide the domain in the  $x$ -direction, leaving narrow columns to be calculated by each process, and ensuring that the maximum number of processes would scale with the total number of grid points in each slab, thus allowing for much larger experiments.

Time integration is done by a third order Runge-Kutta scheme following Wicker and Skamarock (2002). With  $f^n(\phi^n)$  the right-hand side of the appropriate equation of Eqs. (6)–(7) for variable  $\phi = \{\tilde{u}, \tilde{v}, \tilde{w}, e^{1/2}, \tilde{\phi}\}$ ,  $\phi^{n+1}$  at  $t + \Delta t$  is calculated in three steps:

$$\begin{aligned} \phi^* &= \phi^n + \frac{\Delta t}{3} f(\phi^n) \\ \phi^{**} &= \phi^n + \frac{\Delta t}{2} f(\phi^*) \\ \phi^{n+1} &= \phi^n + \Delta t f(\phi^{**}), \end{aligned} \tag{48}$$

with the asterisks denoting intermediate time steps. The size of the timestep  $\Delta t$  is determined adaptively, and is limited by both the Courant-Friedrichs-Lewy criterion (CFL)

$$\text{CFL} = \max \left( \left| \frac{\tilde{u}_i \Delta t}{\Delta x_i} \right| \right), \tag{49}$$

and the diffusion number  $d$  (Wesseling, 1996, see).

$$d = \max \left( \sum_{i=1}^3 \frac{K_m \Delta t}{\Delta x_i^2} \right). \tag{50}$$

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The numerical stability and accuracy depends on the spatial scheme that is used. Therefore, the limiting CFL and  $d$  numbers can be adjusted to further optimize the timestep.

Depending on the desired properties (like high accuracy or monotonicity), several advection schemes are available. With advection in the  $\hat{x}$  direction discretized as

$$\frac{\partial \tilde{u}_i \phi_i}{\partial x} = \frac{F_{i+\frac{1}{2}} - F_{i-\frac{1}{2}}}{\Delta x}, \quad (51)$$

with  $F_{i-\frac{1}{2}}$  the convective flux of variable  $\phi$  through the  $i-\frac{1}{2}$  plane. Since we are using a staggered grid, the velocity is available at  $i-\frac{1}{2}$  without interpolation. Second order central differencing can be used for variables where neither very high accuracy nor strict monotonicity is necessary:

$$F_{i-\frac{1}{2}}^{2nd} = \tilde{u}_{i-\frac{1}{2}} \frac{\phi_i + \phi_{i-1}}{2}, \quad (52)$$

A higher-order accuracy in the calculation of the advection is reached with a sixth order central differencing scheme (see Wicker and Skamarock, 2002):

$$F_{i-\frac{1}{2}}^{6th} = \frac{\tilde{u}_{i-\frac{1}{2}}}{60} [37(\phi_i + \phi_{i-1}) - 8(\phi_{i+1} + \phi_{i-2}) + (\phi_{i+2} + \phi_{i-3})]. \quad (53)$$

By adding a small dissipative term to  $F_{i-\frac{1}{2}}^{6th}$ , a fifth order scheme is created that is nearly monotone:

$$F_{i-\frac{1}{2}}^{5th} = F_{i-\frac{1}{2}}^{6th} - \left| \frac{\tilde{u}_{i-\frac{1}{2}}}{60} \right| [10(\phi_i - \phi_{i-1}) - 5(\phi_{i+1} - \phi_{i-2}) + (\phi_{i+2} - \phi_{i-3})]. \quad (54)$$

For advection of scalars that need to be strictly monotone (for example chemically reacting species) the  $\kappa$  scheme (Hundsdoerfer et al., 1995) has been implemented:

$$F_{i-\frac{1}{2}}^{\kappa} = \tilde{u}_{i-\frac{1}{2}} \left[ \phi_{i-1} + \frac{1}{2} \kappa_{i-\frac{1}{2}} (\phi_{i-1} - \phi_{i-2}) \right], \quad (55)$$

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in case  $\tilde{u} > 0$ . Following Hundsdorfer et al. (1995),  $\kappa_{i-1/2}$  serves as a switch between higher order advection and first order upwind in case of strong upwind gradients of  $\phi$ . This makes the scheme monotone, but also rather dissipative.

## 2.7 Cloud microphysics

5 The cloud-microphysical scheme implemented in DALES is a bulk scheme for precipitating liquid-phase clouds. The hydrometeor spectrum is divided in a cloud droplet and rain drop category. The cloud liquid water specific humidity  $q_c$  is diagnosed using a classic saturation adjustment. The cloud droplet number concentration  $N_c$  is a fixed parameter that can be adjusted according to the degree of pollution of the cloud, meaning that the activation process is not taken in account. Two precipitation schemes have  
 10 been implemented, both 2-moment bulk schemes: rain drop spectra are characterized by the rain drop number concentration  $\tilde{N}_r$  and the rain water specific humidity  $\tilde{q}_r$ . The first one is based on Seifert and Beheng (2001, hereafter SB01) (and will be referred to SB01 scheme) two-moment bulk scheme developed for heavy precipitating warm  
 15 clouds and the second one on Khairoutdinov and Kogan (2000, hereafter KK00), valid only for stratocumulus clouds.

For each prognostic variable modified in microphysics, the source term due to microphysical processes  $\mathcal{S}^{\text{mcr}}$  consists of autoconversion (au), accretion (ac), rain drop selfcollection (sc), break-up (bu), rain sedimentation (ser), cloud droplet sedimentation  
 20 (sec), and rain evaporation (evr):

$$\begin{aligned}
 \mathcal{S}_{q_t}^{\text{mcr}} &= \mathcal{S}_{q_t}^{\text{au}} + \mathcal{S}_{q_t}^{\text{acc}} + \mathcal{S}_{q_t}^{\text{sec}} + \mathcal{S}_{q_t}^{\text{evr}} \\
 \mathcal{S}_{\theta_1}^{\text{mcr}} &= \mathcal{S}_{\theta_1}^{\text{au}} + \mathcal{S}_{\theta_1}^{\text{acc}} + \mathcal{S}_{\theta_1}^{\text{sec}} + \mathcal{S}_{\theta_1}^{\text{evr}} \\
 \mathcal{S}_{N_r}^{\text{mcr}} &= \mathcal{S}_{N_r}^{\text{au}} + \mathcal{S}_{N_r}^{\text{sc}} + \mathcal{S}_{N_r}^{\text{br}} + \mathcal{S}_{N_r}^{\text{ser}} + \mathcal{S}_{N_r}^{\text{evr}} \\
 \mathcal{S}_{q_r}^{\text{mcr}} &= \mathcal{S}_{q_r}^{\text{au}} + \mathcal{S}_{q_r}^{\text{acc}} + \mathcal{S}_{q_r}^{\text{ser}} + \mathcal{S}_{q_r}^{\text{evr}}.
 \end{aligned} \tag{56}$$

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Microphysical tendencies in  $\tilde{\theta}_1$  can be expressed directly in function of  $\tilde{q}_t$  tendencies:

$$S_{\theta_1}^{\text{mcr}} = -\frac{L}{c_{p,d}\Pi} S_{q_t}^{\text{mcr}}. \quad (57)$$

The prognostic thermodynamical variables, microphysical variables, processes and parametrizations are summarized in Fig. 3 and are described in the next sections. The conversion rates that impact rain formation and evolution are parametrized according to KK00 or according to SB01, Seifert and Beheng (2006, hereafter SB06), and Seifert (2008) (hereafter S08). The cloud water specific humidity is diagnosed from the cloud condensation and evaporation scheme.

### 2.7.1 Cloud droplet condensation and evaporation

The cloud water specific humidity  $q_c$  is diagnosed from pressure, temperature and total specific humidity using an “all or nothing” cloud adjustment scheme: it is assumed that there is no cloud water present in an unsaturated grid box, while all moisture above saturation value  $\tilde{q}_s$  is cloud water:

$$q_c = \begin{cases} \tilde{q}_t - \tilde{q}_s & \text{if } \tilde{q}_t > \tilde{q}_s \\ 0 & \text{otherwise.} \end{cases} \quad (58)$$

To calculate  $\tilde{q}_s \equiv \tilde{q}_s(\tilde{T}, p)$ , an implicit equation needs to be solved, because  $\tilde{T}$  is not directly available and has to be diagnosed from the prognostic variables  $\theta_1$  and  $q_t$ .  $\tilde{T}$  is approximated with help of the liquid water temperature  $\tilde{T}_l$ , which is equal to:

$$\tilde{T}_l = \Pi \tilde{\theta}_1. \quad (59)$$

Following Sommeria and Deardorff (1977),  $\tilde{q}_s(\tilde{T}, p)$  is found through a Taylor expansion around  $\tilde{q}_{sl} \equiv \tilde{q}_s(\tilde{T}_l, p)$ :

$$\tilde{q}_s(\tilde{T}, p) = \tilde{q}_s(\tilde{T}_l, p) + (\tilde{T} - \tilde{T}_l) \left. \frac{\partial \tilde{q}_s}{\partial \tilde{T}_l} \right|_{\tilde{T}_l = \tilde{T}} + o(\Delta \tilde{T}_l^2), \quad (60)$$

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and the higher order terms are neglected. For ideal gases, the saturation specific humidity is expressed in the saturation vapor pressure as:

$$\tilde{q}_{sl} = \frac{R_d}{R_v} \frac{e_s}{p - \left(1 - \frac{R_d}{R_v}\right) e_s}. \quad (61)$$

By convention,  $e_s$  is used to denote the saturation vapor pressure; note however, that  $e_s$  is not related to the SFS-TKE  $\tilde{e}$  as defined in Sect. 2.3. The Clausius-Clapeyron relation relates  $e_s$  to the temperature:

$$\frac{de_s}{dT} = \frac{L e_s}{R_v T^2}, \quad (62)$$

with  $R_v = 461.53 \text{ J kg}^{-1} \text{ K}^{-1}$  denoting the gas constant for water vapor. It can be solved in very good approximation as:

$$e_s(\tilde{T}_l) = e_{s0} \exp \left[ a \frac{\tilde{T}_l - T_{\text{trip}}}{\tilde{T}_l - b} \right], \quad (63)$$

with constants  $e_{s0} = 610.78 \text{ Pa}$ ,  $T_{\text{trip}} = 273.16 \text{ K}$ ,  $a = 17.27$  and  $b = 35.86$ . After having substituted in Eqs. (61)–(63) into the truncated Taylor expansion Eq. (60) we obtain for the saturated specific humidity:

$$\tilde{q}_s = \tilde{q}_{sl} \left( 1 + \frac{L^2}{R_v c_{p,d} \tilde{T}_l^2} \tilde{q}_t \right) \left( 1 + \frac{L^2}{R_v c_{p,d} \tilde{T}_l^2} \tilde{q}_{sl} \right)^{-1}, \quad (64)$$

and finally the cloud water specific humidity can be calculated with Eq. (58). If necessary, the procedure can be applied iteratively to obtain increased accuracy.

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## 2.7.2 Cloud droplet sedimentation

The cloud droplet sedimentation process has an impact on cloud evolution by reducing entrainment at stratocumulus cloud top (Ackerman et al., 2004; Bretherton et al., 2007). Its cloud water specific humidity source term can be expressed as the derivative of a sedimentation flux. The latter is parametrized by assuming a Stokes law to calculate the cloud droplets terminal velocity and a log-normal distribution to represent the cloud droplet spectrum (Ackerman et al., 2009), which lead to the following expression:

$$F_{q_c}^{\text{sec}} = k_{\text{St}} \frac{4}{3} \pi \rho_w N_c^{-2/3} q_c^{5/3} \exp(5 \ln(\sigma_{\text{gc}})^2) \quad (65)$$

with  $\rho_w = 1000 \text{ kg m}^{-3}$ ,  $k_{\text{St}} = 1.2 \times 10^8 \text{ m}^{-1} \text{ s}^{-1}$  and the log-normal geometric standard deviation parameter  $\sigma_{\text{gc}}$  is set to 1.3 (Geoffroy et al., 2009).

## 2.7.3 Rain drop processes

The precipitation parametrizations are expressed as functions of local microphysical values. Thus they are valid only for simulations where microphysical fields are explicitly resolved, as is the case in LES. Resolution must not be more than 200 m horizontally and a few ten of meters vertically.

In slightly precipitating clouds, most of the falling mass is contained in particles smaller than  $50 \mu\text{m}$  in radius, also referred to as drizzle. In Khairoutdinov and Kogan (2000) scheme, the limit between the cloud category and the rain category is set at the radius value of  $25 \mu\text{m}$  which permits consideration of drizzle in the precipitating category, which can have significant impact on the evolution of the boundary layer. This scheme is empirically based: it has been tuned with spectra derived from 3-D simulations of stratocumulus clouds using a coupled LES-bin microphysics model. Thus it is valid only for stratocumulus clouds. Because a description of that scheme is fully given in the corresponding paper, it is not described here.

