



**A spreadsheet-based
snow model for
inside-canopy
conditions**

T. Marke et al.

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ESCIMO.spread (v2): parameterization of a spreadsheet-based energy balance snow model for inside-canopy conditions

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Title Page	
Abstract	Introduction
Conclusions	References
Tables	Figures
◀	▶
◀	▶
Back	Close
Full Screen / Esc	
Printer-friendly Version	
Interactive Discussion	



A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



for precipitation phase detection that applies wet-bulb temperature as a criterion to distinguish solid and liquid precipitation. Another model improvement is a new parameterization for cold and liquid water content of the snow cover allowing to consider refreezing of rain or meltwater in the snowpack. Compared to other existing spreadsheet-based snow models (e.g. the glacier and snow melt study model by Brock and Arnold, 2000) ESCIMO.spread (v2) is particularly fast and can easily be modified by simple change of the parameters and formulae with results immediately visualized. The model requires only few input data (hourly recordings of temperature, precipitation, wind speed, relative humidity, global as well as longwave radiation) and is even capable of simulating the evolution of a seasonal snow cover under climate change conditions. The latter is realized by defining trends in precipitation and/or temperature in the models parameter section. While Walter et al. (2005) have presented a spreadsheet energy balance model that requires even less meteorological input data (daily minimum/maximum temperature and precipitation), their approach operates at a daily time step only and does not allow to quantify sub-daily variations in snow cover conditions. Moreover, compared to the canopy model implemented in ESCIMO.spread (v2), the consideration of canopy effects in the Walter et al. (2005) model is reduced to a canopy-induced extinction of solar radiation only. Canopy effects on other meteorological variables or vegetation-snow cover interactions (e.g. the interception of snow in the canopy) are not accounted for. As ESCIMO.spread (v2) is in simple table format and does not include any macros it can be applied by all common spreadsheet programs (e.g., Microsoft Excel, Apple Numbers, OpenOffice Calc) on a variety of platforms (Windows, Linux, Mac OS). Due to its simplicity, ESCIMO.spread (v2) is particularly suitable for application in education (e.g., in practically-oriented student courses) and can even be operated with laptop computers, e.g. to visualize and plausify measured meteorological parameters and the simulated snow cover directly in the field.

With its new features of

- sophisticated precipitation phase detection using a wet bulb temperature threshold,

A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



- cold content and liquid water content calculation with consideration of refreezing of water from melt or rain, and meltwater outflow,
- transformation of standard meteorological observations (precipitation, relative humidity, temperature, wind speed, global radiation) from the open into conditions inside a forest canopy,
- calculation of snow interception and subsequent sublimation, melt or dropping of intercepted snow to the ground and
- calculation of the beneath-canopy snow energy and mass balance,

the new version ESCIMO.spread (v2) reaches beyond the capabilities of most other freely available point-scale snow models and can be expected to set forth the history of ESCIMO.spread as a well-accepted, documented and freely available snow model for application in both science and education. This paper describes the newly implemented algorithms and evaluates the model results against available hydrometeorological observations in and outside the forest canopy at a site in the Black Forest mountain range (south-west Germany, see Fig. 1) with a mostly temperate snow cover in an elevation of 800 m a.s.l. The applied hydrometeorological data have been recorded by a set of low-cost snow monitoring systems (SnoMoS) recently developed by Pohl et al. (2014). The model can be downloaded from www.alpinehydroclimatology.net together with one year of example meteorological recordings and snow observations.

2 The ESCIMO.spread model (v2)

2.1 General description

The new version ESCIMO.spread (v2) builds upon the ESCIMO.spread model as published by Strasser and Marke (2010). It is a 1-D, one-layer process model which calculates snow accumulation and melt for a snow cover assumed to be a single and

A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



homogeneous pack. To do so, it solves the energy and mass balance equations for the snow surface applying simple parameterizations of the relevant processes. The energy balance of the snow surface is calculated for each hourly time step considering short- and longwave radiation, sensible and latent heat fluxes, energy conducted by solid or liquid precipitation as well as sublimation/resublimation and a constant soil heat flux (Strasser and Marke, 2010). Thereby, absorbed and reflected shortwave radiation is calculated from incoming shortwave radiation on the basis of the snow albedo, which is estimated for each hourly time step using an albedo ageing curve approach. Solid precipitation increases the amount of snow water equivalent (SWE) on the land surface, while liquid precipitation is (up to a certain maximum amount depending on actual SWE) added to the liquid water storage of the snowpack. Melt is calculated from the energy balance remainder if air temperature exceeds 273.16 K, whereas no melt is assumed to occur for air temperatures below 273.16 K. The model results are visualized in the form of diagrams for the majority of model variables, together with three quantitative measures of goodness of fit.

