Response to Reviewers

Title: Simple Process-Led Algorithms for Simulating Habitats (SPLASH v.1.0): Robust Indices of Radiation, Evapotranspiration and Plant-Available Moisture


Manuscript Number: GMD-2016-49

We thank the reviewers for their constructive and valuable comments. Below, we have responded to each of the comments and concerns. For clarity, we have included a summary of the comments in italics. We would also like to point out the following additional changes to the manuscript:

1. We have added a new co-author, Rebecca T. Thomas, who is responsible for coding and producing the global SPLASH simulations.
2. We moved the variable substitutions, \( r_u \) and \( r_v \), from Eq. 8 into Eq. 7 to improve the clarity of the later equation derivations.
3. We have updated the definition of the net surface radiation, \( H_N \), in Section 2.1.2 for clarity, such that it now encapsulates a full 24-hour period (not just daytime), and have separated it into its positive, \( H_{N^+} \), and negative, \( H_{N^-} \), components. We have updated the notation of negative net radiation, from \( H_N^* \) to \( H_{N^-} \), in Eqns. 16 and 18. The distinction between \( H_N \) and \( H_{N^+} \) is also made in Eq. 24 in Section 2.4 as well as in Table 1. We have also made corrections to Eq. 16 for the negative net radiation.
4. Eq. 35 in the Discussions has been moved to Section 2, as it seemed more appropriate to make note of it here.
5. Figures 1 and 2 have been updated for improved clarity.
6. Figure 4 has been updated to match the style of the new Figure 3.
7. The results section has been broken into two subsections: one for local and one for global simulations.

Reviewer #1

Comment #1.1:

"[...] the process-based model is greatly simplified; therefore you need to validate whether the model output is reliable or not using globally or locally observed data."

The reviewer suggests that despite the lack of observations of evapotranspiration and plant-available moisture, there are observations of surface and top-of-the-atmosphere radiation available (e.g., CERES) that could be compared to the SPLASH model radiation simulations. The reviewer notes that observations of monthly radiation, evapotranspiration, and soil moisture should accompany Fig. 4, else another site should be selected where data is available.

In our search for daily hydro-meteorological data, we discovered the publication by Ben Livneh and others (doi: 10.1038/sdata.2015.42) consisting of 1/16th degree gridded data over North America (i.e., Southern Canada, United States, and Mexico) that includes observations (e.g., precipitation, maximum and minimum air temperature) from NOAA National Climate Data Center’s (NCDC) Global Historical Climatology Network (GHCN) along with simulated fluxes (e.g., net radiation, runoff, evapotranspiration, and soil moisture) produced by the VIC model. We have plotted these data along side our own for comparison of local scale trends (see figure below). The figure and analysis have been updated in the manuscript.
The monthly sunshine fraction (Fig 3a) is unchanged. Fig. 3b now shows the SPLASH net radiation flux along with the VIC model net radiation from Livneh et al. (2015), converted from units of W m$^{-2}$ to MJ m$^{-2}$ (red line). The SPLASH net radiation exhibits slightly higher values, particularly in the later season, likely due to the skewness of the monthly sunshine fraction. Fig. 3c results have changed slightly due to a correction in the negative net radiation equation. Fig. 3d now includes daily precipitation from Livneh et al. (red line), which has some higher peak values during the winter months, but otherwise is relatively consistent in the timing of rain events to the daily WATCH precipitation (black line) throughout the year. Fig. 3e shows results from Livneh et al. (2015) for the three-layer VIC model (red lines) and SPLASH (black line). SPLASH consistently models soil moisture between the first and second layers (red solid and dashed lines, respectively), except for the rainy season where the SPLASH soil moisture is higher in magnitude, now between the second and third layers (red dashed and dotted lines, respectively). Fig. 3f indicates that SPLASH runoff (black line) is higher in magnitude at peak rainfall events during the wet season compared to the VIC modeled runoff (in red). In Fig. 3g, the red colored region depicts the range of minimum to maximum near-surface air temperature from Livneh et al., which does fairly well at enveloping the daily WATCH mean near-surface air temperature (black line). Lastly, in Fig. 3h, the VIC potential evapotranspiration curve (red line) shows a consistent seasonal course compared to the SPLASH potential evapotranspiration curve (solid black line), albeit with a significantly higher variability during the summer months.

We took the advice of the reviewer and acquired global CERES net surface radiative fluxes, which we have plotted beside the SPLASH model simulations for the months of June and December 2000 (see figure below).
The top row shows the monthly mean net downward surface radiative flux, MJ m$^{-2}$, for June 2000 and the bottom row for December. The left column results are from the CERES EBAF and the right column are from SPLASH. As shown in the figure, SPLASH does a reasonably good job at capturing the latitudinal gradients and temporal shifts of net surface radiation. The hot-spots that SPLASH simulates over the deserts and tundra, which are not seen in the CERES EBAF results, are likely due to the invalidation of the well-watered constant surface emission assumption.

**Comment #1.2:**

"[...] do we really need mean daily meteorological variables for the robust approximations of key quantities?"

The reviewer questions the validity of mixing monthly and daily inputs as is done in the results section (i.e., monthly cloudiness is used in combination with daily precipitation and air temperature). The reviewer notes that for paleoclimate studies, it is common to have mean monthly data and, therefore, recommends implementing a quasi-daily conversion of monthly data to meet the daily input requirements of SPLASH. The necessity of implementing quasi-daily conversion of monthly inputs for paleoclimate studies raises the question of the influence that the quasi-daily conversion will have on the model outputs when compared to regular daily-time stepped meteorological data.

We thank the reviewer for mentioning this, as we did not clarify our choice of input datasets in the manuscript. It is correct that, in most cases, mean monthly meteorology will be the only datasets available for driving the model, and to answer the question, mean daily data is not required. Traditionally, as it was noted, quasi-daily methods are used to convert mean monthly quantities to proximal daily values. This is still a viable and recommended method for producing the input datasets. However, for explanatory purposes, we chose to use daily WATCH meteorology for precipitation and near-surface air temperature, as they were available to us and also to help illustrate how daily variability propagates through the model.

Further to the point in regards to quasi-daily methods, the choice of methodology can/will affect model results. As there are varied methods for producing quasi-daily data (e.g., constant, linear interpolation, empirical model, weather generator), we leave this decision and its influence on the model results to the user. We have included text in Section 1 and Section 4 of the manuscript to explain our choice of datasets and to mention the use of quasi-daily methods in lieu of daily meteorological observations. For temperature and cloudiness, we feel that the constant daily or linearly interpolated values will produce reasonably similar results; however, to capture the stochastic nature of precipitation, a weather generator may be required, which may be included in later versions of the SPLASH model.
Comment #1.3:

“P. 1, L5, what are the exact time scales of 'ecologically relevant time scales'? Under your assumption, can we use a traditional climatology (i.e., 30-year average) data?”

By “ecologically relevant” we mean time scales from months to decades. SPLASH can be driven with transient (month-by-month) data, if used in the modeling of e.g. carbon and water fluxes or tree rings. It can equally be driven with a multi-decadal average climatological seasonal cycle, if used in the modeling of geographic distributions of functional traits or species. We have added these examples of ecological time scales in the text of Section 1.

Comment #1.4:

“P. 3, L6, change 'air temperature' to 'near-surface air temperature' or 'air temperature at the height of 2m'”

We thank the reviewer for noting this correction; it has been addressed in the text of Section 1. This distinction has also been noted in the abstract.

Comment #1.5:

“P. 3, L27, one term/character missing for 'the analytical integral of the minimum ... over a single day’”

We thank the reviewer for noting this. This was a typesetting mistake and has been corrected.

Comment #1.6:

“P. 5, L25, surface shortwave albedo is constant (0.17) for the model simplification, but is it okay for the energy balance at the local scale? The surface condition including albedo must be different locally. Therefore, the anomaly (actual surface albedo at local minus 0.17) largely impacts on surface net radiation and thus evapotranspiration and soil moisture in some regions, right?”

This is a good point. Changes in the shortwave albedo will influence the net radiation, which in turn influences actual evapotranspiration and condensation and, therefore, influences soil moisture. A quick sensitivity analysis of shortwave albedo showed only subtle changes in these variables when the value was halved (i.e., 0.08); however, this was not comprehensively analyzed in this study.

The purpose of the global constants is for model simplification. Their values represent reasonable “global means” and provide an approximation should the researcher have no other information to go on. However, these values can be specified for a locality, if and when the data are available. Due to the rarity of global datasets of such values, we feel assigning constant values for the sake of simplicity is justifiable. We have added text in Section 5 that clarifies the use of localized datasets in replacement of global constants.

Comment #1.7:

“P. 13, L13, although the model requires daily meteorological input (P. 3, L6-7), why do you use monthly cloud fraction data here? Do you think that cloud fraction is less significant for the calculation, compared to the other meteorological variables? I think that the diversity of the forms of clouds and their strong spatial and temporal variability determine the dynamics of the radiation budget to a significant degree.”

As mentioned in the response to Comment #1.2, we chose mean daily meteorology because of its availability and to help exemplify model results. The use of monthly cloudiness was partly out of convenience and partly due to the absence of a better alternative. We did not use monthly cloudiness because we thought it less significant.
Comment #2.1:

The reviewer suggests “expanding the manuscript to include figures displaying model output for the globe [... as well as some evaluation of the results [...] it is important to provide evidence that the model works across the range of global climate conditions.”

The reviewer further notes the importance of evaluating model performance such that other researchers have an indication as to whether SPLASH over- or under-estimates certain variables or has any spatial biases.

We have investigated the literature for global observations for comparing against the SPLASH model. The first comparison is with CERES EBAF surface net downward radiative flux (described in Reviewer Comment #1.1). The second comparison is with NCEP CPC V2 soil moisture (van den Dool et al., 2003; doi: 10.1029/2002JD003114) and is shown in the following figure.

The top row shows the relative mean daily soil moisture, %, for June 2000 and the bottom row for December. The left column is the NCEP CPC soil moisture and right column is from SPLASH (following a spin-up to equilibrate the soil moisture fields). We have plotted the relative soil moisture instead of the magnitude due to the significant differences in bucket size (i.e., 760 mm in NCEP CPC and 150 mm in SPLASH). The SPLASH soil moisture simulations result in a relatively full bucket throughout the wet regions. We contribute the comparatively empty bucket in the NCEP CPC soil moisture results to its larger bucket size. Nevertheless, the spatial patterns and seasonal shifts of soil moisture is consistent between the two models. There is a bias in the SPLASH soil moisture in the north-eastern region of Russia, which may be due to the lack of a long-term spin-up of soil moisture.
Addendum:

In our Consolidated Response to referee comments, we noted a discrepancy between the simulated patterns of soil moisture in northeastern Russia and the patterns shown in the NCEP re-analysis. The SPLASH-simulated soil moisture Figure presented there contained unrealistically sharp boundaries, which we had previously attributed to a spin-up issue. However, we have now traced the problem to an incorrect specification of the specific heat capacity of air at low air temperatures, which we have now corrected. We have also imposed a mask over the Greenland ice sheet where the simulation of soil moisture does not make sense. The following Figure is the corrected version, which will be included in the revised manuscript.