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SB01 assumes the limit at the separating mass value  $x_0$  of  $2.6 \times 10^{-10}$  kg which corresponds to a separating radius  $r_0$  of the order of  $40 \mu\text{m}$ . Thus the SB01 scheme is more suitable for heavily precipitating clouds, in which most of the falling mass is contained in millimeter size particles. SB01 scheme is both theoretically and empirically based. The parametrized collection rates are expressed in function of the microphysical variables by analytical integration of the stochastic collection equation (SCE) and assuming analytical distributions to represent the hydrometeor spectra. A correction function is added to the autoconversion and accretion rate, that take in account the evolution of the cloud droplet spectra due to conversion of cloud water in rain water.

The rain drop size distribution (RDSD) is assumed to be a Gamma distribution:

$$n_r(r) = N_0 \lambda_r^{\mu_r+1} r^{\mu_r} e^{-\lambda_r r} \quad (66)$$

$N_0$  and the slope parameter  $\lambda_r$  can be expressed as a function of the prognostic variables and  $\mu_r$ . In autoconversion and accretion parametrizations,  $\mu_r$  has been set to the Marshall and Palmer (1948) value (i.e. 0) and is fixed because the parametrizations have been tuned with such a value using spectra derived from 1D simulations using a coupled LES-bin microphysics model. Other processes are analytically calculated. The value of the shape parameter  $\mu_r$  is parametrized in function of the rain water content (Geoffroy et al., 2010):

$$\mu_r = 0.5 / (1000 \rho \tilde{q}_r)^{0.7} - 1 \quad (67)$$

#### 2.7.4 Autoconversion from cloud droplets to rain drops

Autoconversion is the process that initializes the rain drop spectra. After analytical integration of the SCE and adding of correction function, SB06 obtain the following parametrized expressions:

$$S_{q_r}^{\text{au}} = \frac{k_{\text{au}}}{20x_0} \frac{(v_c+2)(v_c+4)}{(v_c+1)^2} q_c^2 x_c^2 \left( 1 + \frac{\zeta_{\text{au}}(\chi)}{(1-\chi)^2} \right) \rho_0 \quad (68)$$

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$$S_{q_t}^{\text{au}} = -S_{q_r}^{\text{au}} \quad (69)$$

with:

$$\chi_1 = \tilde{q}_r / (q_c + \tilde{q}_r) \quad (70)$$

$$\zeta_{\text{au}}(\chi_1) = 400 \chi_1^{0.7} (1 - \chi_1^{0.7})^3 \quad (71)$$

$k_{\text{au}} = 10.58 \times 10^9 \text{ m}^3 \text{ kg}^{-2} \text{ s}^{-1}$  (Pinsky and Khain, 2002, SB06) and  $x_c$  is the mean mass of the cloud droplet distribution.  $v_c$  is parametrized according to Geoffroy et al. (2009):

$$v_c = 1580 \rho q_c - 0.28 \quad (72)$$

New drizzle drops are assumed to have a radius equal to the separating radius  $r_0$ . Thus the rain number concentration source term due to autoconversion is:

$$S_{N_r}^{\text{au}} = \frac{S_{q_r}^{\text{au}}}{\frac{4\pi\rho_w r_0^3}{3\rho}} \quad (73)$$

## 2.7.5 Accretion of cloud droplets

The growth rate of rain drops by collecting cloud droplets is roughly a function of cloud and rain water contents. Cloud water specific humidity conversion rates are (SB06):

$$S_{q_r}^{\text{acc}} = k_{\text{acc}} q_c \tilde{q}_r \zeta_{\text{acc}}(\rho_0 \rho)^{\frac{1}{2}} \quad (74)$$

$$S_{q_t}^{\text{acc}} = -S_{q_r}^{\text{acc}} \quad (75)$$

with:

$$\zeta_{\text{acc}}(\chi_1) = \left( \frac{\chi_1}{\chi_1 + 5 \times 10^{-5}} \right)^4 \quad (76)$$

and  $k_{\text{acc}} = 5.25 \text{ m}^3 \text{ kg}^{-1} \text{ s}^{-1}$ .

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## 2.7.6 Rain drop selfcollection

The rain number concentration decreases because of the selfcollection process, i.e. interaction between rain drops together to form larger rain drops. Its parametrization is expressed as the following (SB06):

$$5 \quad S_{N_r}^{\text{sc}} = -k_{\text{sc}} \widetilde{N}_r \widetilde{q}_r \left(1 + \frac{\kappa_{\text{sc}}}{\lambda_r}\right)^{-9} (\rho_0 \rho)^{\frac{1}{2}} \quad (77)$$

with  $k_{\text{sc}} = 7.12 \text{ m}^3 \text{ kg}^{-1} \text{ s}^{-1}$  and  $\kappa_{\text{sc}} = 60.7 \text{ kg}^{-1/3}$ .

## 2.7.7 Break-up of rain drops

The break-up of rain drops into smaller rain drops is applied for spectra with a mean volume radius  $r_{\text{vr}}$  larger than  $150 \mu\text{m}$  following (SB06):

$$10 \quad S_{N_r}^{\text{br}} = -S_{N_r}^{\text{sc}} (k_{\text{br}}(r_{\text{vr}} - r_{\text{eq}}) + 1) \quad (78)$$

with  $k_{\text{br}} = 2000 \text{ m}^{-1}$  and  $r_{\text{eq}} = 550 \mu\text{m}$ . When  $r_{\text{vr}}$  becomes larger than  $r_{\text{eq}}$  the break-up process becomes predominant over the selfcollection process. The strong increase of the break-up process for large mean volume radius is not taken in account.

## 2.7.8 Rain drop sedimentation

15 Assuming the Rogers et al. (1993) dependency of rain drop terminal velocity in function of the drop radius, the flux of the rain number concentration and the flux of the rain water specific humidity are Stevens and Seifert (2008):

$$F_{N_r}^{\text{ser}} = \left(a - b(1 + c/\lambda_r)^{\mu_r + 1}\right) N_r \quad (79)$$

$$F_{q_r}^{\text{ser}} = \left(a - b(1 + c/\lambda_r)^{\mu_r + 4}\right) q_r \quad (80)$$

20 with  $a = 9.65 \text{ m s}^{-1}$ ,  $b = 9.8 \text{ m s}^{-1}$ ,  $c = 1200 \text{ m}^{-1}$  (see S08).

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## 2.7.9 Rain drop evaporation

The tendency of the rain water specific humidity due to evaporation is expressed by integration of the drop growth rate by vapor diffusion formulation (S08):

$$\begin{aligned}
 S_{q_r}^{\text{evr}} &= 4\pi \frac{\rho_w}{\rho} G(T, P) \frac{\tilde{q}_t - \tilde{q}_s (N_r \lambda_r)^{\mu_r + 1}}{\tilde{q}_s \Gamma(\mu_r + 1)} \\
 &\times \left[ a_v \Gamma(\mu_r + 2) \lambda^{-(\mu_r + 2)} + b_v \text{Sc}^{\frac{1}{3}} \left( \frac{a}{v_a} \right)^{1/2} \Gamma\left(\mu_r + \frac{5}{2}\right) \right. \\
 &\times \left. \lambda^{-(\mu_r + \frac{5}{2})} \left( 1 - \frac{1}{2} \frac{b}{a} \left( \frac{\lambda_r}{c + \lambda_r} \right)^{(\mu_r + \frac{5}{2})} - \dots \right) \right] \quad (81)
 \end{aligned}$$

where

$$G(T, P) = \frac{1}{\rho_w} \left[ \frac{R_v T}{e_{vs}(T) D_v} + \frac{L}{k_a T} \left( \frac{L}{R_v T} - 1 \right) \right]^{-1}, \quad (82)$$

Sc, the Schmidt number,  $a_v$  and  $b_v$  are ventilation factor coefficients with the following values:  $\text{Sc}=0.71$ ,  $a_v=0.78$  and  $b_v=0.308$ .

The tendency of the rain drop number concentration due to evaporation is assumed to be (S08):

$$S_{N_r}^{\text{evr}} = \gamma \frac{\tilde{N}_r}{\tilde{q}_r} S_{q_r}^{\text{evr}} \quad (83)$$

with  $\gamma=0.7$  (A. Seifert, personal communication, 2008). Note that a 0 value of  $\gamma$  means that no rain drop disappear during evaporation. A value larger than 1 would be possible if a large number of little rain drops totally evaporate in presence of large drops.

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## 2.8 Radiation Schemes

The net radiative heating consists of the (downward pointing) radiative flux divergence integrated over all wavelengths  $\nu$ :

$$S_{\theta_1}^{\text{rad}} = \int_0^\infty \frac{\partial F^{\text{rad}}(\nu)}{\partial z} d\nu \quad (84)$$

Two approaches towards radiation modeling are implemented in DALES: Fully resolving the radiative transfer, or parameterizing the vertical component of the longwave radiation and of the shortwave radiation through computationally cheap analytic approximations of the Mie theory, that maintain sufficient accuracy for most purposes. In the parameterized radiation scheme, radiative transfer is computed at every single column of the LES model, neglecting horizontal radiative transfer.

### 2.8.1 Full radiation

Due to the integration over many wavebands, running a full radiation code is generally to costly in a tool that lies emphasis on solving the turbulent flow of the ABL, rather than solving the radiative aspects. To reduce the costs, radiative computations can be performed at a coarser spatial or temporal resolution. However, this can introduce biased errors. In DALES the Monte Carlo Spectral Integration (Pincus and Stevens, 2008) is followed, where at each grid point and at each time step the radiative flux is approximated by the radiative flux of one randomly chosen waveband, or even a randomly chosen part of that waveband where all absorption coefficients are similar. The radiative transfer model as described by Fu and Liou (1992); Fu et al. (1997) is used for solving the radiative flux. This incorporates a parametrization for the cloud water optical properties. To calculate the radiative effects of gasses, a k-distribution is used. Radiative transfer is computed with a  $\delta$  four-stream solver in both the infrared and solar parts of the spectrum.

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## 2.8.2 Parameterized longwave radiation

For longwave radiation the absorptivity is controlled by the liquid water path (LWP),

$$\text{LWP}(x, y, z_1, z_2) = \rho_{\text{air}} \int_{z_1}^{z_2} q_c(x, y, z) dz, \quad (85)$$

The net longwave radiative flux  $F_L^{\text{rad}}$  is linked to the liquid water path through an analytic formula,

$$F_L^{\text{rad}}(x, y, z) = F(z_{\text{top}}) e^{-k \text{LWP}(x, y, z, z_{\text{top}})} + F(0) e^{-k \text{LWP}(x, y, 0, z)} \quad (86)$$

where  $k$  is the absorption coefficient, and  $F(z_{\text{top}})$  and  $F(0)$  represent the total net longwave radiative flux divergence at the top of the cloud and the cloud base, respectively. Larson et al. (2007) discuss the validity of this parameterization in detail. They conclude that when the parameterization constants are optimized for individual stratocumulus cases like the ones set up by Duynkerke et al. (1999, 2004), and Stevens et al. (2005), the formula can yield remarkably accurate fluxes and heating rates.

To study the role of longwave radiative cooling on mixed-layer turbulence, but in the absence of latent heat release effects that occur in a real liquid water cloud, one can add a passive scalar field to the model. This so-called “smoke” cloud has an initial concentration set to unity in the boundary layer and zero above (Bretherton et al., 1999b). The liquid water path in the longwave radiation Eq. (86) is then replaced by the smoke path, which can be computed by substituting  $q_c$  by the smoke concentration  $s$  in Eq. (85). For a smoke absorptivity  $k=0.02 \text{ m}^2 \text{ kg}^{-1}$  one obtains similar cooling rates as in stratocumulus (Bretherton et al., 1999b). It should be noted that unlike liquid water, smoke is a conserved quantity. This means that if smoke is transported by turbulence into the inversion layer, it will cause a local cooling tendency in this layer.

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### 2.8.3 Parameterized shortwave radiation

In the shortwave band the cloud optical depth  $\tau$  is the most important parameter defining the radiative properties of clouds,

$$\tau(x, y, z) = \frac{3 \text{LWP}(x, y, z, z_t)}{2 r_e \rho_w}. \quad (87)$$

5 Here  $r_e$  defines the cloud droplet effective radius, i.e. the ratio of the third moment to the second moment of the droplet size distribution (Stephens, 1984). Although  $r_e$  depends on the height in the cloud layer and has a maximum value at the top of the cloud, a constant value is used. A typical number for marine boundary layer clouds is  $r_e = 10 \mu\text{m}$ , which was observed for stratocumulus over the Pacific Ocean off the coast of California during FIRE I (Duda et al., 1991).