2.2 Precipitation phase detection

The new version of ESCIMO.spread (v2) includes an improved distinction between liquid and solid precipitation. As air temperature T_a is often an insufficient indicator for the precipitation water phase (Steinacker, 1983), wet-bulb temperature T_w is used in ESCIMO.spread (v2) as a combined measure of air temperature and humidity to distinguish rain from snowfall. Figure 2 shows the relation between air temperature, wet-bulb temperature and relative humidity for different altitudes to account for the dependence of wet-bulb pressure on air pressure. Each of the displayed lines in Fig. 2 could be interpreted as a borderline to separate liquid and solid precipitation assuming a certain threshold wet-bulb temperature. Largest differences between air temperature and wet-bulb temperature occur at low air humidities, clearly pronouncing the added value associated to application of wet-bulb temperature as a criterion for phase detection.

Generally, wet-bulb temperature can be derived by solving the psychrometric equation

$$e_a(T_a) - e_s(T_w) - A \times (T_a - T_w) = 0 \quad (1)$$

for T_w (K), where A (PaK^{-1}) is the psychrometric constant, and $e_a(T_w)$ (Pa) and $e_s(T_w)$ (Pa) the vapor pressure of the air and the saturation vapor pressure at wet-bulb temperature, respectively. As there is no explicit solution to the psychrometric equation (Campbell and Norman, 1998) and iterations are unfavourable in a spreadsheet model, a pragmatic assumption has been made: for a broad range of combinations of air temperature and relative humidity values, lookup tables have been generated outside the spreadsheet model using an iterative solution scheme for Eq. (1). Beside temperature and humidity, wet-bulb temperature also depends on air pressure p_z (Pa) which is required to calculate the psychrometric constant, A , as (Kraus, 2004)

$$A = \frac{\rho_z \times c_p}{0.622 \times L_v}, \quad (2)$$

where c_p is the specific heat capacity of air at constant pressure ($1004 \text{ J kg}^{-1} \text{ K}^{-1}$) and L_v (J kg^{-1}) represents the latent heat of vaporization. In ESCIMO.spread (v2) the temperature dependence of the psychrometric constant is neglected since this dependency is by far less important compared to that associated to air pressure at higher altitudes (Kraus, 2004; Campbell and Norman, 1998). Air pressure, p (hPa), at a given elevation, z (m), can be derived from standard atmospheric pressure, p_0 ($= 1013 \text{ hPa}$), by integration of the hydrostatic equation assuming a linear decrease of temperature with increasing altitude ($\gamma = -0.0065 \text{ K m}^{-1}$)

$$p_z = p_0 \left[\frac{T_a}{T_a - \gamma \times z} \right]^{-\frac{g}{\gamma \times R}}, \quad (3)$$

where R is the gas constant of dry air ($287 \text{ J kg}^{-1} \text{ K}^{-1}$) and g is gravity (ms^{-2}). To account for the air pressure dependence, the implemented lookup tables have been

GMDD

8, 8155–8191, 2015

A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



prepared for several elevation bands with a 500 m interval. Figure 3 shows a comparison of wet-bulb temperatures calculated using the lookup table approach to those achieved with an iterative solution for different elevations. The differences between both approaches shown for a common snowfall situation are relatively small. Therefore, the lookup table approach allows a sufficiently accurate estimation of wet-bulb temperature in the model. The threshold for wet-bulb temperature as required for precipitation phase detection in ESCIMO.spread (v2) is one of the user-defined input parameters and is here set to 273.16 K.

2.3 Cold and liquid water content

A conceptual parameterization for the liquid water and cold content of the snowpack based on Braun (1984) and Blöschl and Kirnbauer (1991) has been added to the model. Heat losses resulting from a negative energy balance are thereby used to build up a cold content, which first needs to be reduced to zero by positive energy inputs before actual melt can occur. Conversely, melting snow is not immediately removed from the snowpack, but a certain amount of liquid water can be retained (and possibly refreeze again). This approach accounts for the delay between beginning surface melt and drainage of a snow cover.

Three model parameters control these processes: the water holding capacity, HC_w (-), the cold holding capacity, HC_c (-) (both specified as a fraction of the total snowpack weight), and the refreezing factor, F_r (-), which is the fraction of the computed heat loss used for refreezing and building up the cold content. By default, these parameters are set as recommended by Blöschl and Kirnbauer (1991) to $HC_w = 0.1$, $HC_c = 0.03$, and $F_r = 0.5$.

The potential melt, M_p (mm) is calculated using the available melt energy, M (J), as

$$M_p = \frac{M}{c_i}, \quad (4)$$

where c_i is the melting heat of ice ($3337 \times 10^5 \text{ J kg}^{-1}$).

A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



In the case of a negative energy balance ($M < 0$), this “negative melt” is first used to calculate a refreezing of liquid water in the snowpack, RF (mm), in the form of

$$RF = \min \{C_{lw,t-1}, -M_p \times F_r\}, \quad (5)$$

where $C_{lw,t-1}$ (mm) is the liquid water content of the previous time step. C_{lw} for a given time step is then derived as

$$C_{lw} = C_{lw,t-1} + P_r - RF, \quad (6)$$

where P_r (mm) is liquid precipitation.