Comment #2.2:

“[… ] SPLASH has been designed so that it can be used for palaeoclimate applications [… however, the …] code is currently set up so that the orbital parameters used for palaeoclimate simulations are not input as variables but are specified […] as parameters.”

The reviewer suggests that these palaeoclimate-specific parameters (i.e., obliquity, eccentricity and longitude of the perihelion) be clarified in the manuscript and that a description be added as to how these parameters may be changed in the code for palaeo applications. The reviewer further requests that an explanation be presented on how users are meant to input dates for palaeoclimate studies (e.g., does the Meeus (1991) Julian day algorithm work for negative years?)

Thank you for making this clear to us. The code has been updated to identify the paleoclimate variables; however, it is up the user as to which is the appropriate method for updating these values—we made reference to Berger (1978) and Berger & Loutre (1991) for possible algorithms in Section 2.1.1. We added a note on our use of Berger (1978) to calculate the constant orbital parameters in our results.

In Meeus (1991), the Julian day algorithm is valid for positive and negative Gregorian calendar years, but not for negative Julian days (i.e., invalid for dates on or before noon –4712 January 1). Other methods exist, such as ignoring leap years. We understand the difficulty of tracking individual dates into the far past and admit that the current implementation in SPLASH only partially addresses the needs of paleoclimatological studies. The limitation of the Meeus algorithm has been explicitly defined in the source code.
Comment #2.3:

"The reviewer requests clarity be added to the required climatic drivers in regards to the use of fractional bright sunshine hours or fractional cloud cover, which currently reads as interchangeable quantities—actual model input requires fractional bright sunshine hours."

The ambiguity regarding the required climatic drivers has been addressed in the text of the abstract. We feel that in the introduction and results, our explanation of the difference between fractional bright sunshine hours and cloudiness fraction is sufficient for readers to understand.

Comment #2.4:

"The reviewer notes that Eq. 1 is missing the runoff term (as described in §2.6)."

Originally, we had lumped runoff with the correction of daily soil moisture. However, it is clearer to the reader, especially in regards to the bucket model, if we include the runoff parameter in Eq. 1, instead of as a by-product of our soil moisture calculation. Therefore, we have corrected Eq. 1 in Section 2, such that it now includes runoff as a parameter. To accommodate this, we have added a new subsection for runoff in the methodology (Section 2.6) and have amended the text in Section 2.7.

Comment #2.5:

"The reviewer points out that one of the objections made against a soil-dependent bucket size (i.e., that the seasonal course of soil moisture is insensitive to the exact value specified) may not be sufficient, as the seasonal pattern of soil moisture does, in fact, change with respect to a changing bucket size, which could have implications on the annual water deficit (and/or phenology)."

The absolute values of modeled soil moisture are influenced by the bucket size, which may also influence the soil moisture memory (this may be important in long-term simulations). However, the evapotranspiration (ET) rate is determined by the atmospheric demand and the fractional volumetric soil moisture content. In many applications, the quantity $\alpha$ (the ratio of supply-limited to non-supply-limited ET) is used as an index of water availability, and it has commonly been found that this value is insensitive to the bucket size. We have amended our defense of using a soil-independent bucket size and have moved it from Section 2.3 to the Discussions.

Comment #2.6:

"The reviewer suggests uploading the release version of the code (in all of its forms) to the journal in addition to hosting it on the Git repository to make certain of its future preservation."

Thank you for this comment. We plan to upload our source code to GMD for accessibility and preservation.

Comment #2.7:

"Page 13, line 13: Add the name of the specific WATCH data set you used."

Thank you for this comment; we have addressed this ambiguity by noting the use of the WATCH Forcing Data methodology applied to the ERA-Interim, first release, 2012.
Comment #2.8:

Page 17, line 10: I would add text describing the typographical error (e.g., Eq. 7 of Gallego-Sala et al. (2010) used 273.3 instead of 237.3).

Thank you for this comment; as per your recommendation, the typographical error has been noted.

Comment #2.9:

Figure 3: Change "CRU TS" to "CRU TS3.21" in the caption text.

Thank you for this suggestion. Both the CRU TS and WFDEI data source versions have been added to the figure caption to improve their clarity.

Comment #2.10:

There are a number of debugging comments in the code I downloaded from the Git repository, such as "consistency check – XXX PROBLEM: THIS LEADS TO DIFFERENCE WITH OTHER VERSIONS XXX." To prevent confusion for the user, remove these comments if the issues have been resolved. If the issues have not been resolved, provide enough detail in the code comment so that the user can determine how the issue may affect their results.

Our apologies; the old debugging comments have been removed from the source code.

Comment #2.11:

In various places in the code the user is referred to particular equations in the documentation file (splash_doc.pdf) that accompanies the code. However, in some cases the equation referenced in the code does not match the equation in the documentation. For example, the calculation of daily photosynthetic photon flux density (PPFD) in the FORTRAN code refers the reader to Equation 57 in the documentation file, which is an equation for the bulk modulus of water. Check that the references in the code to the splash_doc.pdf file are correct.

Thank you for noting this. Outdated references to the documentation have been removed from the source code.

Comment #2.12:

It would help the user if all of the variable names were defined in the code. For example, in the code where PPFD is defined as a type real variable the accompanying comment defines PPFD as "daily PPFD (mol/m^2)" instead of "daily photosynthetic photon flux density (mol/m^2)."

Thank you for this suggestion. Variable names throughout the code have been written out in the comments to help identify them.

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Abstract. Bioclimatic indices for use in studies of ecosystem function, species distribution, and vegetation dynamics under changing climate scenarios depend on estimates of surface fluxes and other quantities, such as radiation, evapotranspiration and soil moisture, for which direct observations are sparse. These quantities can be derived indirectly from meteorological variables, such as near-surface air temperature, precipitation and cloudiness. Here we present a consolidated set of Simple Process-Led Algorithms for Simulating Habitats (SPLASH) allowing robust approximations of key quantities at ecologically relevant time scales. We specify equations, derivations, simplifications and assumptions for the estimation of daily and monthly quantities of top-of-the-atmosphere solar radiation, net surface radiation, photosynthetic photon flux density, evapotranspiration (potential, equilibrium and actual), condensation, soil moisture, and runoff, based on analysis of their relationship to fundamental climatic drivers. SPLASH, as presented here, is designed for application at discrete locations; however, the same methodology can naturally be applied to spatial grids. The climatic drivers include a minimum of three meteorological inputs: precipitation, air temperature, and either fraction of bright sunshine hours or fractional cloud cover. Indices, such as the moisture index, the climatic water deficit, and the Priestley-Taylor coefficient, are also defined. The SPLASH code is transcribed in C++, FORTRAN, Python, and R. One year of results from a specific location are provided at the local and global
In the original, highly cited BIOME model (Prentice et al., 1992); the general forest succession model (FORSKA) described by Prentice et al. (1993); and the Simple Diagnostic Biosphere Model (Knorr and Heimann, 1995). Despite the subsequent development of more complex Dynamic Global Vegetation Models (Cramer et al., 2001; Sitch et al., 2003; Woodward and Lomas, 2004; Quillet et al., 2010; Prentice and Cowling, 2013; Fisher et al., 2014) and Land Surface Models, the relatively simple algorithms in STASH continue to have many applications, including to new areas such as the distribution of plant functional traits (Harrison et al., 2010; Meng et al., 2015), assessment of climate-change impacts on specific biomes (Gallego-Sala and Prentice, 2012), large-scale water resources assessments (e.g. Ukkola et al., 2015) and simple first-principles modeling of primary production (Wang et al., 2014). The continuing utility of these algorithms owes much to their robustness, which in turn depends on the implicit assumption that vegetation functions predictably—so that, for example, evapotranspiration occurs at a potential rate under well-watered conditions, and is reduced as soil water is drawn down. STASH is thus unsuitable to answer questions like the effect of imposed vegetation changes on runoff, or modeling vegetation-atmosphere feedbacks. Much more complex models that dynamically couple soil, vegetation and atmospheric boundary layer processes exist for such applications; however, their complexity brings a burden in terms of lack of robustness and, potentially, large inter-model differences (Prentice et al., 2014).

Despite their long history of use, no single publication documents the algorithms of the STASH model. This work aims to fill that gap to allow for the continued development and use of these algorithms. As the new incarnation of STASH, SPLASH provides the same physically-based soil moisture accounting scheme with updated and corrected analytical expressions for the calculation of daily radiation, evapotranspiration, and soil moisture. Included in this documentation are the equation derivations,
variable definitions, and information regarding model assumptions and limitations. One notable improvement is that we have discontinued the approximation of constant angular velocity in the orbit of Earth around the Sun. This version is thus suitable for palaeoclimate applications, whereby orbital precession (as well as changes in obliquity and eccentricity) influences the seasonal distribution of insolation. SPLASH also includes explicit consideration of elevation effects on biophysical quantities.

Key model outputs include daily insolation (incoming solar radiation at the top of the atmosphere) and net surface radiation ($H_o$ and $H_N$, respectively); daily photosynthetic photon flux density ($Q_n$); daily condensation, soil moisture and runoff ($C_n$, $W_n$, and $RO$); and daily equilibrium, potential and actual evapotranspiration ($E_n^a$, $E_n^p$, and $E_n^a$). Unlike the STASH model, SPLASH explicitly distinguishes potential and equilibrium evapotranspiration, recognizing that under well-watered conditions the excess of the former over the latter is a requirement for foliage to be cooler than the surrounding air, as has long been observed under high environmental temperatures (e.g. Linacre, 1967).

Input values of latitude, $\phi$ (rad), elevation, $z$ (m), mean daily near-surface air temperature, $T_{air}$ ($^\circ$C), and fractional hours of bright sunshine, $S_f$ (unitless), are necessary used for calculating the daily quantities of net radiation and evapotranspiration. Daily observed Daily precipitation, $P_n$ (mm d$^{-1}$), is necessary used for updating daily soil moisture. $T_{air}$ and $P_n$ may be derived from various sources, including the freely available daily-averaged air temperature and precipitation reanalysis data from the Water and Global Change (WATCH) program’s meteorological forcing data set (Weedon et al., 2014). Meteorological variables are also available in the Climatic Research Unit (CRU) gridded monthly time series datasets (Harris et al., 2014), which may be downscaled to daily quantities by means of quasi-daily methods (e.g., linear interpolation). Cloud cover fraction, for example the simulated quantities given in the CRU TS3.21 dataset (Harris et al., 2014), may be used to approximate $S_f$. Penman’s one-complement approximation based on the cloudiness fraction is regarded here as a sufficient estimate of $S_f$ (Penman, 1948). The piecewise linear method of Hulme et al. (1995)—an adaptation of the Doorenbos-Pruitt estimation procedure (Doorenbos and Pruitt, 1977)—as used in the development of the CRU cloudiness climatology (New et al., 1999) gives similar results.