10 Cloud droplets scatter most of the incident radiation into the forward direction. This asymmetry in the distribution of the scattering angle is measured by the first moment of the phase function, and is commonly referred to as the asymmetry factor  $g$  which is taken  $g = 0.85$ . The radiative transfer for shortwave radiation in clouds is modeled by the delta-Eddington approximation (Joseph et al., 1976). In this approach the highly asymmetric phase function is approximated by a Dirac delta function and a two term expansion of the phase function.

15 The ratio of the scattering coefficient  $Q_s$  to the extinction coefficient  $Q_e$  is called the single scattering albedo  $\omega_0 = Q_s/Q_e$ , and is unity for a non-absorbing medium. Following Fouquart (1985),

$$\omega_0 = 1 - 9 \times 10^{-4} - 2.75 \times 10^{-3} (\mu_0 + 1) e^{-0.09\tau_t}, \quad (88)$$

with  $\tau_t$  the total optical depth in a subcloud column. Although this expression gives single scattering albedos that are very close to unity, absorption in boundary layer clouds can not be neglected due to the large number of scattering events.

25 The delta-Eddington equations are exactly the same as the Eddington equations (Joseph et al., 1976) with transformed asymmetry factor  $g$ , single-scattering albedo  $\omega_0$

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and optical depth  $\tau$ :

$$g' = \frac{g}{1+g} \quad (89)$$

$$\omega'_0 = \frac{(1-g^2)\omega_0}{1-\omega_0 g^2} \quad (90)$$

$$\tau' = (1-\omega_0 g^2)\tau \quad (91)$$

- 5 For constant  $\omega_0$  and  $g$  the delta-Eddington equation can be solved analytically (Shettle and Weinman, 1970; Joseph et al., 1976):

$$F_s^{\text{rad}}(x, y, z) = F_0 \frac{4}{3} \left[ \rho (C_1 e^{-k\tau'(x, y, z)} - C_2 e^{k\tau'(x, y, z)}) - \beta e^{-\frac{\tau'(x, y, z)}{\mu_0}} \right] + \mu_0 F_0 e^{-\frac{\tau'(x, y, z)}{\mu_0}} \quad (92)$$

with:

$$k = [3(1-\omega'_0)(1-\omega'_0 g')]^{1/2}, \quad (93)$$

$$10 \quad \rho = \left( \frac{3(1-\omega'_0)}{1-\omega'_0 g'} \right)^{1/2}, \quad (94)$$

$$\beta = 3\omega'_0 \mu_0 \frac{1+3g'(1-\omega'_0)\mu_0^2}{4(1-k^2\mu_0^2)}, \quad (95)$$

and  $\mu_0 = \cos\alpha_0$  for a solar zenith angle  $\alpha_0$ . The values of the constants  $C_1$  and  $C_2$  in Eq. (92) are calculated from the boundary conditions. A prescribed value for the total downward solar radiation (parallel to the beam) determines the upper boundary condition at the top of the cloud  $F_0$ . In addition, it is assumed that at the ground surface a fraction of the downward radiation reaching is reflected back by a Lambertian ground surface with albedo  $A_g$ . See for further details Shettle and Weinman (1970) and Joseph et al. (1976). The delta-Eddington solution is applied in every column using the local cloud optical depth. A study by de Roode and Los (2008) on the cloud albedo bias

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effect showed a good agreement between results obtained with the delta-Eddington approach and from the I3RC Monte-Carlo model (Cahalan et al., 2005) that utilizes the full three-dimensional structure of the cloud field.

## 2.9 Other forcings and sources

5 Large-scale forcings, such as the mean geostrophic wind  $\mathbf{u}_g$ , the large-scale subsidence  $w_s$ , and the horizontal advective scalar transport can be applied through forcings and sources that can be dependent on height and time. The effects of large-scale subsidence are enforced on the slab-averaged scalar profiles through a prescribed subsidence velocity  $w_s(z, t)$ :

$$10 \quad \mathcal{S}_\varphi^{\text{subs}} = -w_s \frac{\partial \langle \tilde{\varphi} \rangle}{\partial z} \quad (96)$$

Optionally, the slab-averaged prognostic variables can be nudged with a relaxation time scale  $t^{\text{rel}}$  to a prescribed (time depending) value  $\varphi^{\text{rel}}$ :

$$\mathcal{S}_\varphi^{\text{rel}} = -\frac{1}{t^{\text{rel}}} \left( \langle \varphi \rangle - \varphi^{\text{rel}} \right), \quad (97)$$

15 analogous to large-scale forcings in single column models. The application of  $\mathcal{S}_\varphi^{\text{rel}}$  to the horizontal mean  $\langle \varphi \rangle$ , instead of to the individual values of  $\varphi$ , ensures that room for variability within the LES domain remains, and the small-scale turbulence will not be disturbed by the nudging.

## 2.10 Flow over tilted surfaces

20 To simulate flow over a sloped surface under an angle  $\alpha$  ( $>0$ ), a coordinate transformation is performed; computations are then done in a system  $(s, y, n)$ , with  $s$  and  $n$  are the coordinates along and perpendicular to the slope, respectively. Under the assumption that the flow can be considered homogeneous along the slope (see Sect. 3.5), only the

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buoyancy force is directly dependent on  $s$ . To account for this, an additional forcing is introduced:

$$\mathcal{F}_{u_s}^{\text{slope}} = -\frac{g}{\theta_0}(\theta_v - \theta_0)\sin\alpha \quad (98)$$

$$\mathcal{F}_{u_n}^{\text{slope}} = \frac{g}{\theta_0}(\theta_v - \theta_0)(\cos\alpha - 1) \quad (99)$$

5 To accommodate the periodic horizontal boundary conditions for slope flow, we follow Schumann (1990) in splitting each scalar field  $\varphi$  in an ambient component  $\varphi_a$  that incorporates the  $z$  dependency of the mean state, and a deviation  $\varphi_d$  with respect to  $\varphi_a$ .

$$\varphi = \varphi_a + \varphi_d \quad (100)$$

10 Given a mean height depending profile  $\Phi(z)$ ,

$$\varphi_a = (n\cos\alpha - s\sin\alpha)\Phi(z). \quad (101)$$

The deviation  $\varphi_d$  is now homogeneous in the horizontal direction, and periodic boundary conditions can be applied on it. Currently, this splitting procedure is only implemented in DALES for the liquid potential temperature  $\theta_l$ , focussing slope flow studies  
 15 exclusively on the dry boundary layer for now.

## 2.11 Chemically reactive scalars

DALES is equipped with the necessary tools to study the dispersion of atmospheric compounds using the Eulerian and Lagrangian framework and their chemical transformation. The Lagrangian framework is explained in Sect. 2.12.2. In the Eulerian approach, a line or surface source of a passive or a reactant is included to mimic  
 20 the emission of an atmospheric compound in the ABL flow allowing the calculation and analysis of the diagnostic scalar fields (Nieuwstadt and de Valk, 1987). If the

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atmospheric compounds react, the source or sink term in Eq. (7) needs to be included in the numerical calculation. For a generic compound  $\varphi_1$ , this reaction term reads:

$$S_{\varphi_1} = \mathcal{P}(t, \varphi_m) - \mathcal{L}(t, \varphi_m)\varphi_1 \quad m = 1, \dots, n. \quad (102)$$

The respective terms  $\mathcal{P}(t, \varphi_m)$  and  $\mathcal{L}(t, \varphi_m)$  are nonnegative and represent production and loss terms for atmospheric compound  $\varphi_1$  reacting on time  $t$  with the  $n$  number of species  $\varphi_m$  it is reacting with.

In DALES, we solve the term  $S_{\varphi_1}$  using the chemical solver TWOSTEP extensively described and tested by Verwer (1994) and Verwer and Simpson (1995). In short, this chemical solver is an implicit method with second-order accuracy based on the two-step backward differentiation formula. Since in atmospheric chemistry we are dealing with chemical system characterized by a wide range of chemical time scales, *i.e.*, *stiff system of ordinary differential equations*, the two-step solver is able to adjust the time step depending on the chemical reaction rate.

A simple chemical mechanism can serve us as an introduction of the specific form of  $\mathcal{P}(t, \varphi_k)$  and  $\mathcal{L}(t, \varphi_k)$ . Atmospheric chemistry mechanism are composed by first- and second-order reactions. Third-order reactions normally involve water vapor or an air molecule, for instance nitrogen or oxygen. Due to the much larger concentration of these compounds than the reactant concentration, third-order reaction rates are normally expressed as a pseudo second-order reaction, *i.e.*,  $k_{2\text{nd}} = k_{3\text{rd}}[M]$  where  $[M]$  is a molecule of  $\text{H}_2\text{O}$  or air. Therefore, a generic atmospheric chemical mechanism composed by a first- and a second-order reaction reads:



where  $a$ ,  $b$  and  $c$  are atmospheric compound concentrations,  $j$  and  $k$  are the first- and second-order reaction rate. For reactant  $a$  the  $\mathcal{L}$  and  $\mathcal{P}$  are, respectively:

$$\mathcal{L} = -j \quad (103)$$

$$\mathcal{P} = kbc. \quad (104)$$

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The photodissociation rate  $j$  depends on the ultraviolet actinic flux and specific photodissociation properties of the atmospheric compound. Therefore, in DALES  $j$  is a function on the diurnal variability (latitude, day of the year) and the presence of clouds. The  $j$ -values are updated every time step. The cloud influence on the actinic flux is implemented using a function that depends on the cloud optical depth (Eq. 87) (Vilà-Guerau de Arellano et al., 2005). The reaction rate  $k$  depends on the absolute temperature, on the water vapor content and the pressure. Depending on the reaction, several reaction rate expressions can be specified at DALES. Moreover, the generally very low concentrations of chemical species in the atmosphere allows us to neglect the heating contribution of the reactions on the liquid water potential temperature  $\tilde{\theta}_l$ , or on the water content  $\tilde{q}_t$  and  $\tilde{q}_r$ .

For the chemical solver, it is essential that the concentration of the species is non-negative. Therefore, the entire numerical discretization for the reactants, spatial and temporal integration of advection and diffusion and temporal integration of the chemistry, has to satisfy the following three numerical properties: it has to be conservative, monotone and positive definite. Of the advection schemes that are implemented in DALES, the kappa scheme fulfills these properties.

The chemistry module is designed to be very flexible in order to allow study of different chemical mechanisms. Required input parameters include the number of inert scalars, and of chemical species, their initial vertical profiles and surface fluxes, and a list of chemical reactions, together with the reaction rate functions. More information on the chemistry module can be found at Vilà-Guerau de Arellano et al. (2005) and Vilà-Guerau de Arellano et al. (2009).

## 2.12 Statistics

In DALES, standard output includes time series and slab-averaged profiles of the main variables, the (co-) variances, and of the resolved and SFS-modeled fluxes. The modular set-up of the code facilitates inclusion of many other statistical routines, specifically aimed at the purposes of a particular research question. Sharing such code with the

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community leaves the code base with a rich palette of statistics, including specific routines that focus on the details of, for example, radiation, cloud microphysics, or the surface layer. Although not exhaustive, a few examples of the statistical capabilities of DALES are given below.

### 5 2.12.1 Conditional sampling

Conditionally averaged profiles can be found by defining a mask  $M$ , which is equal to 1 or 0, depending on whether a set condition is true or false, respectively. Frequently used sampling conditions are, for instance, clouds ( $q_1 > 0$ ), areas of updrafts ( $\tilde{w} > 0$ ), areas of positive buoyancy ( $\tilde{\theta}_v > \langle \theta_v \rangle$ ), and any combination of these conditions. New definitions of the mask  $M$  are possible with small adjustments of the code.

### 10 2.12.2 Lagrangian statistics

While the Eulerian formulation of the LES favors a Eulerian frame of reference for statistics, many problems can greatly benefit from a Lagrangian point of view. This holds in particular for studies of entrainment and detrainment, since these problems can often be stated as a study on the past and the future of a parcel of in-cloud air. To this end a Lagrangian Particle Dispersion Model (LPDM) has been implemented into DALES. Within this module, massless particles move along with the flow. Since each of the particles is uniquely identifiable, the origins and headings of the particles (and of the air) can be captured.