The remaining amount of energy is used to increase the cold content C_c (the cold content is represented with negative values, while the liquid water content is represented with positive values) as

$$C_c = \max \{C_{ct-1} + (M_p \times F_r + RF), -HC_c \times SWE_{t-1}\}. \quad (7)$$

In the case of a positive energy balance ($M > 0$), first the sum of cold content reduction and actual melt is calculated as

$$M_v = \min \{SWE_{t-1} - C_{ct-1}, M_p\}. \quad (8)$$

C_c is then reduced with

$$C_c = \min \{C_{ct-1} + M_v, 0\} \quad (9)$$

and actual melt, M_a (mm), is calculated as

$$M_a = M_v - (C_c - C_{ct-1}). \quad (10)$$

C_{lw} is then updated in the form of

$$C_{lw} = \min \{C_{lw,t-1} + M_a, SWE_{t-1} \times HC_w\} \quad (11)$$

and the outflow (i.e. the excess water that is actually removed from the snowpack), O (mm), is finally calculated as

$$O = \max \{(C_{lw,t-1} + P_r + M_a) - SWE_{t-1} \times HC_w, 0\}. \quad (12)$$

A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

Title Page	
Abstract	Introduction
Conclusions	References
Tables	Figures
◀	▶
◀	▶
Back	Close
Full Screen / Esc	
Printer-friendly Version	
Interactive Discussion	



2.4 Modification of meteorological conditions inside the forest canopy

The canopy model newly implemented in ESCIMO.spread (v2) by Liston and Elder (2006b) has already been successfully applied under alpine conditions (see Strasser, 2008 or Strasser et al., 2011). The development of the approach was motivated by the fact that meteorological observations inside forest canopies only sparsely exist necessitating the estimation of inside-canopy conditions from available meteorological observations in the open. The method requires information on leaf area index and canopy height which can either be derived from field measurements or be taken from literature for a wide range of plant species (e.g. from Breuer et al., 2003, or Liston and Elder, 2006b).

Wind speed inside the canopy u_c (m s^{-1}) is derived from above-canopy wind speed u (m s^{-1}) as (Cionco, 1978)

$$u_c = u \exp(-a(1 - z/h)), \quad (13)$$

where h (m) is the canopy height and z (m) is the canopy reference level assumed to be $0.6h$ (Liston and Elder, 2006b; Essery et al., 2003).

The canopy flow index, a (–) is calculated as a function of the effective leaf area index, LAI^* ($\text{m}^2 \text{m}^{-2}$), and a scaling factor, β ($= 0.9$) that is introduced by Liston and Elder (2006b) to make LAI^* compatible with the canopy flow index proposed by Cionco (1978):

$$a = \text{LAI}^* \beta \quad (14)$$

LAI^* includes stems, leaves and branches as described by Chen et al. (1997).

To consider the extinction of solar radiation by the forest canopy, top-of-canopy incoming shortwave radiation, Q_{si} , is reduced following the Beer–Lambert law as

$$Q_{\text{sif}} = Q_{\text{si}} \tau_v, \quad (15)$$

where Q_{sif} is the incoming shortwave radiation impinging on the snow surface beneath the canopy (Hellström, 2001). τ_v representing the fraction of Q_{si} reaching the land surface is derived as

$$\tau_v = \exp(-k\text{LAI}^*), \quad (16)$$

with k being a vegetation-dependent extinction coefficient (Liston and Elder, 2006b). Aiming at a best fit to observed radiation inside forest canopies of different species (e.g. spruce, subalpine fir, pine) at a site in the U.S. Department of Agriculture (USDA) Fraser Experimental Forest near Fraser (Colorado, USA), Liston and Elder (2006b) have yielded best overall performance using a k value of 0.71, which is also used for the simulations here.

Incoming longwave radiation inside the canopy is assumed to be composed by a fraction F_g (-) directly reaching the ground through gaps in the forest stand and a fraction F_c (-) emitted by the forest canopy. The canopy-emitted fraction is calculated following Liston and Elder (2006a) as

$$F_c = a + b\ln(\text{LAI}^*) \quad (17)$$

where a (-) and b (-) are constants with values of 0.55 and 0.29, respectively. A value of F_g can be derived as

$$F_g = 1 - F_c, \quad (18)$$

with both calculated fractions used to estimate inside-canopy incoming longwave radiation Q_{lif} (W m^{-2}) from

$$Q_{\text{lif}} = (F_g Q_{\text{li}}) + (F_c \sigma T_c^4), \quad (19)$$

where Q_{li} (W m^{-2}) represents the top-of-canopy incoming longwave radiation. The latter is provided as input for ESCIMO.spread (v2) and is here estimated as a function of