We present SPLASH comprehensively re-coded in a modular framework to be readable, understandable and reproducible. To facilitate varied application requirements (including computational speed), four versions of the code (C++, FORTRAN, Python, and R) are available in an online repository (see Code Availability). The algorithms as presented here focus on application to individual site locations, but a natural extension is towards spatially distributed grid-based datasets.

In line with the intention of the original STASH algorithms, we also present bioclimatic indices at the monthly and annual timescales to exemplify the analytical applications of the SPLASH model outputs.

2 Methodology

The implementation of the soil-moisture accounting scheme follows the steps outlined by Cramer and Prentice (1988), where daily soil moisture, $W_n$ (mm), is calculated based on the previous day’s moisture content, $W_{n-1}$, incremented by daily precipitation, $P_n$ (mm d$^{-1}$), and condensation, $C_n$ (mm d$^{-1}$), and reduced by daily actual evapotranspiration, $E_n^a$ (mm d$^{-1}$) and
runoff, $RO$ (mm):

$$W_n = W_{n-1} + P_n + C_n - E^n_a - RO,$$

where $P_n$ is a model input, $C_n$ is estimated based on the daily negative net radiation, and $E^n_a$ is the analytical integral of the minimum of the instantaneous evaporative supply and demand rates over a single day. and $RO$ is the amount of soil moisture in excess of the holding capacity. An initial condition of $W_n$ is assumed between zero and the maximum soil moisture capacity, $W_m$ (mm), for a given location and is equilibrated over an entire year by successive model iterations (i.e., model spin-up).

Under steady-state conditions, the SPLASH model preserves the water balance, such that $\sum (P_n + C_n) = \sum (E^n_a + RO)$.

To solve the simple ‘bucket model’ represented by Eq. 1, the following steps are taken at the daily timescale: calculate the radiation terms, estimate the condensation, estimate the evaporative supply, estimate the evaporative demand, calculate the actual evapotranspiration, and update the daily soil moisture. Daily quantities may be aggregated into monthly and annual totals and used in moisture index calculations.

2.1 Radiation

2.1.1 Top-of-the-atmosphere solar radiation

The calculation of $C_n$ and $E^n_a$ begin with modeling the extraterrestrial solar radiation flux, $I_o$ (W m$^{-2}$). The equation for $I_o$ may be expressed as the product of three terms (Duffie and Beckman, 2013):

$$I_o = I_{sc} d_r \cos \theta_z,$$

where $I_{sc}$ (W m$^{-2}$) is the solar constant, $d_r$ (unitless) is the distance factor, and $\cos \theta_z$ (unitless) is the inclination factor. Values for $I_{sc}$ may be found in the literature (e.g., Thekaekara and Drummond, 1971; Willson, 1997; Dewitte et al., 2004; Fröhlich, 2006; Kopp and Lean, 2011); a constant for $I_{sc}$ is given in Table 2.

The distance factor, $d_r$, accounts for additional variability in $I_o$ that reaches the Earth. This variability is due to the relative change in distance between Earth and the Sun caused by the eccentricity of Earth’s elliptical orbit, $e$ (unitless), and is calculated as (Berger et al., 1993):

$$d_r = \left( \frac{1 + e \cos \nu}{1 - e^2} \right)^2,$$

where $\nu$ (rad) is Earth’s true anomaly. True anomaly is the measure of Earth’s location around the Sun relative to its position when it is closest to the Sun (perihelion).

The last term, $\cos \theta_z$, attenuates $I_o$ to account for the Sun’s height above the horizon (measured relative to the zenith angle, $\theta_z$), accounting for the off-vertical tilt of Earth’s rotational axis, $\varepsilon$ (i.e., obliquity). The inclination factor is calculated as (Duffie and Beckman, 2013):

$$\cos \theta_z = \sin \delta \sin \phi + \cos \delta \cos \phi \cos h,$$
where \( \phi \) (rad) is the latitude, \( \delta \) (rad) is the declination angle, and \( h \) (rad) is the hour angle, measuring the angular displacement of the Sun east or west of solar noon (\(-\pi \leq h \leq \pi\)). Declination is the angle between Earth’s equator and the Sun at solar noon (\(h = 0\)), varying from \(+\varepsilon\) at the June solstice to \(-\varepsilon\) at the December solstice; the changing declination is responsible for the change in seasons. For the purposes of ecological modeling, \( \delta \) may be assumed constant throughout a single day. See e.g. Woolf (1968) for the precise geometric equation representing \( \delta \):

\[
\delta = \arcsin(\sin \lambda \sin \varepsilon),
\]

where \( \lambda \) (rad) is Earth’s true longitude (i.e., the heliocentric longitude relative to Earth’s position at the vernal equinox) and \( \varepsilon \) (rad) is obliquity (i.e., the slowly varying tilt of Earth’s axis). Several other methods are widely used for the estimation of \( \delta \) for a given day of the year (e.g., Cooper, 1969; Spencer, 1971; Swift, 1976) but are not recommended because they do not account for the change in Earth’s orbital velocity with respect to the distance between Earth and the Sun, while Eq. 5 does. The relationship between true longitude, \( \lambda \), and true anomaly, \( \nu \), is by the angle of the perihelion with respect to the vernal equinox, \( \tilde{\omega} \) (rad) (Berger, 1978):

\[
\nu = \lambda - \tilde{\omega}.
\]

While the three orbital parameters (i.e., \( e \), \( \varepsilon \), and \( \tilde{\omega} \)) exhibit long-term variability (on the order of tens of thousands of years), they may be treated as constants for a given epoch (e.g., \( e = 0.0167 \), \( \varepsilon = 23.44^\circ \), and \( \tilde{\omega} = 283.0^\circ \) for 2000 CE), which can or they may be calculated using the methods of Berger (1978) or Berger and Loutre (1991) for palaeoclimate studies. Berger (1978) presents a simple algorithm to estimate \( \lambda \) for a given day of the year (see Appendix A).

The daily top-of-the-atmosphere solar radiation, \( H_o \) (J m\(^{-2}\)), may be calculated as twice the integral of \( I_o \) measured between solar noon and the sunset angle, \( h_s \), assuming that all angles related to Earth on its orbit are constant over a whole day:

\[
H_o = \int_{\text{day}} I_o = 2 \int_{h=0}^{h_f} I_o = \frac{86400}{\pi} I_{sc} \int_{r_v}^{} (r_u h_s + r_v \sin h_s),
\]

where \( r_u = (\sin \delta \sin \phi) \) and \( r_v = (\cos \delta \cos \phi) \), both unitless.

The sunset angle can be calculated as the hour angle when the solar radiation flux reaches the horizon (i.e., when \( I_o = 0 \)) and can be found by substituting Eq. 4 into Eq. 2, setting \( I_o \) equal to zero, and solving for \( h \):

\[
h_s = \arccos \left(\frac{-r_u}{r_v}\right).
\]

where \( r_u = (\sin \delta \sin \phi) \) and \( r_v = (\cos \delta \cos \phi) \), both unitless. To account for the undefined negative fluxes produced by Eq. 2 for \( h \geq h_s \) and \( h \leq -h_s \), \( I_o \) should be set equal to zero during these nighttime hours. To account for the occurrences of polar day (i.e., no sunset) and polar night (i.e., no sunrise), \( h_s \) should be limited to \( \pi \) when \( r_u/r_v \geq 1 \) and zero when \( r_u/r_v \leq -1 \).

### 2.1.2 Net surface radiation

The daytime (positive) net surface radiation, \( H_N \) (J m\(^{-2}\)), is the integral of the net surface radiation flux received at the land surface, \( I_N \) (W m\(^{-2}\)), which is classically defined as the difference between the net incoming shortwave radiation flux, \( I_{SW} \)
\[ I_N = I_{SW} - I_{LW}. \] (9)

The calculation of \( I_{SW} \) is based on the reduction in \( I_o \) due to atmospheric transmittivity, \( \tau \) (unitless), and surface shortwave albedo, \( \beta_{sw} \) (unitless):

\[ I_{SW} = (1 - \beta_{sw}) \tau I_o. \] (10)

A constant value for \( \beta_{sw} \) is given in Table 2. Atmospheric transmittivity may be expressed as a function of elevation (to account for attenuation caused by the mass of the atmosphere) and cloudiness (to account for atmospheric turbidity). At higher elevations, there is less atmosphere through which shortwave radiation must travel before reaching the surface. To account for this, Allen (1996) presents an equation based on the regression of Beer’s radiation extinction function at elevations below 3000 m with an average sun angle of 45°, which can be expressed as:

\[ \tau = \tau_o \left( 1 + 2.67 \times 10^{-5} z \right), \] (11)

where \( z \) (m) is the elevation above mean sea level and \( \tau_o \) (unitless) is the mean sea-level transmittivity, which can be approximated by the Ångstrom-Prescott formula:

\[ \tau_o = c + d S_f, \] (12)

where \( c \) and \( d \) are empirical constants (unitless) and \( S_f \) is the fraction of daily bright sunshine hours (0 ≤ \( S_f \) ≤ 1). Values for \( c \) and \( d \) are given in Table 2.

The calculation of \( I_{LW} \) is based on the difference between outgoing and incoming longwave radiation fluxes attenuated by the presence of clouds, which may be empirically estimated by (Linacre, 1968):

\[ I_{LW} = \left[ b + (1 - b) S_f \right] (A - T_{air}), \] (13)

where \( A \) and \( b \) are empirical constants and \( T_{air} \) (°C) is the mean near-surface air temperature. The outgoing longwave radiation flux used to derive Eq. 13 assumes a constant ground emissivity, which is accurate under well-watered conditions. The incoming longwave radiation flux is modeled based on clear-sky formulae derived by Linacre (1968). Values for \( A \) and \( b \) are given in Table 2.

Similarly to \( H_m \), may be calculated can be decomposed into its net positive, \( H^+_{N} \) (J m⁻²), and net negative, \( H^-_{N} \) (J m⁻²), components (i.e., \( H_N = H^+_{N} + H^-_{N} \)). Assuming \( I_{LW} \) is constant throughout the day and making substitutions for \( I_o \) into Eq. 10, an expression for \( H^+_{N} \) may be derived as twice the integral of \( I_N \) between solar noon (i.e., \( h = 0 \)) and the net surface radiation flux cross-over hour angle, \( h_n \) (rad):

\[ H^+_{N} = 2 \int_{h=0}^{h_n} I_N = 86400 \pi \left[ (r_w r_u - I_{LW}) h_n + r_w r_v \sin h_n \right], \] (14)

\[ H^-_{N} \]
where \( r_w = (1 - \beta_{sw}) \tau I_{sw} d_r \) (W m\(^{-2}\)).