20 The position of a particle  $\mathbf{x}_p$  is determined using:

$$\frac{dx_{i,p}}{dt} = \tilde{u}_i(\mathbf{x}_p; t) + u'_i(\mathbf{x}_p; t), \quad (105)$$

where  $\tilde{\mathbf{u}}$  is the LES-resolved velocity linearly interpolated to the particle position, and  $\mathbf{u}'$  is an additional random term that represents the SFS-velocity contribution. This term

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is especially important in regions where the SFS-TKE is relatively large, such as near the surface or in the inversion zone. The calculation of  $\mathbf{u}'$  follows Weil et al. (2004), and was tailored for use in LES with TKE-closure. It is implemented in DALES as follows:

$$du'_i = -\frac{3f_s C_0 \epsilon u'_i}{4e} dt + \frac{1}{2} \left( \frac{u'_i}{e} \frac{de}{dt} + \frac{2}{3} \frac{\partial e}{\partial x_i} \right) dt + (f_s C_0 \epsilon)^{1/2} d\xi_i. \quad (106)$$

$C_0$  is the Langevin-model constant (Thomson, 1987) and has been set to 6;  $f_s$  is the slab-averaged ratio between SFS-TKE and total TKE.  $d\xi$  is a Gaussian noise to mimic the velocity field associated with the subfilter turbulence.

Boundary conditions are periodic in the horizontal directions, and emulate the LES boundary conditions at the top and bottom of the domain. Particles are reflected ( $w_p$  changes sign) should they hit the top or bottom. For time integration, the third order Runge Kutta scheme is again used, but a significant dependency on the choice of the time integration scheme has not been observed. The LPDM was validated by Heus et al. (2008) for a cumulus topped boundary layer and additionally by Verzijlbergh et al. (2009) for a scalar point source emission in different clear and cloud-topped flows.

### 2.12.3 Transport, tendencies and turbulence

To study the mechanisms driving the development of the ABL, tendency statistics are included that diagnose slab average profiles of every forcing and source term in Eqs. (6) and (7). Where necessary, the individual terms of the underlying equations can also be diagnosed, such as for the SFS-TKE, radiation or microphysical components. Fluxes and co-variances of the main variables are also calculated.

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For a full understanding of the turbulence in the boundary layer, a key budget is that of the resolved turbulent kinetic energy  $E$ , which is:

$$\begin{aligned}
 \left\langle \frac{\partial E}{\partial t} \right\rangle &= \left\langle \frac{\partial}{\partial t} \left[ \frac{1}{2} (\tilde{u}''^2 + \tilde{v}''^2 + \tilde{w}''^2) \right] \right\rangle \\
 &= - \left[ \langle \tilde{u}'' \tilde{w}'' \rangle \frac{\partial \langle \tilde{u}'' \rangle}{\partial z} + \langle \tilde{v}'' \tilde{w}'' \rangle \frac{\partial \langle \tilde{v}'' \rangle}{\partial z} \right] + \frac{g}{\theta_0} \langle \tilde{w}'' \theta_v'' \rangle - \frac{\partial \langle \tilde{w}'' E \rangle}{\partial z} \\
 &\quad - \frac{1}{\rho_0} \frac{\partial \langle \tilde{w}'' \pi \rangle}{\partial z} - \varepsilon_r,
 \end{aligned} \tag{107}$$

with the double primed variables signifying the deviation from the slab average,  $\varepsilon_r$  the dissipation from the resolved scales and in absence of other forcings. Due to the staggered grid used in DALES each variable entering in the budget terms is evaluated at a different position. In order to correctly build up the different terms, several interpo-

lations have to be performed, which have to be consistent with the spatial discretization of the model. Due to these numerical issues, the budget is not fully closed, although the residual is small compared to the physical terms (see Fig. 4 for the budget of  $E$  in a sheared CBL). In order to further reduce this residual term, a method based on Gao et al. (1994) is currently in development.

### 3 Applications and evaluation of DALES

#### 3.1 Dry boundary layers

##### 3.1.1 Convective boundary layer

One of the most elementary test-cases for an atmospheric LES is the dry convective boundary layer (CBL). In a CBL a positive heat flux at the surface destabilizes the air resulting in a vigorous turbulence which mixes (thermo)dynamic quantities like heat and

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momentum over the entire depth of the boundary layer, and which comprises eddies that vary over a wide range of scales, i.e. from the depth of the boundary layer ( $\sim$ km) down to the Kolmogorov-scale ( $\sim$ mm). But because the largest scales of motion control most of the vertical transport (e.g. the vertical fluxes of heat and momentum), it is reasonable to fully resolve the large scales on a resolution of  $\sim$ 10–100 m, and account for the scales of motion smaller than the grid scale using the subgrid model (such as Eq. 29).

Probably the most defining feature of a CBL is the fact that the mixed-layer is not confined by a rigid lid (such as Rayleigh-Bénard convection), but that it is capped by an inversion, a sudden strong increase of the potential temperature. As such the mixed-layer depth  $z_i$  is not fixed, but grows in time: thermals impinging on the inversion cause overlying free tropospheric air to be entrained into the mixed-layer, the depth of which therefore increases. The rate of growth is called the entrainment rate  $w_e$ , a key unknown in weather, climate and air quality models. Large-Eddy Simulation provides a powerful tool to make a comprehensive study of entrainment (see e.g., Sullivan et al., 1998; Fedorovich et al., 2004a) and investigate the dependencies on for example the inversion jump  $\Delta\theta_v$ , the surface heat flux  $\overline{w'\theta'}$  and the actual mixed-layer depth  $z_i(t)$ . Rather than studying the entrainment rate directly, one can also focus on the entrainment *flux* of heat, in particular the value of the heat flux at the inversion. This approach is followed below for DALES.

To test the performance of DALES for dry convective boundary layers, we simulated two of the cases studied by Sullivan et al. (1998), one with a weak inversion  $\Delta\theta_v \sim 0.5$  K (their case number W06), and one with a strong inversion  $\Delta\theta_v \sim 5$  K (case S24). The corresponding surface heat flux, initial mixed-layer depth  $z_i(0)$  and stratification  $d\langle\theta_v\rangle/dz$  of the overlying layer, are given in Table 3. In both cases there is no mean wind and hence no (mean) shear. Note that W06 was initiated without an inversion jump. For S24 the initial inversion thickness amounted to 120 m (linear interpolation between 300 K and 308 K over 120 m). Both simulations were conducted on a grid of  $N_x=N_y=64, N_z=96$ , using the same resolution as in the original simula-

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tions,  $\Delta x = \Delta y = 100$  m,  $\Delta z = 20$  m. Time-step was variable, and for the advection of all variables the fifth-order scheme (see Sect. 2.6) was chosen.

In Fig. 5 we present the results averaged from hour 3 to 4. Turbulence statistics are normalized using the convective velocity scale

$$w_* = \left( \frac{g}{\Theta_0} \widetilde{w'\theta'} z_i \right)^{1/3} \quad (108)$$

where  $\widetilde{w'\theta'}$  is the surface heat flux in K m/s, and  $z_i$  the actual depth of the mixed layer (see Table 3). The figures are formatted such that they can be directly compared with the original study by Sullivan et al. (1998). Although W06 was initiated without an inversion, the CBL dynamics is such that it creates its own inversion, as can be seen in Fig. 5a showing the characteristic “steepening” of the temperature profile in the entrainment zone. The strength of the resulting inversion is the same as observed by Sullivan et al. (1998). The same holds for case S24 (Fig. 5d). For both cases also the normalized heat flux profiles display the usual negative value of roughly  $-0.15$  in the entrainment zone, indicative of the entrainment process (Fig. 5b,e). The SFS contribution to the heat flux is rather small in the mixed-layer and near the inversion. The SFS contribution to TKE, on the other hand, extends over the entire layer (Fig. 5c,f); again the magnitude and shape of the SFS-TKE are in very good agreement with the results reported by Sullivan et al. (1998).

### 3.1.2 Sheared convective boundary layer

To analyze the influence of wind-shear characteristics on the evolution of the CBL, long simulations and large domains are necessary to fulfill a quasi-stationary flow pattern that matches with the prescribed surface fluxes, and to resolve the expected pattern for forced convection (Khanna and Brasseur, 1998). With DALES, resolutions up to 25 m and 6 m in the horizontal and vertical directions, respectively were considered.

The studies of the sheared CBL focus on the influence of the wind shear on the boundary layer growth due to the modification of the entrainment fluxes (Pino et al.,

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2003); on identification and parameterization of the main physical mechanisms that control the entrainment heat flux (Kim et al., 2006; Pino et al., 2006b); on the role of shear and the inversion strength in the decay of convective turbulence during sunset (Pino et al., 2006a); and most recently on how to parameterize the different terms of the TKE budget by means a first order jump mixed layer model (Pino and Vilà-Guerau de Arellano, 2008). In an intercomparison study of the sheared CBL in different wind regimes by Fedorovich et al. (2004b), a previous version of the model showed larger entrainment fluxes than the other codes in the intercomparison, producing a warmer and dryer boundary layer. In comparison with this older version of DALES, DALES 3.2 shows smaller entrainment fluxes, due to the improved numerical scheme (see Fig. 6).

Among the results mentioned above we would like to emphasize first the influence of the shear in the boundary layer growth by using LES and observations (Pino et al., 2003), and second the influence of the wind shear in the characteristics length scales during afternoon decaying convective turbulence (Pino et al., 2006a). It was shown there that the enhancement of the entrainment heat flux caused by the wind shear at the inversion zone is responsible for an increased boundary layer height. Neglecting this wind shear would results in a significant underestimation in parameterizations of the boundary layer height (see Fig. 7).

### 3.2 Stable boundary layers

In the context of LES, one of the characteristics of stable boundary layers (SBLs) is the mere absence of large eddies (see e.g. the spectra presented in Kaimal and Finnigan, 1994). The stable stratification suppresses vertical motion and transfers turbulent kinetic energy into potential energy. Part of that potential energy is released back as turbulent kinetic energy but part is dissipated through the dissipation of temperature variance. Due to these two aspects, the role of the subfilter-scale model tends to be much larger in LES of SBLs than it is for convective or neutral (but sheared) boundary layers. This implies that for the SBL generally much higher resolutions are used than for other simulations.

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The first application of (a previous version of) DALES to stable boundary layers was reported by Galmarini et al. (1998) where a slightly different version of the subfilter-scale model was used.

In the context of the GEWEX Atmospheric Boundary Layer Study (GABLS, Holtslag, 2006), a series of model intercomparisons has been organized for SBL cases. In all intercomparisons a single-column model intercomparison case was defined, whereas an LES case was defined in the first and third intercomparison. The first case (Beare et al., 2006) was inspired by the setup of the simulations of Kosović and Curry (2000): an arctic moderately SBL (with  $z_i/L \approx 2$ , where  $z_i$  is the height of the SBL and  $L$  the Obukhov length given in Eq. 36). The domain size was set to 400 m in all three directions. The roughness length  $z_0$  was set to 0.1 m. For heat the same roughness length was applied, and a constant cooling rate of 0.25 K/h for the surface temperature.

In total 11 models participated in the intercomparison, being run at resolutions from 12.5 m down to 1 m for some models. DALES participated in the intercomparison at resolutions of 12.5 and 6.25 m. For coarse resolutions the subfilter-scale model plays an important role. The results are shown in Fig. 8. The results of DALES are clearly within the range of the other models, although the mean shear is stronger than in most models close to the surface and weaker at higher levels in the SBL. Furthermore, the strength of the low-level jet seems to be slightly less than in the other models.

### 3.3 Cloud topped boundary layer

If there is sufficient moisture in the convective boundary so that the total specific humidity  $q_t$  exceeds its saturation value  $q_s$ , condensation processes will initiate and clouds will start to form. Since  $q_s$  increases exponentially with temperature and as temperature decreases with 10 K/km in the convective boundary layer, clouds typically start to form at the top of the convective boundary layer. They are often referred to as boundary layer clouds, as long as the capping inversion at the top of the boundary layer is strong enough to encapsulate them. As a result they have a limited vertical extend of around

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3 km which makes the use of LES as a virtual laboratory highly suitable to study the dynamics of boundary layer clouds.

Stratocumulus and shallow cumulus are the two main types of boundary layer clouds that have been simulated extensively in the past with DALES and the schematics of these different types of boundary clouds are depicted in Fig. 9.

Stratocumulus clouds are low-lying, stratiform clouds often covering the sky completely, with a thickness of only several hundreds of meters, capped by a strong inversion. The turbulence that maintains the well-mixed profiles of the conserved variables  $q_t$  and  $\theta_t$  is mainly driven from the top of the stratocumulus deck due the longwave radiative cooling in addition to local cooling and heating due to condensation and evaporation of cloud droplets.

In contrast, shallow cumulus clouds occur as a population of separated small cauliflower shaped clouds with a cloud base height at around 1 km and a maximum vertical extend of around 2 km. These clouds generally only cover 10 to 30% of the sky. Shallow cumulus clouds usually form on top of the dry rising thermals in the sub-cloud layer and are dynamically characterized by strong vertical motions due to the condensational heating resulting in inner cloud cores that are positively buoyant with respect to the (dry) environment. As a result the stratification of the mean profile of  $\langle\theta_v\rangle$  is stable with respect to vertical displacements of unsaturated test parcels and unstable with respect to saturated test parcels. This effect, often referred to as conditional instability, has no counterpart in any other part of convection and is responsible for the strong intermittant behaviour of cumulus updrafts.