A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

Title Page	
Abstract	Introduction
Conclusions	References
Tables	Figures
◀	▶
◀	▶
Back	Close
Full Screen / Esc	
Printer-friendly Version	
Interactive Discussion	



temperature and cloud cover as proposed by Liston and Elder (2006a) due to a lack of observations. σ represents the Stefan Boltzmann constant and T_c (K) the inside-canopy temperature. Assuming a linear dependency on canopy fraction, T_c is derived from top-of-canopy temperature T_a (K) as proposed by Obled (1971):

$$T_c = T_a - F_c(T_a - (R_c(T_a - T_{\text{mean}}) + T_{\text{mean}} - \delta T)), \quad (20)$$

where T_{mean} (K) is the mean daily air temperature, R_c (-) is a dimensionless scaling parameter set to 0.8 and δT ($-2\text{K} \leq \delta T \leq +2\text{K}$) is a temperature offset defined to be (Durot, 1999)

$$\delta T = \frac{T_{\text{mean}} - 273.16}{3}. \quad (21)$$

Durot (1999) has further shown that relative humidity inside the canopy, RH_c (%), is often higher compared to the open due to sublimation and evaporation of melted snow. We therefore propose to modify top-of-canopy humidity RH (%) with consideration of the canopy fraction in the form of (Durot, 1999)

$$\text{RH}_c = \max\{\text{RH}(1 + 0.1F_c), 100\}. \quad (22)$$

2.5 Simulating canopy effects on the snow cover

The following describes the newly implemented approaches to describe snow interception through the forest canopy as well as melt-induced unloading of intercepted snow from the canopy.

Interception of snow precipitation P (mm), at time t is derived introducing a canopy-intercepted load, I (mm), expressed as (Pomeroy et al., 1998)

$$I = I_{t-1} + 0.7(I_{\text{max}} - I_{t-1})(1 - \exp(-P/I_{\text{max}})), \quad (23)$$

where $t - 1$ represents the previous time step and I_{max} is the maximum interception storage calculated as (Hedstrom and Pomeroy, 1998)

$$I_{\text{max}} = 4.4\text{LAI}^*. \quad (24)$$

A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Sublimation of intercepted snow Q_{cs} (mm) is calculated as described by Liston and Elder (2006b) as

$$Q_{cs} = C_e / \Psi_s dt, \quad (25)$$

where dt (s) is the time increment (here: 3600 s), Ψ_s (s^{-1}) is the sublimation-loss rate coefficient for an ice sphere and $C_e(-)$ represents the canopy exposure coefficient. Ice spheres are assumed to be characterized by a constant radius of 500 μm as proposed by Liston and Elder (2006b).

The canopy exposure coefficient is calculated as

$$C_e = k_c (l/l_{\max})^{-0.4}, \quad (26)$$

where k_c (-) is a dimensionless coefficient related to the shape of the intercepted snow deposits (Liston and Elder, 2006b). Sublimation at the canopy scale is hence estimated based on sublimation from individual ice spheres. Analysing observed (Montesi et al., 2004) and modelled sublimation rates for a 2.7 m-tall subalpine fir tree at the USDA Fraser Experimental Forest, Liston and Elder (2006b) have found that the application of $k_c = 0.010$ seems to best reproduce observed sublimation rates at both, higher and lower elevated tree sites. This value is very close to the value of $k_c = 0.011$ derived by Pomeroy et al. (1998) for Canadian Boreal Forest and is used as k_c value for the calculations with ECIMO.spread (v2) here. This parameter can be easily adapted by changing the respective setting in the parameter section of the model.

The sublimation-loss rate coefficient Ψ_s is calculated from the particle mass m (kg) in the form of

$$\Psi_s = (dm/dt)/m, \quad (27)$$

where the particle mass is given by

$$m = \frac{3}{4} \pi \rho_i r^3, \quad (28)$$

A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

Title Page	
Abstract	Introduction
Conclusions	References
Tables	Figures
◀	▶
◀	▶
Back	Close
Full Screen / Esc	
Printer-friendly Version	
Interactive Discussion	



with ρ_i (kg m^{-3}) being ice density and r (m) representing the radius of a spherical ice particle (assumed to be $500 \mu\text{m}$ as proposed by Liston and Elder, 2006b).

Mass loss from an ice particle is described as a function of intercepted solar radiation, humidity gradients between the ice surface and the surrounding atmosphere, the size of the considered ice particle and a ventilation term, following Thorpe and Mason (1966) and Schmidt (1972):

$$\frac{dm}{dt} = \frac{2\pi \left(\frac{RH}{100} \right) - S_p \Omega}{h_s \Omega + \frac{1}{D \rho_v S h}}, \quad (29)$$

where h_s is the latent heat of sublimation ($2.8355 \times 10^6 \text{ J kg}^{-1}$).