Here, \( h_n \) is the hour angle when \( I_{SW} \) equals \( I_{LW} \) (i.e., when and can be found by setting \( I_N = 0 \) and in Eq. 9 and solving for \( h \). following the substitution of same substitutions as used for \( h_s \) in Eq. 10 and Eq. 13 for \( I_{SW} \) and \( I_{LW} \), respectively, may be expressed as:

\[
h_n = \arccos \left( \frac{I_{LW} - r_w r_u}{r_w r_v} \right).
\]

To account for the occurrences when the net surface radiation flux does not cross the zero datum, \( h_n \) should be limited to \( \pi \) when \((I_{LW} - r_w r_u)/(r_w r_v) \leq -1 \) (i.e., net surface radiation flux is always positive) and zero when \((I_{LW} - r_w r_u)/(r_w r_v) \geq 1 \) (i.e., net surface radiation flux is always negative).

The night time (negative) net surface radiation, \( H_N^+ \) (J m\(^{-2}\)), consists of two parts. Complementary to \( H_N^+ \), \( \bar{H}_N^- \) may be calculated as twice the integral of \( I_N \) for \( h_s \leq h \leq h_n \) and \( I_{LW} \) for \( h_n \leq h \leq \pi \). The calculation consists of doubling the half day integrals.

which may be expressed between \( h_n \) and solar midnight, defined by the piecewise function of \( I_N \) between \( h_n \) and \( h_s \) and \(-I_{LW} \) between \( h_s \) and solar midnight (i.e., \( h = \pi \)), given as (note that \( H_N^- - H_N^+ \) is a negative quantity):

\[
H_N^- = 2 \left( \int_{h_n}^{h_s} I_N - \int_{h_s}^{\pi} I_{LW} \right) = \frac{86400}{\pi} \left[ r_w r_v \left( \sin h_s - \sin h_n \right) + r_w r_u \left( h_s - h_n \right) - I_{LW} \left( \pi - h_n \right) \right].
\]

Figure 1 shows an example of a half-day \( I_N \) curve used in the integrals defined in Eqns. 14 and 16. \( I_N \), which is at its peak at solar noon, crosses zero at \( h_n \) and reaches a minimum at \( h_s \). After sunset (i.e., \( h > h_s \)), when \( I_{SW} \) is zero, \( I_N \) is equal to \(-I_{LW} \). \( H_N^- - H_N^+ \) is represented as twice the integral under the positive net radiation curve (solid gray line), above the zero line (dashed black line), and between the vertical lines of solar noon and \( h_n \). \( \bar{H}_N^- - H_N^+ \) is represented as twice the integral below the zero line and above the negative net radiation curve (the two dash-dotted lines, dashed gray line).

2.1.3 Photosynthetically active radiation

The daily photosynthetically active radiation in units of photon flux density, \( Q_n \) (mol m\(^{-2}\) d\(^{-1}\)), is calculated based on the number of quanta received (moles of photons) within the visible light spectrum, which also corresponds to the action spectrum of photosynthesis (Monteith and Unsworth, 1990):

\[
Q_n = 1 \times 10^{-6} \text{fFEC} \left( 1 - \beta_{vis} \right) \tau H_o,
\]

where \( \beta_{vis} \) (unitless) is the visible light albedo and fFEC (µmol J\(^{-1}\)) is the flux-to-energy conversion factor (Ge et al., 2011). This factor takes into account both the portion of visible light within the total solar spectrum, approximately 50% (Stanhill and Fuchs, 1977), and the mean number of quanta in the visible light energy band, approximately 4.6 µmol J\(^{-1}\) (McCree, 1972). The \( 1 \times 10^{-6} \) converts the units of \( Q_n \) from µmol m\(^{-2}\) d\(^{-1}\) to mol m\(^{-2}\) d\(^{-1}\). Values for \( \beta_{vis} \) and fFEC are given in Table 2.
Figure 1. Example of a half-day the net radiation flux curve with time represented along between the x-axis hours of solar noon (i.e., $h = 0$) and solar midnight (i.e., $h = \pi$). The $I_N$ curve is equal to the difference between the net incoming shortwave radiation flux, $I_{SW}$ (solid red line), and the net outgoing longwave radiation flux, $I_{LW}$ (dotted blue line). Positive values of $I_N$, shown decreasing from solar noon to zero at the cross-over hour angle, $h_n$, is denoted with a bold solid gray line, while negative values of $I_N$, which continues to decrease shown decreasing from zero to a minimum at $I_{LW}$ between $h_n$ and the sunset hour angle, $h_s$, and constant between $h_s$ and solar midnight, is denoted with a bold dashed gray line. The dash-dotted solid black horizontal line marks the datum of zero radiation.

2.2 Condensation

The daily condensation, $C_n$, may be expressed as the water-equivalent of the absolute value of negative net radiation, $H_N$:

$$C_n = 1 \times 10^3 E_{con} |H_N|,$$

where $E_{con}$ (m$^3$ J$^{-1}$) is the water-to-energy conversion factor that relates the energy released or required for a unit volume of water to evaporate or condense at a given temperature and pressure, which may be expressed as:

$$E_{con} = \frac{s}{L_v \rho_w (s + \gamma)},$$

where $s$ (Pa K$^{-1}$) is the slope of the saturation vapor pressure-temperature curve, $L_v$ (J kg$^{-1}$) is the latent heat of vaporization of water, $\rho_w$ (kg m$^{-3}$) is the density of water, and $\gamma$ (Pa K$^{-1}$) is the psychrometric constant. Standard values may be assumed for certain parameters (e.g., $L_v = 2.5 \times 10^6$ J kg$^{-1}$; $\rho_w = 1 \times 10^3$ kg m$^{-3}$; $\gamma = 65$ Pa K$^{-1}$); however, equations for the temperature dependence of $s$ and $L_v$ (e.g., Allen et al., 1998; Henderson-Sellers, 1984) and the temperature and pressure dependence of $\rho_w$ and $\gamma$ (e.g., Kell, 1975; Chen et al., 1977; Allen et al., 1998; Tsilingiris, 2008) are available (see Appendix B).

The barometric formula may be used to estimate the atmospheric pressure, $P_{atm}$ (Pa), at a given elevation, $z$ (m), when observations are not available. Assuming a linear decrease in temperature with height, which is a reasonable approximation
within the troposphere (i.e., for \(z < 1.10 \times 10^4\) m), the following equation may be used (Berberan-Santos et al., 1997):

\[
P_{atm} = P_o \left( 1 - \frac{L z}{T_o} \right)^{\frac{M_a}{R L}},
\]

(20)

where \(P_o\) (Pa) is the base pressure, \(T_o\) (K) is the base temperature, \(z\) (m) is the elevation above mean sea level, \(L\) (K m\(^{-1}\)) is the mean adiabatic lapse rate of the troposphere, \(g\) (m s\(^{-2}\)) is the standard gravity, \(M_a\) (kg mol\(^{-1}\)) is the molecular weight of dry air, and \(R\) (J mol\(^{-1}\) K\(^{-1}\)) is the universal gas constant. Values for the constants used in Eq. 20 are given in Table 2.

### 2.3 Evaporative Supply

The evaporative supply rate, \(S_w\) (mm h\(^{-1}\)) is assumed to be constant over the day and can be estimated based on a linear proportion of the previous day’s soil moisture, \(W_{n-1}\) (Federer, 1982):

\[
S_w = S_c \frac{W_{n-1}}{W_m},
\]

(21)

where \(S_c\) (mm h\(^{-1}\)) is the supply rate constant (i.e., maximum rate of evaporation) and \(W_m\) (mm) is the maximum soil moisture capacity. Constant values for \(S_c\) and \(W_m\) are given in Table 2. Although in principle \(W_m\) could be formulated as a property of soil type (as was done, for example, in the original BIOME model), there are several objections to doing so. One is that the seasonal course of soil moisture in the ‘bucket model’ formulation is insensitive to the exact value specified for \(W_m\). Another is that although \(W_m\) has a standard definition (as the difference between field capacity and wilting point) in agronomy, the wilting point in reality depends on plant properties. Yet another is that the effective ‘bucket size’ depends on rooting behavior, which is highly adaptable to the soil wetness profile. Thus, we suggest that no meaningful improvement in realism is likely to be achieved by applying soil type dependent values of \(W_m\).

### 2.4 Evaporative Demand

The evaporative demand rate, \(D_p\) (mm h\(^{-1}\)), is set equal to the potential evapotranspiration rate, \(E^p\) (mm h\(^{-1}\)), as defined by Priestley and Taylor (1972). \(E^p\) usually exceeds the equilibrium evapotranspiration rate, \(E^q\) (mm h\(^{-1}\)), due to the entrainment of dry air in the convective boundary layer above an evaporating surface (Raupach, 2000, 2001). \(E^p\) is related to \(E^q\) by the Priestley-Taylor coefficient, which may be defined as one plus an entrainment factor, \(\omega\) (Lhomme, 1997):

\[
D_p = E^p = (1 + \omega) E^q.
\]

(22)

The constant value used for \(\omega\) is given in Table 2. The calculation of \(E^q\) is based on the energy-water equivalence of \(I_N\), ignoring the soil heat flux, (Lhomme, 1997):

\[
E^q = 3.6 \times 10^6 E_{con} I_N,
\]

(23)

where \(3.6 \times 10^6\) converts the units of \(E^q\) from m s\(^{-1}\) to mm h\(^{-1}\). Note that \(E^q\) is defined only for positive values (i.e., \(E^q = 0\) for \(I_N < 0\)). The Priestley-Taylor potential evapotranspiration is preferred in this context to the general Penman-Monteith equation for actual evapotranspiration (Penman, 1948; Monteith, 1965), which requires knowledge of stomatal and
aerodynamic conductances, or to any of the ‘reference evapotranspiration’ formulae (Allen et al., 1998) that specifically relate to agricultural crops.

Daily equilibrium evapotranspiration, $E_q^n$ (mm d$^{-1}$), is based on the integration of Eq. 23 for values of positive $I_N$, or simply the energy-water equivalence of $H_N^+$. The daily demand, which is equal to the daily potential evapotranspiration, $E_p^n$ (mm d$^{-1}$), may be calculated from $E_q^n$, as in Eq. 22:

$$E_p^n = (1 + \omega) E_q^n.$$  

(25)

2.5 Actual Evapotranspiration

The calculation of daily actual evapotranspiration, $E_a^n$ (mm d$^{-1}$), is based on the daily integration of the actual evapotranspiration rate, $E_a$ (mm h$^{-1}$), which may be defined as the minimum of the evaporative supply and demand rates (Federer, 1982):

$$E_a = \min(S_w, D_p),$$

(26)

where $S_w$ (mm h$^{-1}$) is the evaporative supply rate, defined in Eq. 21, and $D_p$ (mm h$^{-1}$) is the evaporative demand rate, defined in Eq. 22.