Although stratocumulus might appear, due to its well-mixed character, conceptually simpler than shallow cumulus, it is actually harder to simulate stratocumulus clouds in an LES model than to simulate shallow cumulus clouds. This is due to the strong inversion at the top of the stratocumulus deck, where temperature jumps of 10 K over 100 m are not uncommon. Such strong inversions result from the radiative cooling and are difficult to resolve with LES techniques, resulting in unwanted numerical diffusion over this interface which can dominate the transport over the inversion interface. On

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the other hand, in the case of shallow cumulus clouds the interaction with the radiation is not so strong due to the low cloud fraction. As a result, shallow cumulus clouds are not topped by such strong inversions which simplifies the numerical simulations. Another related simplifying factor is that because of the low cloud fraction the interaction between the clouds and the radiation is not so critical that an interactive treatment of both processes would be essential.

DALES has participated in numerous LES intercomparison studies organized over the last 15 years by the GEWEX Cloud System Studies (GCSS). These intercomparison studies have been set up to serve several purposes. It provides a critical evaluation of the participating LES codes and moreover it provides unique data sets to obtain further insights in the dynamics of the cloud topped boundary layer. More specifically these LES data sets have helped in improving the parameterized formulation of these processes in large scale Numerical Weather Prediction (NWP) and Climate models. In the coming 2 subsections examples are presented how research with DALES have contributed to the improved knowledge of the physics and dynamics of shallow cumulus and stratocumulus.

### 3.3.1 Stratocumulus

One of the most critical phenomena in the dynamics of stratocumulus is the entrainment of dry air at the top of the cloud layer. Following the flux-jump relation (Lilly, 1968), the entrainment rate ( $w_e$ ) determines the turbulent flux at the top of the boundary layer ( $\langle \widetilde{w' \varphi'} \rangle_e$ ),

$$\langle \widetilde{w' \varphi'} \rangle_e = -w_e \Delta \langle \varphi \rangle \quad (109)$$

with  $\Delta \langle \varphi \rangle$  the jump across the inversion. This equation is valid for an infinitesimally thin inversion layer and shows the importance of the entrainment rate on the turbulent fluxes at the top of the boundary layer. The representation of turbulent transport by an LES model therefore critically depends on its capability to produce realistic entrainment rates.

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Figure 10 shows the modeled entrainment rates by previous versions of DALES and by other models as reported in intercomparison studies by Duynderkerke et al. (1999), Duynderkerke et al. (2004), and Stevens et al. (2005). Results obtained with the current version of DALES using three different advection schemes are also shown. To facilitate a direct comparison to the previous findings, the new simulations were all run nearly identically to the original case descriptions. The entrainment rates obtained from the LES models represent hourly-average values from the third (ASTEX) or fourth hour of simulations (Eurocs FIRE and DYCOMS RF01).

The entrainment rates from previous versions of DALES were all rather large in comparison to results from other participating LES models used in the intercomparison studies. An initially large entrainment rate generated by the previous version of DALES led to a rapid thinning and subsequent break-up of the DYCOMS RF01 stratocumulus deck. As a consequence the longwave radiative forcing at the boundary-layer top decreased, explaining the small value for the entrainment rate as shown in the figure for “old DALES”. It is clear from the figure that entrainment rates are reduced in the current version of DALES. Entrainment rates in simulations with the monotone kappa advection scheme tend to be slightly larger in comparison to the second or fifth order advection schemes. Only a simulation with the latter scheme is capable to maintain a solid stratocumulus cloud deck for the DYCOMS RF01 case.

### 3.3.2 Shallow cumulus

A number of interesting and well-documented shallow cumulus cases based on observational studies have been simulated by DALES over the last 10 years. These studies include: non-precipitating steady-state marine shallow cumulus based on the Barbados Oceanographic and Meteorological Experiment (BOMEX) (Siebesma et al., 2003) and on the Atlantic Trade Wind Experiment (ATEX) (Stevens et al., 2001), diurnal cycles of shallow cumulus over land observed on June 21, 1997 at the Southern Great Plains (SGP) site (Brown et al., 2002) and during the Small Cumulus Microphysics Study (SCMS) (Neggers et al., 2003a) and more recently precipitating marine shallow

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cumulus (van Zanten et al., 2010) such as observed during the Rain in Cumulus over the Ocean (RICO) field study Rauber et al. (2007). All these cases have been used to critically evaluate the DALES results against observations and to help developing and testing theories, conceptual models and parameterizations of shallow cumulus convection. In this section we will give a short overview of the results of these studies.

The first category of these studies is related to cloud geometrical issues. In Siebesma and Jonker (2000) it has been shown that the simulated cumulus cloud boundaries have self-similar or fractal properties that can be characterized by a fractal dimension  $D_f=7/3$ . These results are in excellent agreement with observational studies and therefore provide a critical test of the capability of DALES to simulate realistic cumulus clouds. Moreover, these results helped in constructing theoretical scaling arguments explaining why cloud boundaries choose to be self-similar with a dimension of  $7/3$ . Another intriguing cloud geometrical topic is related to the question: what is the shape of the cumulus cloud size distribution? It is well known that shallow cumulus cloud ensembles consist of many small clouds and lesser large clouds but the precise shape of the cloud size distribution is still an open issue. Extensive numerical studies with DALES show that the cloud size density of the simulated cloud populations is described well by a power-law from scales smaller than the standard grid-spacing (50 m) up to scales of typically 1000 m with a power-law exponent of  $-1.70$  (Neggers et al., 2003b). This exponent is comparable to values found in observational studies (Cahalan and Joseph, 1989; Rodts et al., 2003). No convincing theory for the powerlaw behaviour nor for the scale break has yet been put forward. Finally, more recently analyses with DALES of up- and downdrafts in and around individual cumulus clouds have shown that strong updrafts in individual cumulus clouds are typically surrounded by so-called subsiding shells with persistent downdrafts (Heus and Jonker, 2008). These downdrafts are driven by negative buoyant forces that result from the evaporative cooling of the cloud water. As they surround the clouds along their entire perimeter, the subsiding shells cover a significant area and are therefore found to be responsible for a large part of the downward mass transport (Jonker et al., 2008).

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The second category studies is related to transport due to cumulus convection which is one of the important processes that needs parameterization in large scale Numerical Weather Prediction (NWP) and climate models. The time evolution of a moist conserved variable  $\varphi$  due to moist convection can be written as

$$5 \quad \frac{\partial \varphi}{\partial t} = - \frac{\partial F_{\varphi}}{\partial z} \quad (110)$$

where  $F_{\varphi}$  is the (upward pointing) turbulent flux. A popular method to parameterize this turbulent flux is through the use of a so called mass flux approach

$$F_{\varphi} \approx \frac{M}{\rho} (\varphi_c - \langle \varphi \rangle). \quad (111)$$

10 where  $\rho$  is the density and the subscript  $c$  refers to cloud averaged values of  $\varphi$  and the mass flux is defined as  $M \equiv \rho a_c w_c$  Betts (1975), i.e. essentially the product of the cloud averaged vertical velocity times  $w_c$  and the fractional cloud area  $a_c$ . Usually a cloud model is derived to obtain equations for  $M$  and  $\varphi_c$

$$\frac{\partial M}{\partial z} = M(\varepsilon - \delta) \quad (112)$$

$$15 \quad \frac{\partial \varphi_c}{\partial z} = -\varepsilon(\varphi_c - \langle \varphi \rangle) \quad (113)$$

20 Within this cloud model the key variables are the fractional entrainment  $\varepsilon$  and fractional detrainment  $\delta$  rate. These inverse length scales are measures of the rate of dilution of the cloud ensemble (entrainment) and the rate of air leaving the cloud ensemble (detrainment) and LES results from DALES have been used extensively to diagnose  $\varepsilon$  and  $\delta$  on the basis of Eqs. (112) and (113) (Siebesma and Cuijpers, 1995). This approach has initiated considerable research in developing theories and models of these exchange mechanisms between clouds and environment. From these studies it has become clear that the fractional entrainment rate can be well estimated by the inverse

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cloud depth (Siebesma et al., 2003). The fractional detrainment rate  $\delta$  is typically larger than  $\varepsilon$  as a result of the fact the cloud fraction  $a_c$  is in general decreasing with height.

Another useful additional equation often used in cloud models is the vertical velocity equation for the cloud ensemble (Simpson and Wiggert, 1969)

$$\frac{1}{2} \frac{\partial w_c^2}{\partial z} = -b \varepsilon_w w_c^2 + aB \quad \text{with} \quad B = \frac{g}{\theta_0} (\theta_{v,c} - \langle \theta_v \rangle) \quad (114)$$

which describes how buoyancy forces and entrainment processes influence the vertical velocity in the clouds. Adjustable prefactors  $a$  and  $b$  are introduced in this equation to incorporate pressure perturbation effects and incloud turbulent effects in an implicit way. By using Eqs. (114) and (112) we can derive alternative expressions for the entrainment that are more linked to the dynamics

$$\varepsilon_w = \left( \frac{B}{w_c^2} \right) - \frac{1}{b} \frac{\partial \ln w_c}{\partial z} + \frac{a}{b} \quad (115)$$

$$\delta_w = \left( \frac{B}{w_c^2} \right) - \frac{(1+b)}{b} \frac{\partial \ln w_c}{\partial z} + \frac{a}{b} - \frac{\partial \ln a_c}{\partial z} \quad (116)$$

In Fig. 11 we compare the entrainment and detrainment rates based on Eqs. (114), (115), and (116) for which estimates of  $a=0.6$  and  $b=1$  are used for a large variety of different LES experiments.

The fact that the results fall reasonably well on the diagonal shows that Eqs. (114) and (112) are consistent, so that the subscripts of  $\varepsilon_w$  and  $\delta_w$  can be removed and Eqs. (115) and (116) can be used as well to interpret the exchange rates. It can also be observed that  $\varepsilon$  can vary considerably between values of  $1 \sim 4 \times 10^{-3} \text{ m}^{-3}$  indicating that parameterizations that use a constant value for  $\varepsilon$  is not a good option. Furthermore it should be noted that the range of variability for  $\delta$  is much larger  $1 \sim 20 \times 10^{-3} \text{ m}^{-3}$ . More detailed analysis shows that this large variability is mainly due to the gradient

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of the cloud fraction with height in Eq. (116). This indicates that, in order to have a good estimate of the mass flux  $M$ , it is more relevant to have a good parameterisation of  $\delta$  rather than for  $\varepsilon$ , a statement already emphasized in de Rooy and Siebesma (2008). In that respect it is surprising to see that most of the research efforts have been concentrated on entrainment rather than on detrainment.

### 3.4 Heterogeneous surfaces

DALES has contributed to the understanding of flow over thermally heterogeneous terrain. The study of van Heerwaarden and Vilà-Guerau de Arellano (2008) addressed the question whether convective cloud formation is more likely to form over a land surface that has a heterogeneous surface flux compared to a land surface that is homogeneously heated.

Heterogeneous land surfaces were simulated by creating two stripes of 3.2 km wide at the land surface in the model, as this is the spatial scale at which heterogeneity is considered to modify the turbulent structure of the overlying CBL the most. The turbulent fluxes at the land surface were prescribed. Both stripes had the same sum of sensible and latent heat, but a different Bowen ratio. The left stripe was characterized by a small Bowen ratio, whereas the right stripe had a large ratio. In this setup the LES model was run for four hours; statistics were calculated over the last hour.

The main findings of the study are summarized in Fig. 12 that shows the relative humidity in the CBL and the wind vectors in a case where the free atmosphere is moist (left panel) and in a case where the free atmosphere is dry.

In both cases a secondary circulation (see wind vectors) distributes heat and moisture towards the area that has a relatively large sensible and a small latent heat flux. At these hot spots, strong but moist thermals rise, resulting in a large relative humidity over the area that has the smallest latent heat flux. In case of a dry free troposphere (right panel), the secondary circulation can transport very dry free tropospheric air downwards to the land surface. Therefore, a very low relative humidity is found over the area that has the largest latent heat flux.

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To conclude, the study showed that heterogeneity results in a situation that is more favorable for cloud formation, regardless of the specific humidity of the free troposphere. Using a similar set up, Górska et al. (2008) use DALES to determine the role of heterogeneity on the carbon dioxide distribution. The study points out of the need to redefine aircraft measurements strategies above non-uniform surfaces.