The diffusivity of water vapour in the atmosphere, D ($\text{m}^2 \text{ s}^{-1}$) is derived following Thorpe and Mason (1966) as:

$$D = 2.06 \times 10^{-5} (T_a / 273)^{1.75}. \quad (30)$$

The molecular weight of water M ($18.01 \text{ kg kmole}^{-1}$), the universal gas constant R ($8313 \text{ J kmole}^{-1} \text{ K}^{-1}$), air temperature T_a (K) and the thermal conductivity of the atmosphere λ_t ($0.024 \text{ J m}^{-1} \text{ s}^{-1} \text{ K}^{-1}$) are used to calculate Ω as proposed by Liston and Elder (2006b):

$$\Omega = \frac{1}{\lambda_t T_a Nu} \left(\frac{h_s M}{RT_a} - 1 \right). \quad (31)$$

The Nusselt number Nu and Sherwood number Sh are both calculated as:

$$Nu = Sh = 1.79 + 0.606 Re^{0.5}, \quad (32)$$

where Re ($0.7 < Re < 10$) is the Reynolds number expressed by:

$$Re = \frac{2ru_c}{\nu} \quad (33)$$

A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



with ν representing the kinematic viscosity of air ($1.3 \times 10^{-5} \text{ m}^2 \text{ s}^{-1}$) and u_c the ventilation velocity inside the canopy, which is set equal to inside-canopy wind speed as proposed by Liston and Elder (2006b).

Following Fleagle and Businger (1981) the saturation density of water vapor ρ_v (kg m^{-3}) is derived as

$$\rho_v = 0.622 \frac{e_s}{R_d T_a} \quad (34)$$

where R_d is the gas constant for dry air ($287 \text{ J K}^{-1} \text{ kg}^{-1}$) and e_s (Pa) is the saturation vapor pressure over ice, estimated following Buck (1981) as:

$$e_s = 611.15 \exp\left(\frac{22.452(T_a - 273.16)}{T_a - 0.61}\right). \quad (35)$$

The shortwave radiation absorbed by a snow particle with radius r is defined to be

$$S_p = \pi r^2 (1 - \alpha_p) S_i, \quad (36)$$

where α_p is the snow albedo, and S_i (W m^2) is the solar radiation at the earth surface, which in case of ESCIMO.spread (v2) is among the required meteorological input parameters.

To account for a melt-induced unloading of intercepted snow from the canopy, a melt-unloading rate L_m (kg m^{-2}) is introduced by Liston and Elder (2006b):

$$L_m = 5.8 \times 10^{-5} (T_a - 273.16) dt. \quad (37)$$

We assume an unloading rate of $5 \text{ kg m}^{-2} \text{ day}^{-1} \text{ K}^{-1}$ whenever temperatures are above freezing, with unloading snow adding to snow accumulation at the land surface. The simulated filling and depletion of the interception storage through snow fall, sublimation and melt induced unload is illustrated in Fig. 4 exemplarily for a period in February 2013.

A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



3 Data and test site description

Snow cover simulations in this study are carried out for the forest site Vordersteinwald in the Black Forest mountain range (south-west Germany) (see Fig. 1). This site is eminently suitable for testing of the newly developed version of ESCIMO.spread as it (i) usually experiences alternation of accumulation and melting periods over the winter season, making the simulation of snow conditions particularly demanding and (ii) has been subject to intense snow surveys over the years 2010–present, including simultaneous observation of meteorological and snow conditions in and outside the forest canopy (Pohl et al., 2014).

The forest stand at the study site is mostly conifer with spruce, fir and pine, representing the most common conifer tree species. To quantify the vegetation effect on snow conditions, the applied snow monitoring systems (SnoMoS) were installed pairwise with one SnoMoS located in the open and another set up in close distance inside the forest canopy (see Fig. 5). The data recorded by these low-cost monitoring sensors include hourly values of snow depth, surface temperature, air temperature and humidity, global radiation, wind speed, and barometric pressure.

The continuous monitoring of snow depth with the SnoMoS was accompanied by bi-weekly snow density surveys that allow translation of snow depth into snow water equivalent. A comprehensive description of the technical specifications and the instrumental setup of the SnoMoS is provided by Pohl et al. (2014). Precipitation recordings for the study site originate from nearby weather station Freudenstadt (DWD, 2015), operated by the German Weather Service (DWD). Precipitation observations have been corrected for differences in terrain elevation between the sites of measurement and model application by applying monthly elevation adjustment factors as proposed by Liston and Elder (2006a). The latter have been taken from Marke (2008) who has investigated altitudinal differences in precipitation for the Upper Danube Watershed. No interpolation using other station data has been carried out due to the closeness of the study site (3 km distance) to station Freudenstadt. Hemispherical images were

GMDD

8, 8155–8191, 2015

A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



equivalent in this period might be a false interpretation of rainfall as snowfall. While the model acceptably reproduces snow accumulation between 10 and 30 January in the open, a noticeable overestimation of snow water equivalent can be observed in the results using the modified outside-canopy meteorological conditions. While a period of snow accumulation can be observed in the observations and simulations for the open in March, inside the canopy this increase in snow water equivalent is merely predicted by the model and not confirmed by the observations. A comparison of the snow simulations based on observed and simulated meteorological conditions inside the canopy reveals that only little differences exist between both model runs, with the model performance using observed meteorological conditions as model input being slightly better. Taking a closer look at the efficiency criteria in Table 2, the model results for both locations seem to well reflect the observed conditions with model performance inside the canopy being even slightly better than in the open.