The analytical solution to $E_a^n$ may be expressed analogous to the methodology used for solving $H_o$ and $H_N$ and is defined as twice the integral of $E_a$ between solar noon and $h_n$, which comprises two curves: $S_w$ for $0 \leq h \leq h_i$ and $D_p$ for $h_i \leq h \leq h_n$, where $h_i$ (rad) is the hour angle corresponding to the intersection of $S_w$ and $D_p$ (i.e., when $S_w = D_p$):

$$E_a^n = 2 \int_{h=0}^{h_n} E_a = 2 \left( \int_{h=0}^{h_i} S_w + \int_{h_i}^{h_n} D_p \right),$$

(27a)

which may be expressed as:

$$E_a^n = \frac{24}{\pi} [S_w h_i + r_x r_w r_v (\sin h_n - \sin h_i) + (r_x r_w r_u - r_x I_{LW}) (h_n - h_i)],$$

(27b)

where $r_x = 3.6 \times 10^6 (1 + \omega) E_{con}$ (mm m$^2$ W$^{-1}$ h$^{-1}$). The intersection hour angle, $h_i$, is defined by setting Eq. 21 equal to Eq. 22 and solving for $h$:

$$h_i = \arccos \left( \frac{S_w}{r_x r_w r_v} + \frac{I_{LW}}{r_w r_u - r_v} \right).$$

(28)

To account for the occurrences when supply is in excess of demand during the entire day, $h_i$ should be limited to zero when $\cos h_i \geq 1$. For occurrences when supply limits demand during the entire day, $h_i$ should be limited to $\pi$ when $\cos h_i \leq -1$. The intersection hour angle, $h_i$, is defined by setting Eq. 21 equal to Eq. 22 and solving for $h$:
Figure 2. Example of half-day evaporative supply and demand curves with time represented along actual evapotranspiration curve between the solar hours of solar noon (i.e., evaporative, e, \( h = 0 \)) and solar midnight (i.e., \( h = \pi \)). The evaporative demand, \( D_p \) (dashed red line), is at a maximum at solar noon and zero at the cross-over hour angle, \( h_n \). The evaporative supply, \( S_w \) (dotted blue line), is constant throughout the day. The point where supply is equal to demand denotes the intersection hour angle, \( h_i \). Actual evapotranspiration (solid gray line) is defined as the minimum of \( S_w \) and \( D_p \) throughout the day.

Figure 2 shows an example of the half-day evaporative supply and demand rate curves. \( D_p \) (dashed red line) is at a maximum at solar noon and decreases down to zero at \( h_n \), while \( S_w \) (dotted blue line) is constant throughout the day. The point where \( S_w \) equals \( D_p \) is denoted by the vertical bar at \( h_i \). \( E_a \) (bold solid gray line), limited by supply during most of the day, follows the \( S_w \) line between solar noon and \( h_i \). During the time between \( h_i \) and \( h_n \), \( E_a \) no longer limited by supply, follows the \( D_p \) curve. After \( h_n \), both \( D_p \) and \( E_a \) are zero. \( E_a \) is represented by twice the area above the zero line (dash-dotted horizontal solid black line), below the bold solid lines of \( E_a \) line, and between the vertical bars of solar noon and \( h_n \).

2.6 Runoff

The calculation of daily runoff, \( RO \), is based on the excess of daily soil moisture without runoff compared to the holding capacity, \( W_m \), and is given by:

\[
RO = \max(0, W_n^* - W_m),
\]

where \( W_n^* \) (mm) is the daily soil moisture without runoff (i.e., Eq. 1 where \( RO = 0 \)).

2.7 Soil Moisture

With analytical expressions for \( C_n \) and \( E_a^* \) and \( RO \) (i.e., Eqns. 18 and 27b, 18, 27b and 29, respectively), \( W_n \) may now be calculated by Eq. 1. Daily soil moisture in excess of the maximum soil moisture capacity, \( W_m \), is assumed to be runoff, \( RO \) (mm), and may be calculated as:

\[
\]
To account for the occurrences when $W_n$ exceeds $W_m$ or when $W_n$ drops below zero, once $W_n$ is calculated, the following limits have to be applied to $W_n$ following the calculation of $RO$ are checked:

$$0 \leq W_n \leq W_m. \quad (30)$$

The calculation of $RO$ in Eq. 29 should prevent $W_n$ from being greater than $W_m$, thus satisfying the upper limit of Eq. 30. The limiting effect of $S_w$ on $E_n^a$, through Eqns. 27 and 28, should, in most cases, prevent $W_n$ from falling below zero and satisfy the lower limit of Eq. 30; however, due to the assumption that $S_w$ is constant throughout the day, there is the possibility that $E_n^a - E_p^a + RO$ may exceed $W_{n-1} + P_n + C_n$, resulting in negative $W_n$. In these rare cases, in order to maintain the mass balance of the bucket model presented in Eq. 1, $E_n^a$ is reduced by an amount equal to the magnitude of the negative soil moisture.

### 3 Bioclimatic Indices

One application of the SPLASH model is estimating the surface fluxes required for the calculation of bioclimatic indices. Typically described at longer time scales (e.g., monthly or annually), the daily SPLASH fluxes can be aggregated to monthly and annual totals:

$$X_{m,a} = \sum_{d=1}^{N_{m,a}} X_d, \quad (31)$$

where $X$ is a model output parameter at a given day ($X_d$), month ($X_m$), or year ($X_a$) and $N$ is the total number of days to sum over for a given month ($N_m$) or a given year ($N_a$).

The following sections describe three common bioclimatic indices.

#### 3.1 Moisture Index

There exists a long history that includes several variants of the moisture index, $MI$, also commonly referred to as the aridity index, $AI$, or moisture ratio, $MR$ (Thornthwaite, 1948; Budyko, 1961). A current definition describes $MI$ as the ratio of annual precipitation to annual potential evapotranspiration (Middleton and Thomas, 1997), given as:

$$MI = \frac{P_a}{E_p^a}, \quad (32)$$

where $P_a$ (mm $a^{-1}$) is the annual precipitation and $E_p^a$ (mm $a^{-1}$) is the annual potential evapotranspiration as calculated by Eq. 31; $P_a$ and $E_p^a$ may be substituted with their multi-year means (i.e., $\bar{P}_a$ and $\bar{E}_p^a$) if available. Values less than one are indicative of annual moisture deficit.

Model results and input data of daily (a) fraction of bright sunshine hours, $S_f$ (CRU TS); (b) positive net surface radiation, $H$; (c) condensation, $C$; (d) precipitation, $P_n$ (WATCH); (e) soil moisture, $W_n$; (f) runoff, $RO$; (g) mean air temperature, $T_{air}$ (WATCH); and (h) potential (solid line) and actual (dashed) evapotranspiration, $E_p^a$ and $E_{n,a}$, respectively. Days of the year are represented along the x-axis. Data are for one year (2000 CE) in San Francisco, United States.
3.2 Climatic Water Deficit

The climatic water deficit, $\Delta E$, defined as the difference between the evaporative demand (i.e., potential evapotranspiration) and the actual evapotranspiration, has been shown to be a biologically meaningful measure of climate as it pertains to both the magnitude and length of drought stress experienced by plants (Stephenson, 1998). At the monthly timescale, this index is calculated as:

$$\Delta E_m = E^p_m - E^a_m,$$

(33)

where $\Delta E_m$ (mm mo$^{-1}$) is the monthly climatic water deficit, $E^p_m$ (mm mo$^{-1}$) is the monthly potential evapotranspiration and $E^a_m$ (mm mo$^{-1}$) is the monthly actual evapotranspiration. $E^p_m$ and $E^a_m$ are the monthly totals of $E^n_p$ and $E^n_a$, respectively, calculated by Eq. 31. Values of $\Delta E$ may also be computed at the annual timescale.

3.3 Priestley-Taylor Coefficient

The Priestley-Taylor coefficient, $\alpha$, is the ratio of actual evapotranspiration to equilibrium evapotranspiration, which represents the fraction of plant-available surface moisture (Priestley and Taylor, 1972; Sykes et al., 1996; Gallego-Sala et al., 2010). At the monthly timescale, this is defined as:

$$\alpha_m = \frac{E^a_m}{E^p_m},$$

(34)

where $\alpha_m$ is the monthly Priestley-Taylor coefficient, $E^a_m$ is the monthly actual evapotranspiration and $E^p_m$ (mm mo$^{-1}$) is the monthly equilibrium evapotranspiration. Values of $\alpha$ may also be computed at the annual timescale.

4 Results

The methodology described in Sect. 2 was translated into computer application code (C++, FORTRAN, Python and R). The following sections describe the year-long SPLASH simulation results (2000 CE) at the local and global scales along with comparisons with other model results.

4.1 Local Temporal Trends and Bioclimatic Indices

Data were assembled for one year (2000 CE) including daily WATCH precipitation and air temperature and monthly WATCH Forcing Data methodology applied to the ERA-Interim, WFDEI first release, 2012) and monthly cloudiness fraction (CRU TS3.21 cloudiness fraction). Daily WATCH meteorology was chosen because of its availability for this time period and for illustrating daily model results. At each time step, data were extracted from a single 0.5° × 0.5° pixel above San Francisco, United States (i.e., 37.75° N, 122.25° W). The mean daily air temperature was converted from K to °C and the mean daily precipitation was converted from kg m$^{-2}$ s$^{-1}$ to mm d$^{-1}$ assuming a constant density of water (i.e., $\rho_w = 1 \times 10^3$ kg m$^{-3}$). Fractional sunshine hours were assumed equal to the one-complement of cloudiness fraction and were assumed constant over each month. Figures 3a, 3d, and 3g show the experimental data for $S_f$, $P_n$, and $T_{air}$, respectively.
Approximate values were given for the latitude, $37.7^\circ$ (0.658 rad), and elevation above mean sea level, 142 m, and model. Orbital parameters (for paleoclimatology studies) were assumed constant and calculated for the 2000 CE epoch based on the methods of Berger (1978). Model constants were assigned as per Table 2. The daily soil moisture was initialized at zero and allowed to stabilize, which occurred after just two year-long model iterations. After the second iteration, the daily and monthly results showed no appreciable change and are presented in Figs. 3 and 4. To accompany the daily SPLASH results, in Fig. 3, shown in red, are daily station meteorology and surface fluxes based on the three-layer Variable Infiltration Capacity (VIC) model, extracted from the $1/16^\circ$ pixel centered over San Francisco in the datasets provided in Livneh et al. (2015).

Figure The solid black line in Fig. 3b shows the SPLASH $H_N$ curve (MJ m$^{-2}$), which has the characteristic bell-shaped curve of radiation in the northern hemisphere. The slight jumps between months are due to the irregular jumps in the $S_f$ data (i.e., Fig. 3a). The net radiation flux from Livneh et al. (2015), converted from units of W m$^{-2}$ to MJ m$^{-2}$, follows closely in magnitude to the SPLASH results, albeit with higher inter-daily variability. The SPLASH $H_N$ has a slight negative skew, which is due to the skew of $S_f$.