### 3.5 Atmospheric flow over sloping surfaces

Compared with the many successful large-eddy simulations of the boundary layer over flat terrain, as of yet only a few simulations of the ABL over sloping surfaces have been carried out. One of the problems concerning the simulation of slope flow, is that the potential temperature as well as the depth of the flow and the flow velocity change along the slope. Observations of katabatic flow, however, have shown that the flow in a strongly stratified boundary layer and/or over (moderately) steep slopes varies only slightly along the slope (Haiden and Whiteman, 2005). Therefore, although DALES currently only facilitates periodic boundary conditions, we are still able to study homogeneous slope flow. One of the outlooks is to implement open boundary conditions, which would enable the simulation of slope flow under a larger range of circumstances. Nonetheless, DALES has in recent years successfully been used to study homogeneous katabatic flow over moderately steep slopes.

Axelsen and van Dop (2010) performed a model validation by comparing simulation results to observations from two glaciers. They found that the simulated profiles of temperature and downslope velocity were quantitatively in agreement with the observations. An example is given in Fig. 13. Near the surface the downslope velocity increases with height and reaches a maximum at a height of  $z \approx 4$  m. Above the wind maximum height, the downslope velocity decreases with height. The figure shows that near the surface the simulated and observed velocity profiles agree, but above the wind maximum the model underestimates the velocity. The profile of the simulated potential temperature is also seen to agree rather well with the mast measurements, but that there is a systematic offset between the balloon measurements and the simulated

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potential temperature. The agreement between modelled and observed fluxes of momentum and buoyancy were less satisfactory (not shown). However, the observations were influenced by processes such as a valley wind, cross-slope winds, and gravity waves. The latter processes, which enhance the turbulence production and the vertical mixing in the katabatic layer, are not accounted for in the numerical model.

### 3.6 Dispersion and chemically reacting flows

We summarize here the main research results achieved in the field of turbulent dispersion and chemical transformations using DALES. The plume dispersion main characteristics and statistics under different ABL flow conditions have been thoroughly investigated using DALES. Dosio et al. (2003, 2005) and Dosio and Vilà-Guerau de Arellano (2006) investigated the plume dispersion in the dry CBL from a Eulerian and a Lagrangian perspective. Based on DALES results, they derived a parameterization to include the effect of shear on the plume spreading, studied the validity of Taylor's diffusion theory for horizontal and vertical dispersion, and separated the contributions of small- and large-scales on the plume evolution, both from an absolute coordinate system as well as relative to the plume's center of mass. Verzijlbergh et al. (2009) extended this study to determine the influence of stratocumulus and shallow cumulus on the turbulent dispersion properties and related to turbulent structures like skewness of the vertical velocity. As an example, Fig. 14 shows the vertical concentration characteristics and the location of the maximum concentration under different ABL conditions: dry convective boundary layer, stratocumulus and shallow cumulus Verzijlbergh et al. (2009).

Similarly to turbulent dispersion, the chemical transformations in the ABL are influenced by the characteristics of the turbulent flow. This turbulent control is particularly important when the turbulent time scale ( $\tau_t$ ) and the chemistry time scale ( $\tau_c$ ) have similar values, i.e., the order of magnitude of the Damköhler number ( $\tau_t/\tau_c$ ) is  $\mathcal{O}(1)$ . Under this regime, the species are chemically transformed at a different reaction rate

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depending on the way species are introduced in the ABL, premixed or non-premixed, and the turbulent intensity to mix chemical species. Key tropospheric chemical reactions involving species such as nitric oxide and certain biogenic hydrocarbons like isoprene are therefore controlled by turbulence.

5 Following Schumann (1989), Petersen et al. (1999) and Petersen and Holtslag (1999) studied by means of LES how the transport and mixing of reactants in the CBL is influenced by the presence of vigorous thermals and subsidence motions. Based on the DALES results, they suggested a parameterization to represent the fluxes and covariance of reactants in large scale chemistry transport models. The research was  
10 extended to study more complex mechanism under non-uniform emissions of the reactants (Krol et al., 2000; Vilà-Guerau de Arellano et al., 2004). To further study the influence of the reactivity on high-order moments, a spectral analysis showed that the reactant variability (variance) depends strongly on the reaction rate (Jonker et al., 2004). The analysis was done using the DALES simulation of a turbulent flow reacting  
15 according to the scheme (R1)–(R2). These results showed large variations in the characteristic length scale as a function of the Damköhler number and the state of the chemical equilibrium.

To improve parameterizations in large-scale atmospheric chemistry models, Vinuesa and Vilà-Guerau de Arellano (2003, 2005) proposed an expression of an effective reaction rate ( $k_{\text{eff}}$ ) that takes into account explicitly the influence of turbulent mixing on  
20 the reaction rate.

The moist and optically thick boundary layer clouds can also influence atmospheric chemistry. DALES was used to study the combined effect of turbulence and radiation on simple chemical mechanism in a dry smoke cloud (Vilà-Guerau de Arellano and Cuijpers, 2000) and shallow cumulus (Vilà-Guerau de Arellano et al., 2005). Figure 15  
25 shows the cloud water content and the photostationary state ( $\Phi$ ) in a CBL developed over land characterized by the presence of shallow cumulus.  $\Phi$  quantifies the effect of the physical processes (turbulence and radiation) on the atmospheric chemistry. For the reactants nitric oxide (NO), ozone ( $\text{O}_3$ ) and nitrogen dioxide ( $\text{NO}_2$ ), it is defined as

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$\Phi = (k[\text{NO}][\text{O}_3]) / (j[\text{NO}_2])$ . Departure from the value  $\Phi = 1$  indicate perturbations of the chemical equilibrium either by radiation or turbulent processes.

#### 4 Outlook

As was shown in this paper, DALES can provide reliable results for a multitude of atmospheric conditions, and there are many alleys of study that can be pursued with DALES 3.2. In the field of cloudy boundary layers, very fine grid spacing can be used to reliably resolve most of the dynamics within and around the cloud. Simulations on relatively large horizontal domains ( $\sim 25$  km) can mimic the physics in an area similar to a single column of a regional or global model. On that scale, LES is well capable of variability studies that are necessary to improve the GCMs, and to study the impact of GCM grid refinement. For other studies, LES can provide spatial and temporal turbulence characteristics that cannot be easily retrieved from measurements alone. This is always a role that LES can play, but it can be especially important in spatially anisotropic or inhomogeneous situations, such as in the fields of flow over sloped or heterogeneous surfaces.

While there are many plans to use DALES in its current state, ongoing improvement of the code is also planned. In the near future, we aim to be able to run DALES in more diverse and more realistic scenarios than what was shown in this paper. Furthermore, we aim to focus on studies that makes integrated use of several of the features of DALES.

As was shown throughout many parts of the applications section, LES could still benefit from a better representation of anisotropic turbulence around steep gradients and inversion layers, specifically in stable boundary layers, dry convective boundary layers, and stratocumulus layers. Increasing computer power and resolution could resolve these gradients within the coming years, but more intelligent subfilter-scale modeling could also give a significant contribution in solving this problem. This is especially important in critical stratocumulus cases, where entrainment of relatively dry and warm air leads to buoyancy reversal.

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To study the interactions between the various components of the model, we strive to have the modules as interactive as possible. This could for instance lead to better understanding of coupling mechanisms between radiative forcings and the surface conditions, coupling between radiation and chemistry, or between chemistry and cloud and aerosol formation.

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**Table 1.** The main thermodynamical constants used throughout this paper.

$R_v$	Gas constant for water vapor	$461.5 \text{ J kg}^{-1} \text{ K}^{-1}$
$R_d$	Gas constant for dry air	$287.0 \text{ J kg}^{-1} \text{ K}^{-1}$
$L$	Latent heat release for vaporization	$2.5 \times 10^6 \text{ J kg}^{-1}$
$c_{pd}$	Heat capacity for dry air	$1004 \text{ J kg}^{-1} \text{ K}^{-1}$

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**Table 2.** An overview of the parameters used in the SFS scheme of DALES model. Not all parameters are independent.

$\alpha$	$c_f$	$c_{\varepsilon,1}$	$c_{\varepsilon,2}$	$c_m$	$c_{h,1}$	$c_{h,2}$	$c_N$
1.5	2.5	0.19	0.51	0.12	1	2	0.76

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**Table 3.** Simulation details of the two simulated CBLs: weak inversion case (W06) and strong inversion case (S24).  $z_i(0)$  and  $\Delta\theta(0)$  denote the initial mixed-layer depth and initial temperature jump, respectively.

Case	$Q$ $\text{K m s}^{-1}$	$d\langle\theta_v\rangle/dz$ $\text{K m}^{-1}$	$z_i(0)$ m	$\Delta\theta_v(0)$ K	$z_i$ m	$w_*$ $\text{m s}^{-1}$
W06	0.06	0.003	750	0	1230	1.34
S024	0.24	0.003	950	8	1096	2.05

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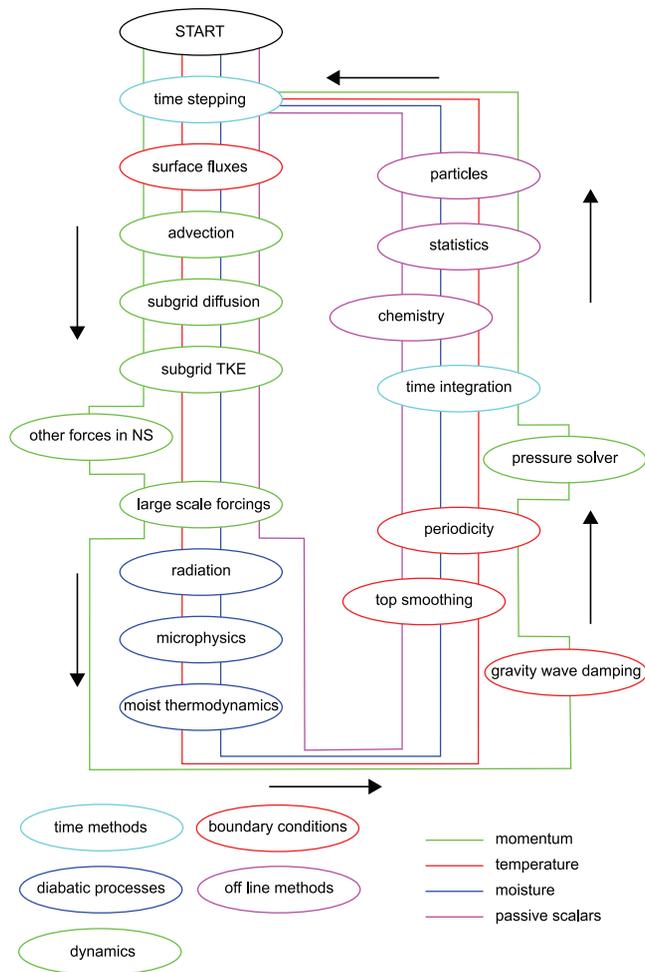
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**Fig. 1.** Flowchart of DALES.

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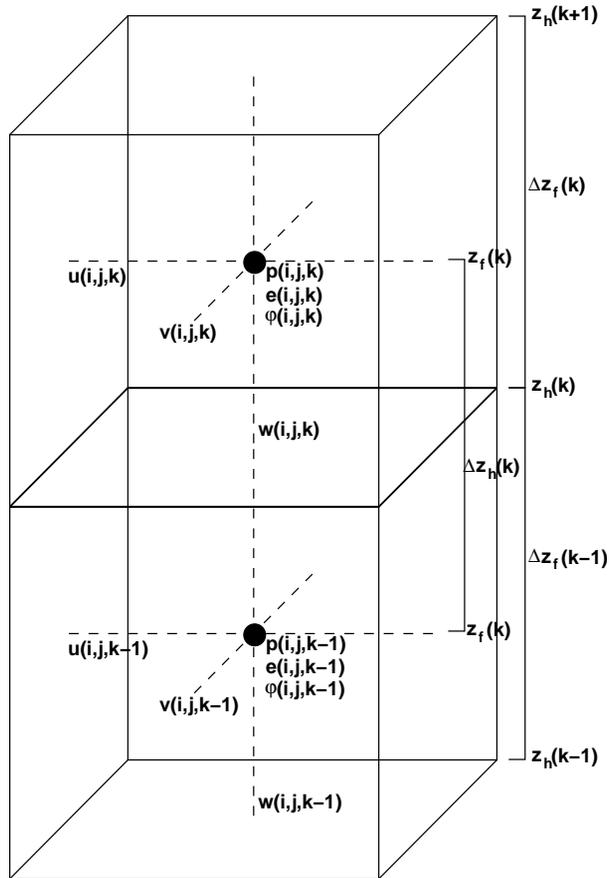
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**Fig. 2.** The Arakawa C-grid as used in DALES. Pressure, SFS-TKE and the scalars are defined at cell-center, the 3 velocity components at the face of the cell. The level of cell center is called the full level (denoted with an “f”); the level where  $w$  is located is called the half level (an “h”). The (variable) vertical grid spacing  $\Delta z$  is defined centered around the belonging level.