5 Conclusions

A new version of the spreadsheet-based point energy balance snow model ESCIMO.spread has been presented (ESCIMO.spread (v2)) that allows an improved precipitation phase detection, consideration of cold and liquid water content in the snow cover, estimation of inside canopy meteorological conditions from meteorological observations in the open and the simulation of snow accumulation and ablation inside a forest canopy. It thereby does not require meteorological observations in the canopy but instead derives inside-canopy meteorological conditions from available observations in the open requiring only LAI and canopy height as plant-specific input parameters. The derived meteorological conditions inside the canopy are not only applicable as input for snow cover simulations but can be expected to be of interest for a variety of scientific disciplines, e.g. forest ecology or pedology. To provide the data required for model application and evaluation, a pair of SnoMoS has been utilized as an innovative technology that allows the collection of important meteorological variables at low

A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



financial costs. Comparison of simulated inside-canopy meteorological conditions to observations at a site in the Black Forest region (Germany) reveals good overall model performance, particularly with respect to global radiation and temperature representing the most important meteorological variables for the estimation of snow melt. A similar picture unfolds when comparing the simulated snow cover evolution in- and outside the canopy to snow observations provided by the SnoMoS. Here, snow cover evolution is well reproduced for both, out- and inside the forest canopy, with slightly higher model performance for inside-canopy conditions, even though the empirical model parameters have not yet been adjusted to (pre)alpine forest species. Making use of the full potential of simultaneous observation of snow and meteorological conditions as provided by the SnoMoS, an effort is currently undertaken to develop parameters for the applied canopy model that are tailored to the specific conditions in (pre)alpine forests. Moreover, despite its physically-based character and advanced model features, ES-CIMO.spread (v2) still oversimplifies some important processes of the snow-vegetation interaction. In the current version the model only considers unloading of intercepted snow as a result of melting. While the fact that wind also induces unloading of intercepted snow is well known, the combined dependence on plant characteristics (e.g. plant structure and plant element flexibility) and meteorological conditions (e.g. snow temperature, wind speed and direction) makes this a complex process hard to consider in numerical models (Liston and Elder, 2006b). The modification of shortwave and longwave radiation assumes a plant specific extinction coefficient and a constant canopy fraction, respectively. While these assumptions can be expected to reasonably reproduce the general observed trends in local radiation, they are not capable to accurately capture the actual radiation conditions whenever canopy densities strongly vary or sun is shining through open areas in the trees as a result of changing solar zenith angles.

Code availability

ESCIMO.spread (v2) can be downloaded free of charge at www.alpinehydroclimatology.net together with one year of sample data including the meteorological and snow observations used in this study. The model has been tested on OpenOffice 4.1.1. as well as on different versions of Microsoft Excel for Windows and Mac.

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A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



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GMDD

8, 8155–8191, 2015

**A spreadsheet-based
snow model for
inside-canopy
conditions**

T. Marke et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Table 1. Performance of ESCIMO.spread (v2) in the modification of outside-canopy global radiation, temperature, relative humidity and wind speed for canopy effects.

Variable	NSME	R^2	IA
Global radiation	0.64	0.66	0.89
Air temperature	0.79	0.82	0.94
Relative humidity	-1.10	0.61	0.74
Wind speed	-0.29	0.60	0.78

A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Table 2. Performance of ESCIMO.spread (v2) in the simulation of snow cover evolution in- and outside the forest canopy. The simulations inside the canopy are based on modified outside-canopy meteorological conditions.

Variable	NSME	R^2	IA
SWE (outside canopy)	0.65	0.77	0.88
SWE (inside canopy)	0.80	0.82	0.95