Figure 3c shows the results for $C_n$, which is based on Eq. 18 as a function of $H_N^*$ and also displays the monthly jumps due to $S_f$. Additional fluctuations in $C_n$ due to the air temperature displays inter-daily variability due its dependence on $T_{air}$ (i.e., Fig. 3g) can also be seen as influenced by the temperature dependency of $C_n$ and, to a lesser degree, monthly variability due to its dependence on $S_f$. The magnitude of $C_n$ varies over the year between 0.45 and 0.8 and contributes about 236 mm annually, which is small when compared to the magnitude of daily rainfall occurrences during the winter months that, in some instances, exceeds 20 mm d$^{-1}$ as shown in Fig. 3d, about 40% of the annual rainfall.

Daily soil moisture, $W_n$, is shown in Fig. 3e. The heavy rains at the beginning of the year (as shown in Fig. 3d) produced saturated soil conditions (i.e., $W_n = W_m = 150$ mm), which gradually reduced as the rainy season came to an end. Small spikes in the soil moisture are seen during the infrequent rain events throughout the spring and summer when soil moisture was maintained at a low level (< 10 mm). At the end of October, when the rains began again, soil moisture began to be replenished. While $W_n$ was saturated, runoff was produced due to the excess in rainfall (shown in begins to replenish. Compared to the VIC soil moisture, the SPLASH soil moisture is consistently between the top two layers (red solid and dashed lines) except during the rainy season where the SPLASH soil moisture magnitude is higher, between the second and third VIC layers (red dashed and dotted lines).

In Fig. 3f, the SPLASH runoff is produced during the rainy season when $W_n$ is saturated. The magnitude of runoff from SPLASH is higher compared to the VIC runoff during the rainy season, which may, in part, be due to the allowance of subsurface flow deep layer drainage in the VIC model; otherwise, the two model results are comparable for the remainder of the year.

Figure 3h shows the overlay of the SPLASH $E_n$ (black solid line) and $E_n^*$ (black dashed line). During the winter and early spring when $D_p$ is relatively low and $S_n$ is non-limiting due to the high soil moisture conditions, $E_n^*$ is shown following the $E_n$ curve. As $D_p$ continued to increase into the summer, $E_n^*$ falls below the $E_n$ curve due to the depletion of soil moisture. The small spikes in soil moisture from rainfall events throughout the late spring and summer can be
seen translated into the $E'_m$ curve. When the rains begin again in the autumn, replenishing soil moisture, replenishing the soil moisture, and $D_p$ had decreased due to the seasonal change in radiation, $E'_m$ is once again shown following the $E'_m$ curve. The same trend shown in Fig. 3h can be seen at the monthly time scale in Fig. 4a, where $E'_m$ (dashed line) is shown following $E'_m$ (solid line) during the first three months, then drops below for the following seven months, and for the last month, once again is following the $E'_m$ curve. The difference between $E'_m$ and $E''_m$ is the climatic water deficit (i.e., Eq. 33), which is shown in Fig. 4b, which highlights the months when supply was limited.

Figure 4c shows the comparison between $E''_m$ (i.e., $E''_m$ with zero entrainment) and $E'_m$. The ratio of $E'_m$ to $E''_m$ is the Priestley-Taylor coefficient, $\alpha_m$ (i.e., Eq. 34). Due to the entrainment factor, $\alpha_m$ may vary between zero (i.e., no moisture) and $1 + \omega$ (i.e., unlimited moisture). During the months when supply is not limiting and $\Delta E_m$ is zero, $\alpha_m$ is at a maximum, as shown in Fig. 4d. Similarly, for the months when supply is limiting and $\Delta E_m$ is positive — a reduction correspond to a dip in $\alpha_m$ occurs.

At the annual timescale, $\Delta E_m$ is 660 mm, which is greater than the annual precipitation (i.e., 620 mm). The annual moisture index (i.e., $MI = 0.493$) and Priestley-Taylor coefficient (i.e., $\alpha_m = 0.598$) are both less than one. These three bioclimatic indices concur that the year was water stressed, which is unsurprising given the frequent occurrence of summer droughts in the western United States.

4.2 Global Simulation of Spatiotemporal Trends

For the global simulation, 0.5° × 0.5° CRU TS3.23 data were assembled for one year (2000 CE), including monthly precipitation (mm mo⁻¹), monthly mean daily air temperature (°C), and monthly cloudiness fraction. Monthly precipitation was converted to daily precipitation by dividing the rainfall equally amongst the days in the month. Fractional sunshine hours were calculated based on the one-complement of cloudiness fraction and assumed constant over the month. Mean daily air temperature was also assumed constant over each day of the month. Half-degree land-surface elevation (m above mean sea level) was provided by CRU TS3.22 (Harris et al., 2014). Once again, orbital parameters were assumed constant over the year and calculated for the 2000 CE epoch based on the methods of Berger (1978) and model constants were assigned as per Table 2.

The SPLASH simulations were driven by the data described above, one pixel at a time, starting each pixel with an empty bucket and terminating when a steady-state of soil moisture was reached between the beginning and the end of the year. Following the steady-state spin-up, the model was driven once again to produce daily simulations of net radiation and soil moisture.

Figure 5b and 5d show the SPLASH results of the mean daily net surface radiation flux (MJ m⁻²) for the months of June and December, respectively. For comparison, the Clouds and the Earth’s Radiant Energy System (CERES) Energy Balanced and Filled (EBAF) average all-sky surface net total flux for June and December 2000 are plotted in Fig. 5a and 5c, respectively. The CERES EBAF net downward radiative flux was converted from W m⁻² to MJ m⁻².

Overall, the SPLASH model produces a reasonable simulation of the latitudinal gradients and seasonal shifts of net surface radiation flux. Locations where the well-watered constant surface albedo assumption fails (e.g., deserts, tundra, and ice sheets),
the SPLASH model simulations are shown to overestimate the CERES EBAF surface radiative flux (especially seen over northern Africa in Fig. 5a and 5b).

Figure 6b and 6d show the SPLASH results of the mean daily relative soil moisture (%) for the months of June and December, respectively. An ice sheet was imposed over Greenland. For comparison, the National Center for Environmental Prediction (NCEP) Climate Prediction Center (CPC) Version 2 mean soil moisture (van den Dool et al., 2003) for June and December 2000 are plotted in Fig. 6a and 6c, respectively. The relative soil moisture in both datasets is computed as the ratio of mm of soil moisture over the total bucket size (i.e., 760 mm in NCEP CPC and 150 mm in SPLASH).

Overall, the SPLASH model simulates soil moisture patterns similar to the NCEP CPC model results. Unlike the NCEP CPC soil moisture, the SPLASH model is shown with a relatively full bucket across wet regions. The lower relative fullness of the NCEP CPC bucket may be contributed to its significantly larger bucket size. Despite the differing magnitudes of soil moisture, the spatial distributions of soil moisture show consistently drier regions in both simulations at both time periods, especially across mid northern latitudes (e.g., eastern North America, northern Africa, and central Asia). Seasonal shifts in soil moisture from June to December are also consistently shown (e.g., southern transition in Africa, eastern transition in South America and northern transition in Australia). There are discrepancies in the spatiotemporal distribution of soil moisture across the high latitude regions (especially Russia). The predominantly saturated conditions in the SPLASH simulations across Russia for December (Fig. 6d) may actually be representative of an increasing snow pack, which could account for these differences.

5 Discussion

The results presented in Sect. 4 are intended to illustrate the dynamic changes in each variable from wet to dry patterns and trends in the SPLASH model outputs across regions and seasons for a single year under steady-state conditions. It should be noted that this work intends only to facilitate the development and application of this model and does not purport to be an in-depth analysis of processes. The SPLASH model results are promising despite the model’s simplifications and limited climatic drivers. At the local scale, the comparison between SPLASH and the VIC model (i.e., Fig. 3b, 3e, 3f, and 3h) shows relatively good agreement in regards to the timing of events and their magnitude (except for runoff). Furthermore, at the global scale, the SPLASH model reasonably captures the latitudinal gradation of net surface radiation flux (where surface emission assumptions are valid) compared to the CERES EBAF results (i.e., Fig. 5) and produces similar spatial patterns of soil moisture, albeit at different magnitudes, compared to the NCEP CPC soil moisture results (i.e., Fig. 6).

While the methodology presented in Sect. 2 makes numerous assumptions and simplifications (e.g., saturation-excess runoff generation, invariant soil properties, physically predictable vegetation function, no infiltration excess runoff, and constant global parameterization), it provides a simple and robust framework for the estimation of radiation components, evapotranspiration, and plant-available moisture requiring only standard meteorological measurements as input. Under steady-state conditions, the SPLASH model preserves the water balance, such that: The SPLASH model currently only produces saturation excess runoff. For more realistic runoff generation, other water balance models allow runoff to occur when the bucket is less than full, for example the empirical relationship of runoff to the weighted relative soil moisture in the simple water balance
model (Orth et al., 2013). Regarding the bucket size, in principle, \( W_{m} \) in Eq. 21 could be formulated as a property of soil type (as was done, for example, in the original BIOME model), there are some objections to doing so. While \( W_{m} \) has a standard definition in agronomy (i.e., the difference between field capacity and wilting point), the wilting point in reality depends on plant properties. Also, the effective ‘bucket size’ depends on rooting behavior, which is highly adaptable to the soil wetness profile. The absolute value of daily soil moisture will be influenced by the bucket size (as shown in Fig. 6) and can have an impact on the local hydrology (e.g., the difference in runoff magnitude in Fig. 3f); however, plant-available moisture indexes, such as \( \alpha \) (i.e., the ratio of supply-limited to non-supply-limited evapotranspiration), have commonly been found to be relatively insensitive to the bucket size. Regarding localized effects, the standard values presented in Table 2 are representative of reasonable global means; however, it is recommended that local parameterization (e.g., shortwave albedo) be used if and when data are available.

Over the years, a common misconception has developed regarding the calculation of daily actual evapotranspiration (as defined by Federer, 1982), whereby the integration of Eq. 26 is mistakenly interpreted as:

\[
E_n^a = \min (S, D),
\]

where \( D \) (mm d\(^{-1}\)) is the total daily demand, given by Eq. 25, and \( S \) (mm d\(^{-1}\)) is the total daily supply over the hours of positive net radiation, which may be given by:

\[
S = \int_{h_n}^{\infty} S_w + \int_{-h_n}^{0} S_w = \frac{24}{\pi} h_n S_w,
\]

where \( h_n \) is the net radiation cross-over angle, given by Eq. 15, and the constant coefficient converts the units of radians to hours. As shown in Fig. 2, \( E_n^a \) is a piecewise function consisting of two curves overlaid throughout the course of a single day that must be accounted for simultaneously; however, even in some recent model developments, \( E_n^a \) is calculated using Eq. 35, including the equilibrium terrestrial biosphere models BIOME3 and BIOME4 (Haxeltine and Prentice, 1996; Kaplan, 2001) and the Lund-Potsdam-Jena Dynamic Global Vegetation Model (Sitch et al., 2003). Only under specific circumstances will Eq. 35 produce correct results. It is the intention of this work to provide a simple analytical solution that correctly accounts for the integration of Eq. 26, which has been provided in the form of Eq. 27b.