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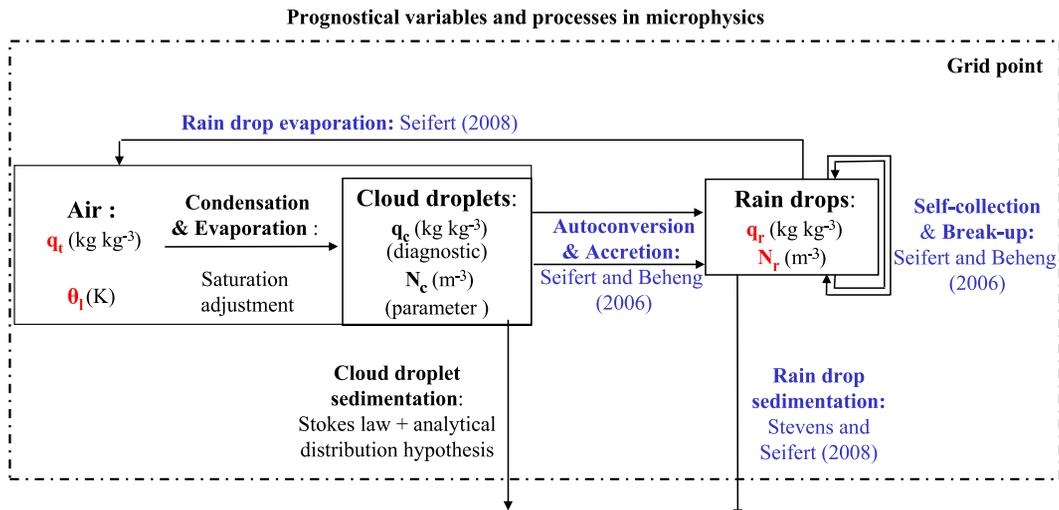
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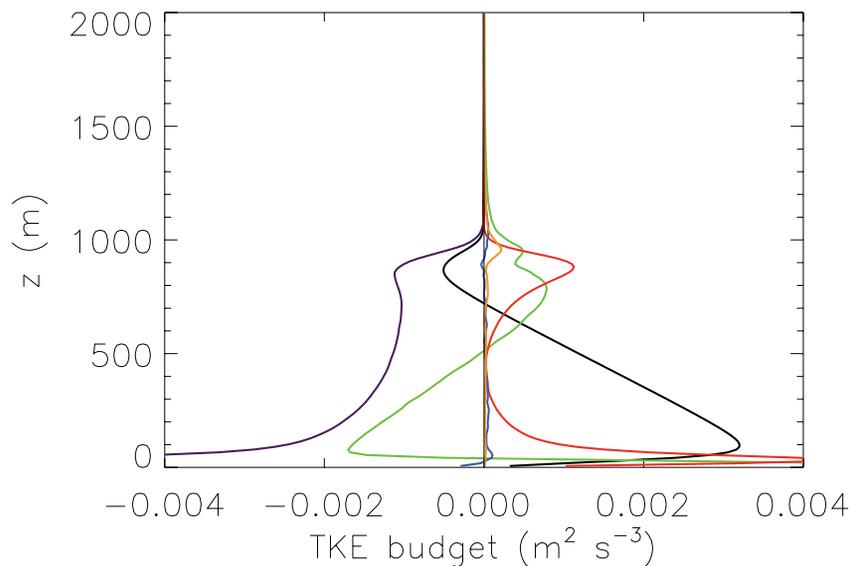




**Fig. 3.** Representation of the prognostic thermodynamical variables  $\tilde{\theta}_l$ ,  $\tilde{q}_t$ , the microphysical parameter and variables  $N_c$ ,  $q_c$ ,  $\tilde{N}_r$ ,  $\tilde{q}_r$ , and the microphysical processes relating these variables.

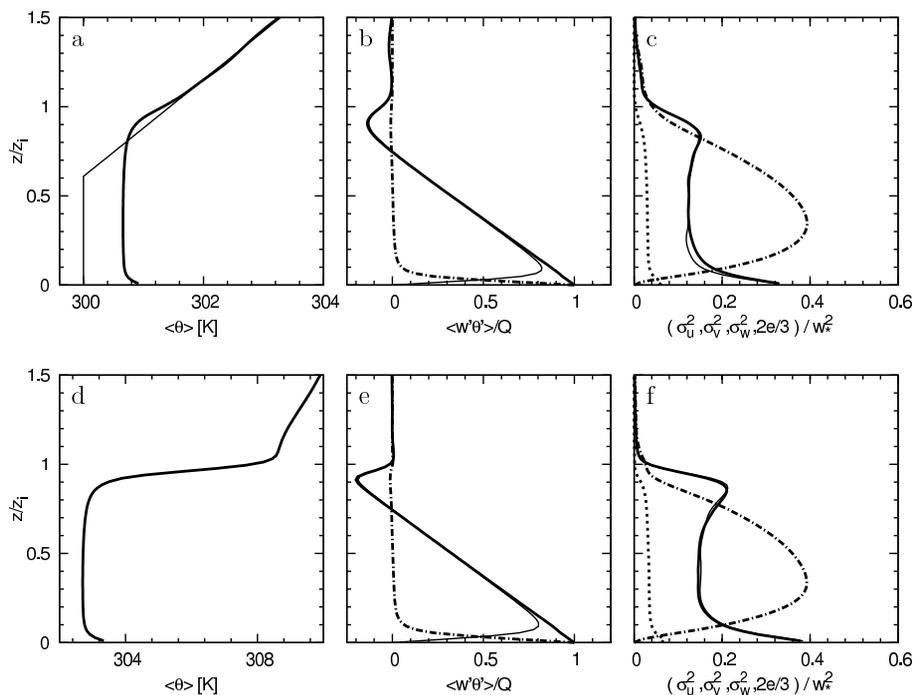
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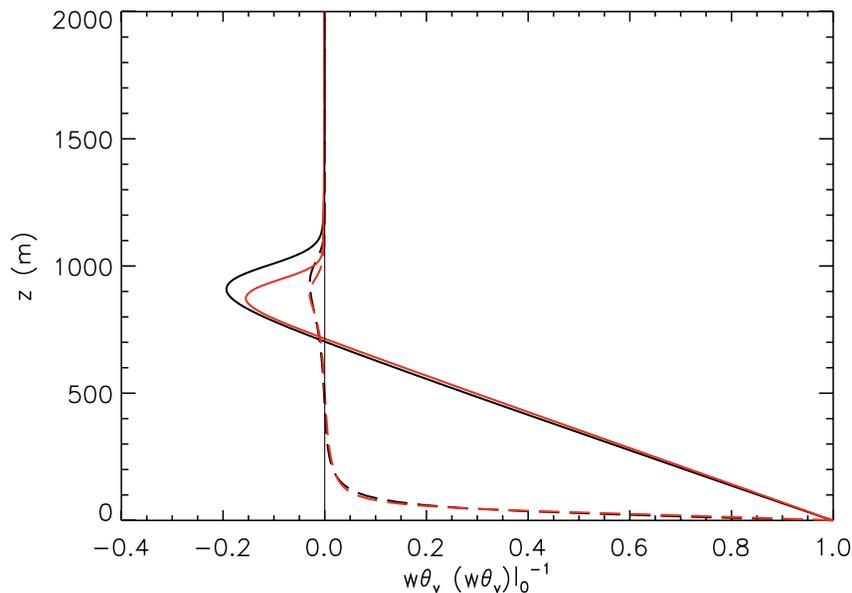
**Fig. 4.** Vertical profile of the various terms of the TKE budget in a sheared CBL: total tendency (orange), buoyancy (black), dissipation (violet), transport and pressure (green), shear (red), and the residual (blue).

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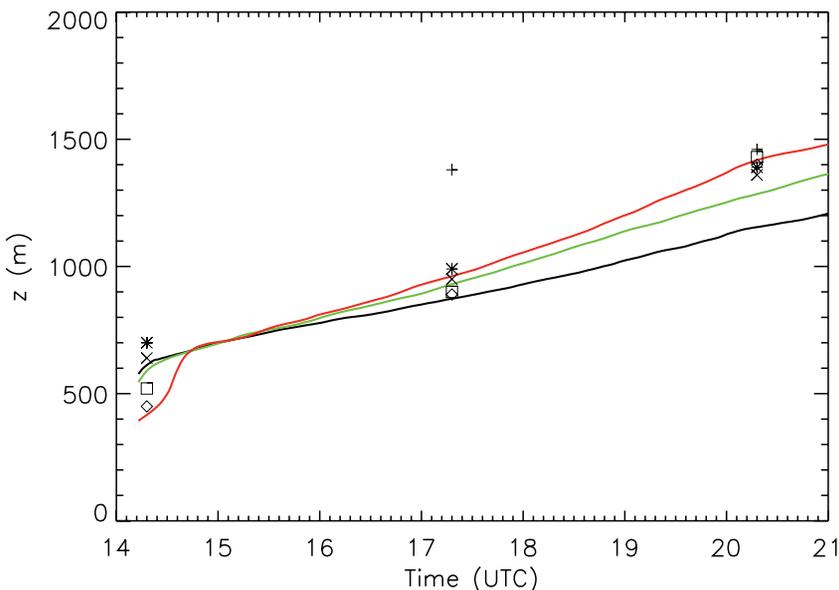
**Fig. 5.** DALES results for a CBL with a weak inversion (a,b,c) and with a strong inversion (d,e,f), reproducing cases W06 and S24 of the study by Sullivan et al. (1998). For extra information see Table 3. All results are averages over hour 3-4. a,d: average temperature profile (thin line in a shows the initial temperature profile). b,e: normalized heat-flux profiles, resolved (thin line), subgrid (dashed) and total (solid line). c,f: turbulence statistics,  $\sigma_u^2 = \overline{u'^2}$  (solid line),  $\sigma_v^2$  (thin line),  $\sigma_w^2$  (dashed line), and subgrid contribution (dotted line).

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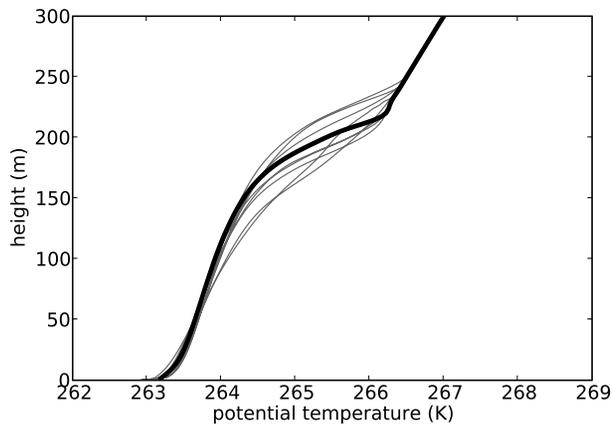
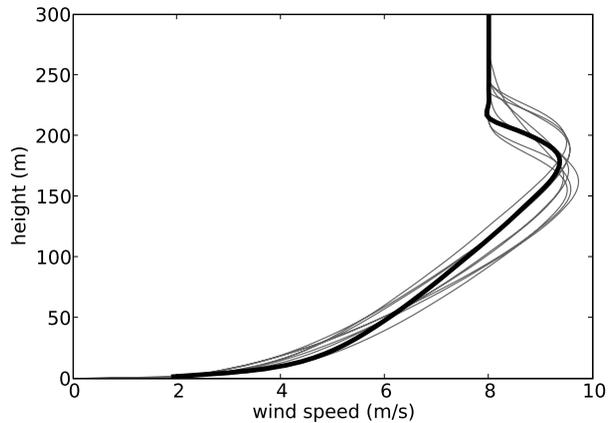
**Fig. 6.** Vertical profile of the total  $\langle \widetilde{w}\widetilde{\theta}_v + \widetilde{w}'\theta_v' \rangle$  (solid line) and subfilter-scale contribution  $\langle \widetilde{w}'\theta_v' \rangle$  (dashed) of the virtual potential temperature flux obtained after four hours of simulation by DALES2.0 (black) and DALES 3.2 (red) with the same physical conditions and advection scheme (2nd order central differences).

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**Fig. 7.** Boundary layer height  $z_i$  observed by radiosondes launched at different facilities of ARM campaign (symbols) and obtained by means of LES: without shear (black), including a constant geostrophic wind of  $10 \text{ ms}^{-1}$  in the east-west direction (green), and prescribing the observed mean wind (red). Adapted from Pino et al. (2003).

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**Fig. 8.** Profiles of mean wind speed (top) and potential temperature (bottom) for the first GABLS1 LES intercomparison (average over 9th hour of simulation). Solid black line: DALES result at 3.125 m resolution and  $c_f=2.0$ ; grey lines: results of other participants at 3.125 resolution.