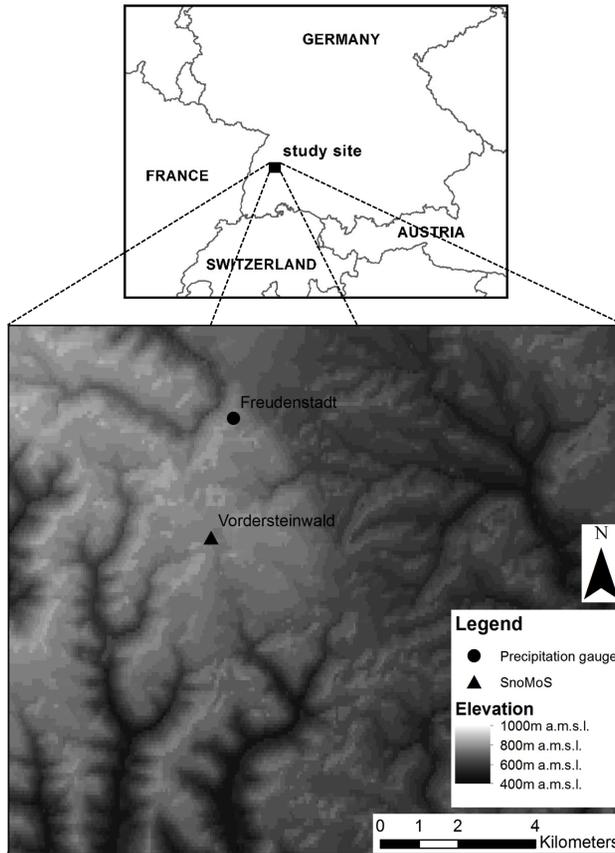


Figure 1. The site Vordersteinwald in the Black Forest mountain range (south-west Germany, 800 m a.s.l.).

GMDD

8, 8155–8191, 2015

A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

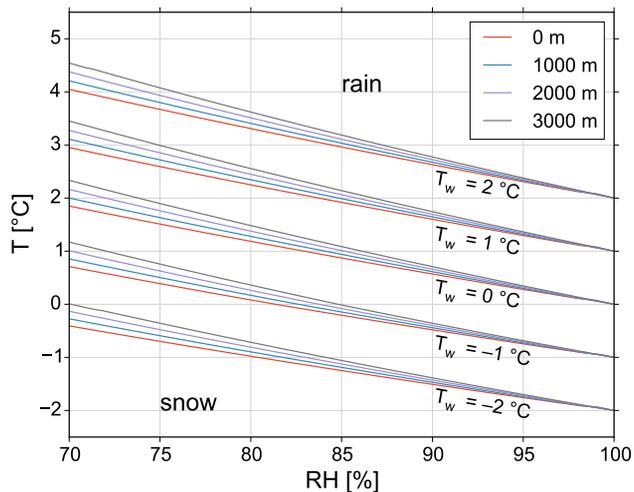


Figure 2. Relation between air temperature, wet-bulb temperature and relative humidity in different altitudes. The latter represent different air pressure levels derived using the hydrostatic equation.

[Title Page](#)[Abstract](#)[Introduction](#)[Conclusions](#)[References](#)[Tables](#)[Figures](#)[◀](#)[▶](#)[◀](#)[▶](#)[Back](#)[Close](#)[Full Screen / Esc](#)[Printer-friendly Version](#)[Interactive Discussion](#)

A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

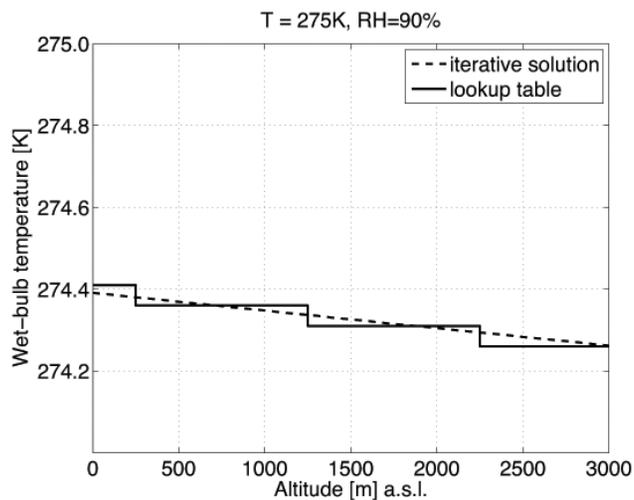


Figure 3. Comparison of iteratively calculated wet bulb temperature to the results of the lookup table approach implemented in ESCIMO.spread (v2).

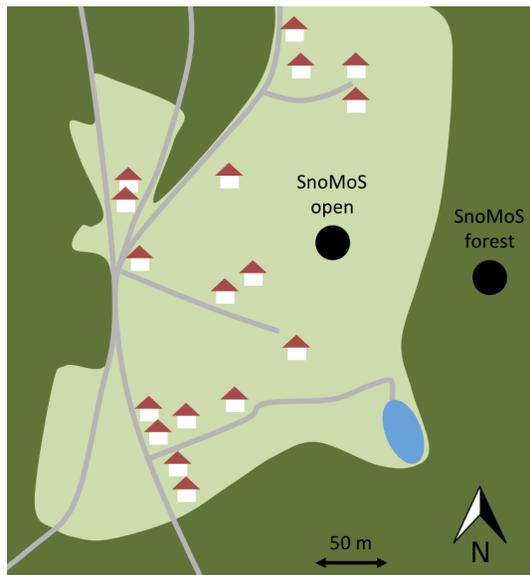


Figure 5. Schematic overview of the SnoMoS setup locations in- and outside the forest canopy at site Vordersteinwald in the Black Forest mountain range (south-west Germany, 800 m a.s.l.). The light green areas indicate grassland, the dark green areas forest, the grey lines streets and the light blue area a lake.