**Code Availability**

The code, in four programming languages (C++, FORTRAN, Python, and R), is available on an online repository under the GNU Lesser General Public License (https://bitbucket.org/labprentice/splash). The repository includes the present release (v1.0) and working development of the code (with Makefiles where appropriate), example data, and the user manual. All four versions of the code underwent and passed a set of consistency checks to ensure similar results were produced under the same input conditions. The following describes the requirements for compiling and executing SPLASH v.1.0.

For the C++ version, the code was successfully compiled and executed using the GNU C++ compiler (g++ v.4.8.2) provided by the GNU Compiler Collection (Free Software Foundation, Inc., 2016). It utilizes the C numerics library (cmath),
input/output operations library (cstdio), and the standard general utilities library (cstdlib) and references the vector container and string type.

For the FORTRAN version, the code was successfully compiled and executed using the PGI Fortran compiler (pgf95 v.16.1-0) provided by The Portland Group - PGI Compilers and Groups (NVIDIA Corporation, 2016) and the GNU Fortran compiler (gfortran v.4.8.4) provided by the GNU Compiler Collection (Free Software Foundation, Inc., 2016).

For the Python version, the code was successfully compiled and executed using Python 2.7 and Python 3.5 interpreters (Python Software Foundation, 2016). It requires the installation of third-party packages, including NumPy (v.1.10.4 by NumPy Developers, 2016) and SciPy (v.0.17.0 by SciPy Developers, 2016) and utilizes the basic date and time types (datetime), logging facility (logging), Unix-style pathname pattern extension (glob), and miscellaneous operating system interfaces (os) modules.

For the R version, the code was successfully compiled and executed using R-3.2.3 “Wooden Christmas-Tree” (The R Foundation for Statistical Computing, 2015).

**Appendix A: Calculating True Longitude**

Berger (1978) presents a method for estimating true longitude, \( \lambda \), for a given day of the year, \( n \), that associates uniform time (i.e., a mean planetary orbit and constant day of the vernal equinox) to Earth’s angular position. The formula is based on classical astronomy and is suitable for calculations in palaeoclimatology. The algorithm begins with the calculation of the mean longitude of the vernal equinox, \( \lambda_{m0} \) (rad), assumed to fall on 21 March:

\[
\lambda_{m0} = 2 \left[ \left( \frac{1}{2} e + \frac{1}{8} e^3 \right) (1 + \beta) \sin \bar{\omega} - \frac{1}{4} e^2 \left( \frac{1}{2} + \beta \right) \sin 2 \bar{\omega} + \frac{1}{8} e^3 \left( \frac{1}{3} + \beta \right) \sin 3 \bar{\omega} \right], \tag{A1}
\]

where \( \beta = \sqrt{1 - e^2} \). The mean longitude, \( \lambda_m \) (rad), is then calculated for a given day based on a daily increment with respect to the day of the vernal equinox (i.e., day 80):

\[
\lambda_m = \lambda_{m0} + 2\pi (n - 80) N_a^{-1}, \tag{A2}
\]

where \( N_a \) is total number of days in the year. The mean anomaly, \( \nu_m \) (rad), is calculated based on the equality presented in Eq. 6:

\[
\nu_m = \lambda_m - \bar{\omega}, \tag{A3}
\]

which is then used to determine the true anomaly by:

\[
\nu = \nu_m + \left( 2e - \frac{1}{4} e^3 \right) \sin \nu_m + \frac{5}{4} e^2 \sin 2\nu_m + \frac{13}{12} e^3 \sin 3\nu_m, \tag{A4}
\]

and is converted back to true longitude by:

\[
\lambda = \nu + \bar{\omega}. \tag{A5}
\]

The resulting \( \lambda \) should be constrained to an angle within a single orbit (i.e., \( 0 \leq \lambda \leq 2\pi \)).
Appendix B: Calculating Temperature and Pressure Dependencies

The four variables used to calculate the water-to-energy conversion factor, $E_{con}$, given in Eq. 19 have temperature and/or pressure dependencies that may be solved using the equations presented here.

The temperature-dependent equation for the slope of the saturation vapor pressure-temperature curve, $s$, can be expressed as (Allen et al., 1998):

$$s = \frac{2.503 \times 10^6 \exp \left( \frac{17.27 T_{air}}{T_{air} + 237.3} \right)}{(T_{air} + 237.3)^2},$$

(B1)

where $s$ ranges from about 11 to 393 Pa K$^{-1}$ for $T_{air}$ between $-20$ and 40 °C. Please be aware of the typographical error in this formula as presented in Eq. 7 of Gallego-Sala et al. (2010) where 237.3 is misrepresented as 273.3.

The temperature-dependent equation for the latent heat of vaporization, $L_v$, may be expressed as (Henderson-Sellers, 1984):

$$L_v = 1.91846 \times 10^6 \left[ \frac{T_{air} + 273.15}{(T_{air} + 273.15) - 33.91} \right]^2,$$

(B2)

where $L_v$ ranges from about $2.558 \times 10^6$ to $2.413 \times 10^6$ J K$^{-1}$ for $T_{air}$ between $-20$ and 40 °C.

The temperature and pressure dependence of the density of water, $\rho_w$, may be expressed as (Chen et al., 1977):

$$\rho_w = \rho_o \frac{K_o + C_A P_{atm}^* + C_B P_{atm}^*}{K_o + C_A P_{atm}^* + C_B P_{atm}^* - P_{atm}^*}.$$

(B3)

The equation for $\rho_o$ is based on the work of Kell (1975): 

$$\rho_o = \sum_{i=0}^{8} C_i T_{air}^i.$$ 

(B4)

The equation for $K_o$ is also based on the work of Kell (1975):

$$K_o = \sum_{i=0}^{5} C_i T_{air}^i.$$ 

(B5)

The equations for $C_A$ and $C_B$ are given as (Chen et al., 1977):

$$C_A = \sum_{i=0}^{4} C_i T_{air}^i.$$ 

(B6)

$$C_B = \sum_{i=0}^{4} C_i T_{air}^i.$$ 

(B7)

The coefficients for $T_{air}$ in Eqns. B4 through B7 are given in Table 3.
The temperature and pressure dependence of the psychrometric constant, $\gamma$, may be expressed as (Allen et al., 1998):

$$\gamma = \frac{C_p M_a P_{atm}}{M_v L_v},$$  
(B8)

where $C_p (\text{J kg}^{-1} \text{K}^{-1})$ is the temperature-dependent specific heat capacity of humid air; $M_a$ (kg mol$^{-1}$) and $M_v$ (kg mol$^{-1}$) are the molecular weights of dry air and water vapor, respectively; $L_v$ (J kg$^{-1}$) is the latent heat of vaporization of water; and $P_{atm}$ (Pa) is the atmospheric pressure. Constants for $M_a$ and $M_v$ are given in Table 2. The temperature dependence of $C_p$ may be assumed negligible (e.g., $C_p = 1.013 \times 10^3 \text{ J kg}^{-1} \text{ K}^{-1}$) or calculated by (Tsilingiris, 2008):

$$C_p = \sum_{i=0}^{5} C_i T_{air}^i,$$  
(B9)

where the coefficients of $T_{air}$ are given in Table 3.

Author contributions. I. C. Prentice, M. T. Sykes, and W. Cramer developed the original model theory and methods. A. V. Gallego-Sala, B. J. Evans, H. Wang, and T. W. Davis contributed to model improvements. R. T. Thomas, R. J. Whitley, B. D. Stocker, and T. W. Davis transcribed the new model code and ran simulations. The manuscript was prepared with contributions from all authors.

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...
References


Table 1. Nomenclature.

<table>
<thead>
<tr>
<th>Instantaneous</th>
<th>Daily</th>
<th>Monthly</th>
</tr>
</thead>
<tbody>
<tr>
<td>( S_w )</td>
<td>( W_n ) soil moisture, mm</td>
<td>( E'_m ) equilibrium evapotranspiration, mm mo(^{-1})</td>
</tr>
<tr>
<td>( D_p )</td>
<td>( P_n ) precipitation, mm d(^{-1})</td>
<td>( E''_m ) potential evapotranspiration, mm mo(^{-1})</td>
</tr>
<tr>
<td>( E_q )</td>
<td>( C_n ) condensation, mm d(^{-1})</td>
<td>( E'_m ) actual evapotranspiration, mm mo(^{-1})</td>
</tr>
<tr>
<td>( E_p )</td>
<td>( RO ) runoff, mm</td>
<td>( H_o ) solar irradiation, J m(^{-2}) d(^{-1})</td>
</tr>
<tr>
<td>( E_a )</td>
<td>( E_n ) equilibrium evapotranspiration, mm d(^{-1})</td>
<td>( H_N ) net surface radiation, J m(^{-2}) d(^{-1})</td>
</tr>
<tr>
<td>( I_o )</td>
<td>( H_{SN}^+ ) nighttime positive net surface radiation, J m(^{-2}) d(^{-1})</td>
<td></td>
</tr>
<tr>
<td>( I_N )</td>
<td>( H_{SN}^- ) negative net surface radiation, J m(^{-2}) d(^{-1})</td>
<td></td>
</tr>
<tr>
<td>( I_{SW} )</td>
<td>( Q_n ) photosynthetically active radiation, mol m(^{-2}) d(^{-1})</td>
<td></td>
</tr>
<tr>
<td>( I_{LW} )</td>
<td>( S_f ) fraction of bright sunshine hours, unitless</td>
<td></td>
</tr>
<tr>
<td>( E'_m )</td>
<td>( T_{air} ) mean air temperature, °C</td>
<td></td>
</tr>
<tr>
<td>( E''_m )</td>
<td>( \Delta E_m ) climatic water deficit, mm mo(^{-1})</td>
<td></td>
</tr>
<tr>
<td>( E'_m )</td>
<td>( \alpha_m ) Priestley-Taylor coefficient, unitless</td>
<td></td>
</tr>
<tr>
<td>( E''_m )</td>
<td></td>
<td></td>
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25
### Table 1 (continued). Nomenclature.