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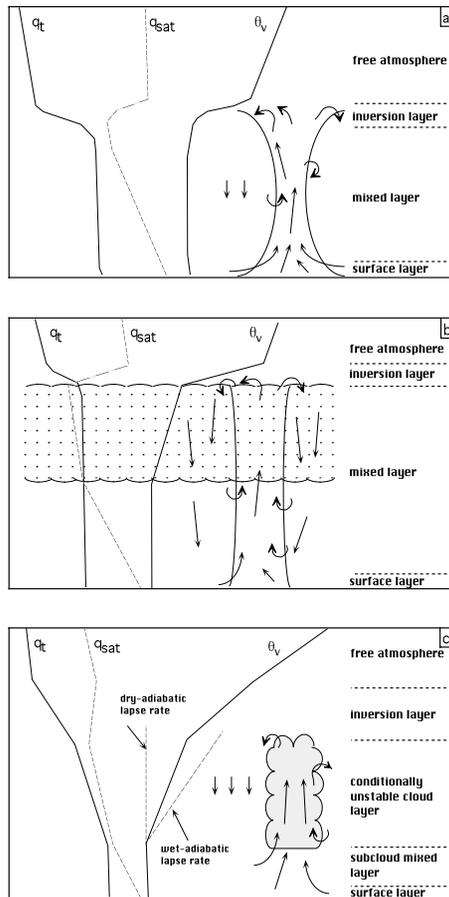
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**Fig. 9.** Schematic overview of the different types of boundary layer clouds.

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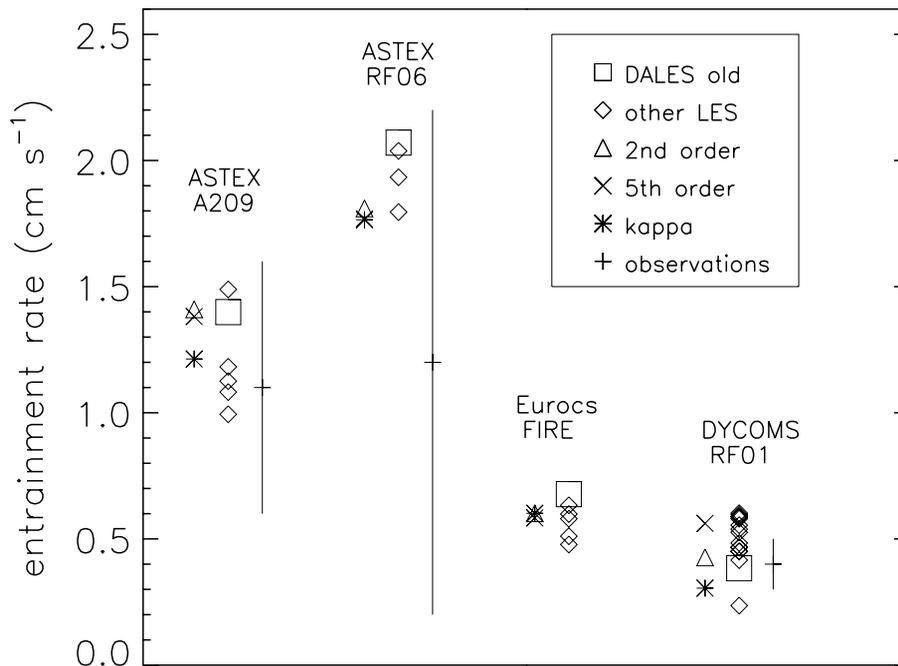
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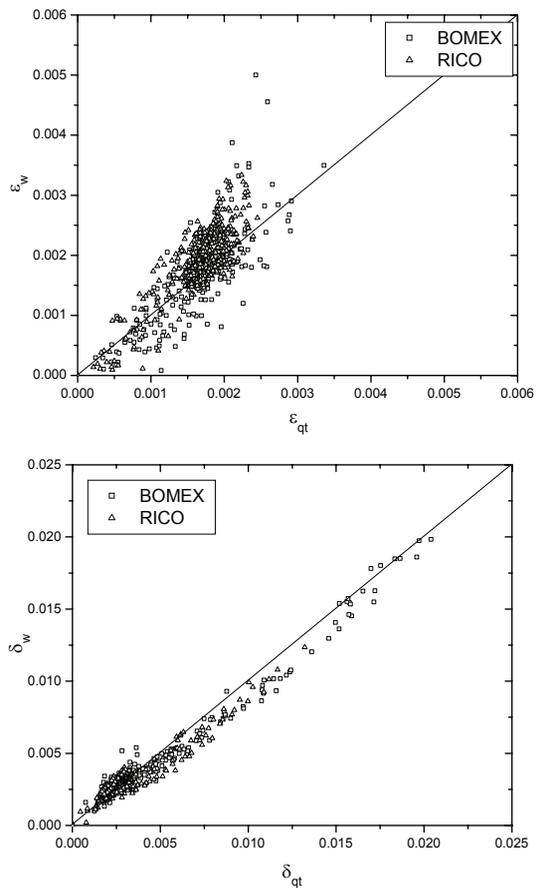
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**Fig. 10.** Hourly-averaged entrainment rates from LES models for four different GCSS cases. The “DALES old” and “other LES” indicate entrainment results as obtained from the original modeling results. The observed entrainment rates with their uncertainties are also plotted. Because the Eurocs FIRE case is based on a monthly mean climatology, no observed entrainment rates are available for this case. The “DALES old” results were all obtained with the kappa scheme. The results with the current DALES version were obtained with the kappa, second-order and fifth-order advection schemes as indicated in the legend.

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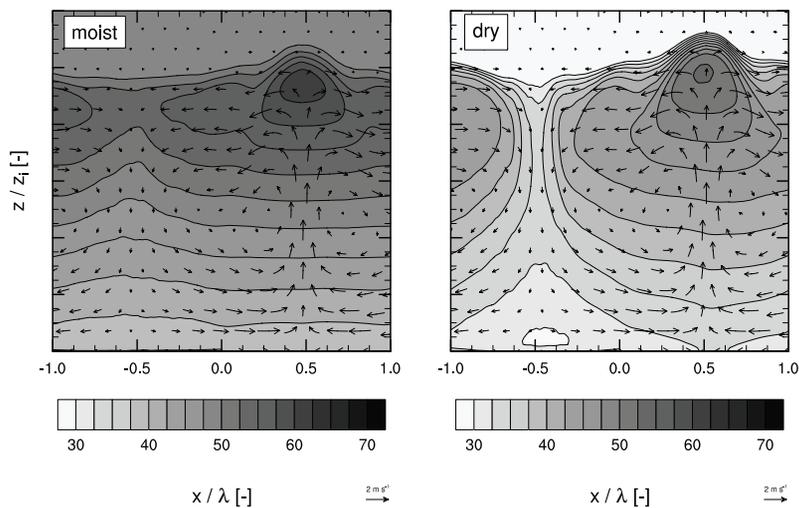



**Fig. 11.** Comparison of LES derived fractional entrainment and detrainment rates  $\varepsilon_q$  and  $\delta_q$  using  $\phi=q_t$  based on Eq. (114) (horizontal axis) versus LES estimates of these rates  $\varepsilon_w$  and  $\delta_w$  based on the vertical velocity equation (Eqs. 115 and 116).

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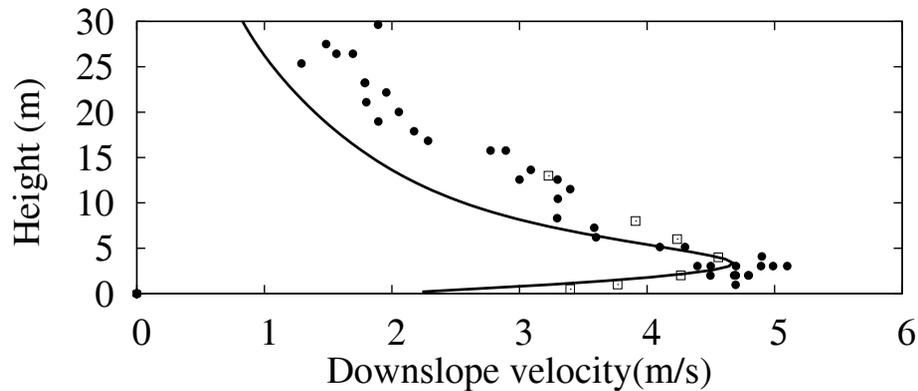
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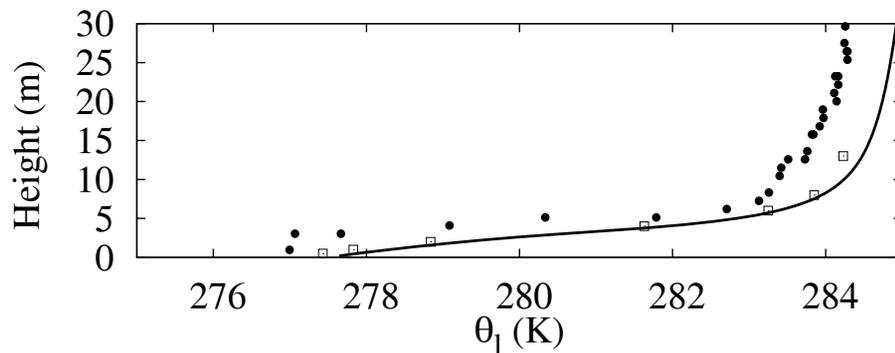
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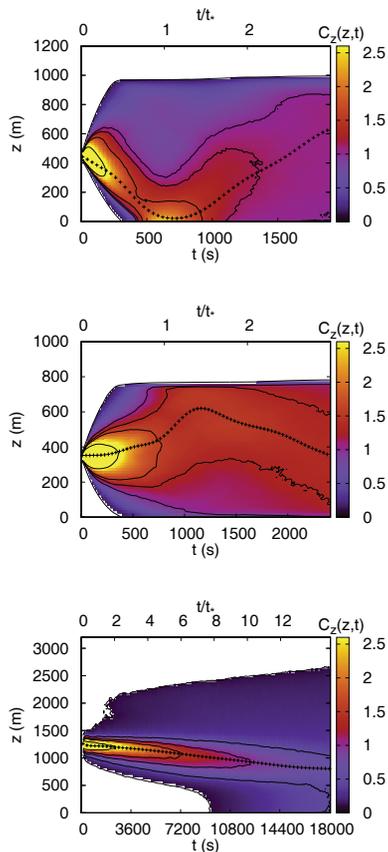
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(a)



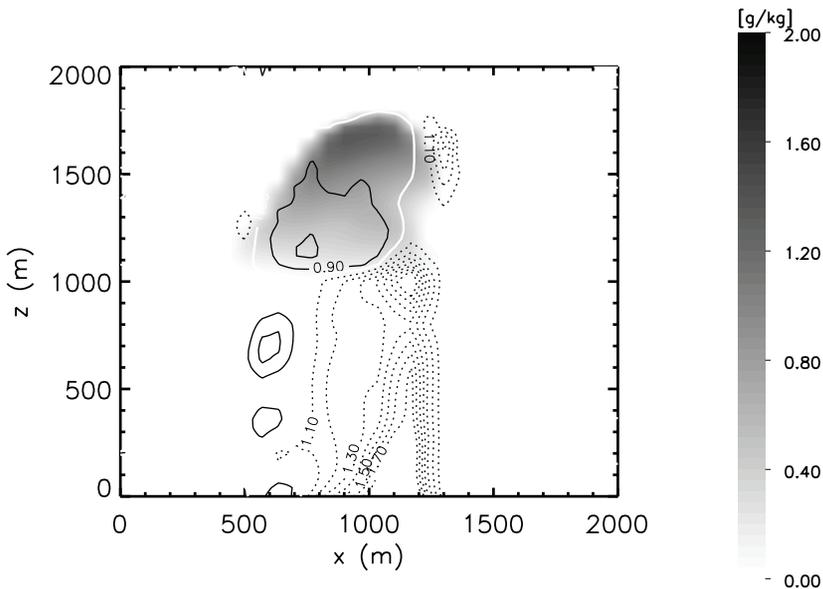
(b)

**Fig. 13.** Mast profiles (squares), balloon data (dots) and LES profiles of downslope velocity **(a)** and potential temperature **(b)** in flow over a sloped surface.



**Fig. 14.** Evolution on time of the vertical concentration (crosswind integrated) of a plume released as a function of the releasing time (dimensional and non-dimensional) for (top) dry convective conditions, (middle) stratocumulus topped boundary layer, and (bottom) shallow cumulus topped boundary layer. Concentration has been multiplied by a factor 1000 to obtain a convenient scale. The crosses indicates the position of the maximum concentration.

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**Fig. 15.** Instantaneous vertical cross section of the cloud water  $q_c$  (g/kg) content and the photostationary state ( $\Phi$ ) calculated using the NO, NO<sub>2</sub> and O<sub>3</sub> mixing ratios. At chemical equilibrium  $\Phi=1$ .

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