A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

Title Page	
Abstract	Introduction
Conclusions	References
Tables	Figures
◀	▶
◀	▶
Back	Close
Full Screen / Esc	
Printer-friendly Version	
Interactive Discussion	



A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

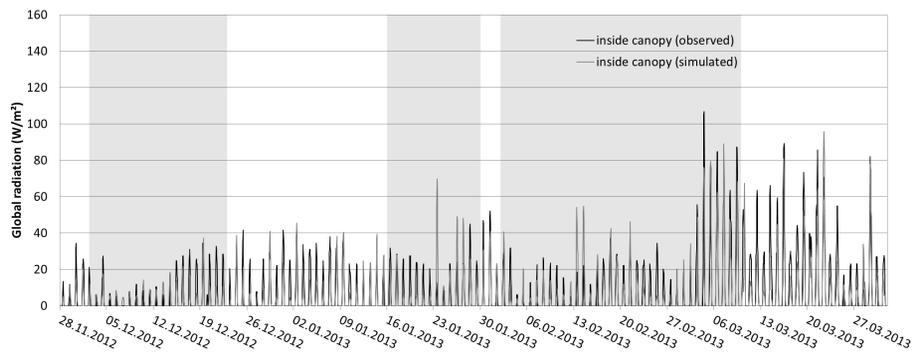


Figure 6. Simulated and observed global radiation for the winter period 2012/13. The grey areas indicate periods with presence of a snow cover.

[Title Page](#)[Abstract](#)[Introduction](#)[Conclusions](#)[References](#)[Tables](#)[Figures](#)[Back](#)[Close](#)[Full Screen / Esc](#)[Printer-friendly Version](#)[Interactive Discussion](#)

A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

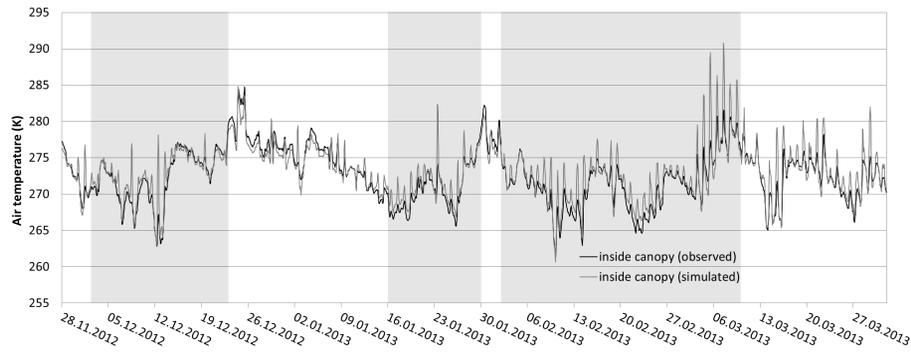


Figure 7. Simulated and observed temperature inside the forest canopy for the winter period 2012/13. The grey areas indicate periods with presence of a snow cover.

[Title Page](#)[Abstract](#)[Introduction](#)[Conclusions](#)[References](#)[Tables](#)[Figures](#)[Back](#)[Close](#)[Full Screen / Esc](#)[Printer-friendly Version](#)[Interactive Discussion](#)

A spreadsheet-based snow model for inside-canopy conditions

T. Marke et al.

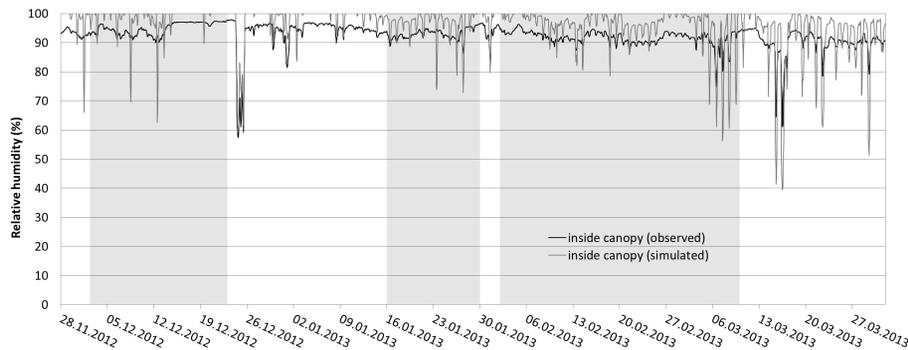


Figure 8. Simulated and observed relative humidity inside the forest canopy for the winter period 2012/13. The grey areas indicate periods with presence of a snow cover.

[Title Page](#)[Abstract](#)[Introduction](#)[Conclusions](#)[References](#)[Tables](#)[Figures](#)[Back](#)[Close](#)[Full Screen / Esc](#)[Printer-friendly Version](#)[Interactive Discussion](#)