<table>
<thead>
<tr>
<th>Symbol</th>
<th>Description</th>
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</thead>
<tbody>
<tr>
<td>$\cos \theta_z$</td>
<td>inclination factor, unitless</td>
</tr>
<tr>
<td>$\delta$</td>
<td>declination angle, rad</td>
</tr>
<tr>
<td>$d_r$</td>
<td>distance factor, unitless</td>
</tr>
<tr>
<td>$\varepsilon$</td>
<td>obliquity, rad</td>
</tr>
<tr>
<td>$e$</td>
<td>eccentricity, unitless</td>
</tr>
<tr>
<td>$E_{\text{con}}$</td>
<td>water to energy conversion factor, m$^3$ J$^{-1}$</td>
</tr>
<tr>
<td>$\gamma$</td>
<td>psychrometric constant, Pa K$^{-1}$</td>
</tr>
<tr>
<td>$h$</td>
<td>hour angle, rad</td>
</tr>
<tr>
<td>$h_i$</td>
<td>intersection of evaporative rates hour angle, rad</td>
</tr>
<tr>
<td>$h_n$</td>
<td>net radiation crossover hour angle, rad</td>
</tr>
<tr>
<td>$h_s$</td>
<td>sunset hour angle, rad</td>
</tr>
<tr>
<td>$i$</td>
<td>day of month (1–31)</td>
</tr>
<tr>
<td>$\lambda$</td>
<td>true longitude, rad</td>
</tr>
<tr>
<td>$L_v$</td>
<td>latent heat of vaporization of water, J kg$^{-1}$</td>
</tr>
<tr>
<td>$\nu$</td>
<td>true anomaly, rad</td>
</tr>
<tr>
<td>$n$</td>
<td>day of year (i.e., 1–365)</td>
</tr>
<tr>
<td>$N_a$</td>
<td>total number of days in a year (e.g., 365)</td>
</tr>
<tr>
<td>$N_m$</td>
<td>total number of days in a given month (e.g., 31)</td>
</tr>
<tr>
<td>$\tilde{\omega}$</td>
<td>longitude of perihelion, rad</td>
</tr>
<tr>
<td>$\phi$</td>
<td>latitude, rad</td>
</tr>
<tr>
<td>$P_{\text{atm}}$</td>
<td>atmospheric pressure, Pa</td>
</tr>
<tr>
<td>$\rho_w$</td>
<td>density of water, kg m$^{-3}$</td>
</tr>
<tr>
<td>$r_u$</td>
<td>$\sin \delta \sin \phi$, unitless</td>
</tr>
<tr>
<td>$r_v$</td>
<td>$\cos \delta \cos \phi$, unitless</td>
</tr>
<tr>
<td>$r_w$</td>
<td>$(1 - \beta_{sw}) \tau I_{sc} d_r$, W m$^{-2}$</td>
</tr>
<tr>
<td>$r_x$</td>
<td>$3.6 \times 10^6 (1 + \omega) E_{\text{con}}$, mm m$^2$ W$^{-1}$ h$^{-1}$</td>
</tr>
<tr>
<td>$s$</td>
<td>slope of saturated vapor pressure-temperature curve, Pa K$^{-1}$</td>
</tr>
<tr>
<td>$\tau$</td>
<td>transmittivity, unitless</td>
</tr>
<tr>
<td>$\tau_o$</td>
<td>transmittivity at mean sea level, unitless</td>
</tr>
<tr>
<td>$z$</td>
<td>elevation above mean sea level, m</td>
</tr>
</tbody>
</table>
Figure 3. Daily (a) fraction of bright sunshine hours, $S_f$ (based on CRU TS3.21); (b) SPLASH (black) and VIC (red) modeled net surface radiation, $H_N$; (c) SPLASH modeled condensation, $C_n$; (d) precipitation, $P_n$ (based on WFDEI 2012 in black and station measurements in red); (e) SPLASH (black) and VIC three-layer (layer 1 in solid red, layer 2 in dashed red, and layer 3 in dotted red) modeled soil moisture, $W_n$; (f) SPLASH (black) and VIC (red) modeled runoff, $RO$; (g) mean near-surface air temperature, $T_{air}$ (based on WFDEI, 2012 in black) and station max and min near-surface air temperature (red colored region); and (h) SPLASH and VIC potential (black and red solid lines, respectively) and SPLASH actual (black dashed) evapotranspiration, $E_{ap}^P$ and $E_{ap}^A$, respectively. Data plotted in red are based on Livneh et al. (2015). Days of the year are represented along the x-axis. Data are for one year (2000 CE) extracted from a pixel centered over San Francisco, United States.
Figure 4. Model results of monthly (a) potential (solid line) and actual (dashed line) evapotranspiration, \( E_{pm} \) and \( E_{am} \), respectively; (b) climatic water deficit, \( \Delta E_m \); (c) equilibrium (solid) and actual (dashed line) evapotranspiration, \( E_{qm} \) and \( E_{am} \), respectively; and (d) monthly Priestley-Taylor coefficient, \( \alpha_m \). Months of the year are represented along the x-axis. Results are of one year (2000 CE) for San Francisco, United States.
Figure 5. Mean net downward surface radiation flux, \( \text{MJ m}^{-2} \), for (a) June 2000 (CERES EBAF); (b) June 2000 (SPLASH); (c) December 2000 (CERES EBAF); December 2000 (SPLASH).

Figure 6. Mean relative soil moisture, \( \% \), for (a) June 2000 (NCEP CPC); (b) June 2000 (SPLASH); (c) December 2000 (NCEP CPC); (d) December 2000 (SPLASH). The relative soil moisture is based on the total bucket size (i.e., 760 mm for NCEP CPC and 150 mm for SPLASH).
### Table 2. Constants and Standard Values.

<table>
<thead>
<tr>
<th>Variable</th>
<th>Units</th>
<th>Description</th>
</tr>
</thead>
<tbody>
<tr>
<td>A</td>
<td>107 °C</td>
<td>empirical constant, Eq. 13 (Monteith and Unsworth, 1990)</td>
</tr>
<tr>
<td>$\beta_{sw}$</td>
<td>0.17</td>
<td>shortwave albedo, Eq. 10 (Federer, 1968)</td>
</tr>
<tr>
<td>$\beta_{vis}$</td>
<td>0.03</td>
<td>visible light albedo, Eq. 17 (Sellers, 1985)</td>
</tr>
<tr>
<td>b</td>
<td>0.20</td>
<td>empirical constant, Eq. 13 (Linacre, 1968)</td>
</tr>
<tr>
<td>c</td>
<td>0.25</td>
<td>cloudy transmittivity, Eq. 12 (Linacre, 1968)</td>
</tr>
<tr>
<td>d</td>
<td>0.50</td>
<td>angular coefficient of transmittivity, Eq. 12 (Linacre, 1968)</td>
</tr>
<tr>
<td>fFEC</td>
<td>2.04 μmol J$^{-1}$</td>
<td>flux-to-energy conversion, Eq. 17 (Meek et al., 1984)</td>
</tr>
<tr>
<td>g</td>
<td>9.80665 m s$^{-2}$</td>
<td>standard gravity, Eq. 20 (Allen, 1973)</td>
</tr>
<tr>
<td>$I_{sc}$</td>
<td>1360.8 W m$^{-2}$</td>
<td>solar constant, Eq. 2 (Kopp and Lean, 2011)</td>
</tr>
<tr>
<td>L</td>
<td>0.0065 K m$^{-1}$</td>
<td>mean adiabatic lapse rate, Eq. 20 (Allen, 1973)</td>
</tr>
<tr>
<td>$M_a$</td>
<td>0.028963 kg mol$^{-1}$</td>
<td>molecular weight of dry air, Eq. 20 (Tsilingiris, 2008)</td>
</tr>
<tr>
<td>$M_v$</td>
<td>0.01802 kg mol$^{-1}$</td>
<td>molecular weight of water vapor, Eq. B8 (Tsilingiris, 2008)</td>
</tr>
<tr>
<td>$\omega$</td>
<td>0.26</td>
<td>entrainment factor, Eq. 22 (Priestley and Taylor, 1972)</td>
</tr>
<tr>
<td>$P_o$</td>
<td>101325 Pa</td>
<td>standard sea-level pressure, Eq. 20 (Allen, 1973)</td>
</tr>
<tr>
<td>R</td>
<td>8.31447 J mol$^{-1}$ K$^{-1}$</td>
<td>universal gas constant, Eq. 20 (Moldover et al., 1988)</td>
</tr>
<tr>
<td>$S_c$</td>
<td>1.05 mm h$^{-1}$</td>
<td>supply rate constant, Eq. 21 (Federer, 1982)</td>
</tr>
<tr>
<td>$T_o$</td>
<td>288.15 K</td>
<td>base temperature, Eq. 20 (Berberan-Santos et al., 1997)</td>
</tr>
<tr>
<td>$W_m$</td>
<td>150 mm</td>
<td>soil moisture capacity, Eq. 21 (Cramer and Prentice, 1988)</td>
</tr>
</tbody>
</table>

### Table 3. Coefficients of $T_{air}$.

<table>
<thead>
<tr>
<th>$\rho_o$ (kg m$^{-3}$)</th>
<th>$K_o$ (bar)</th>
<th>$C_A$ (unitless)</th>
<th>$C_B$ (bar$^{-1}$)</th>
<th>$C_p$ (J kg$^{-1}$ K$^{-1}$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>$C_0$</td>
<td>9.998395 × 10$^2$</td>
<td>1.96520 × 10$^4$</td>
<td>3.26138</td>
<td>7.2061 × 10$^{-5}$</td>
</tr>
<tr>
<td>$C_1$</td>
<td>6.78826 × 10$^{-2}$</td>
<td>1.48183 × 10$^2$</td>
<td>5.223 × 10$^{-4}$</td>
<td>5.8948 × 10$^{-6}$</td>
</tr>
<tr>
<td>$C_2$</td>
<td>-9.08659 × 10$^{-3}$</td>
<td>-2.29995</td>
<td>1.324 × 10$^{-4}$</td>
<td>8.6990 × 10$^{-8}$</td>
</tr>
<tr>
<td>$C_3$</td>
<td>1.02213 × 10$^{-4}$</td>
<td>1.28100 × 10$^{-2}$</td>
<td>7.655 × 10$^{-7}$</td>
<td>-1.0100 × 10$^{-9}$</td>
</tr>
<tr>
<td>$C_4$</td>
<td>-1.35439 × 10$^{-6}$</td>
<td>-4.91564 × 10$^{-5}$</td>
<td>8.584 × 10$^{-10}$</td>
<td>4.3220 × 10$^{-12}$</td>
</tr>
<tr>
<td>$C_5$</td>
<td>1.47115 × 10$^{-8}$</td>
<td>1.03553 × 10$^{-7}$</td>
<td>—</td>
<td>—</td>
</tr>
<tr>
<td>$C_6$</td>
<td>-1.11663 × 10$^{-10}$</td>
<td>—</td>
<td>—</td>
<td>—</td>
</tr>
<tr>
<td>$C_7$</td>
<td>5.04407 × 10$^{-13}$</td>
<td>—</td>
<td>—</td>
<td>—</td>
</tr>
<tr>
<td>$C_8$</td>
<td>-1.00659 × 10$^{-15}$</td>
<td>—</td>
<td>—</td>
<td>—</td>
</tr>
</tbody>
</table>