We would like to thank the two anonymous Referees for their comments on the manuscript. In addition, we would like to thank Dr Sylwester Arabas for his review and suggestions. All comments have been addressed in the revised version of the manuscript. Below we are submitting our replies to the Reviewers, the list of additional changes and the revised version of the manuscript.

Review of libcloudph++ 1.1: aqueous phase chemistry extension of the Lagrangian cloud microphysics scheme by Jaruga and Pawlowska (GMD-2018-96), Anonymous Referee #1

- P. 8, l. 2 - p. 9, l. 7, Fig. 1: The liquid water content is one of the most important parameters to steer aqueous sulfur chemistry (and probably in cloud physics in general). Therefore, it is necessary that the same liquid water content is reproduced by the presented model to validate it with the intercomparison results of Kreidenweis et al. (2003), in which seven other models are capable of reproducing the same liquid water content. The authors state that the reason for this discrepancy results from the unspecified pressure profile in the study of Kreidenweis et al. (2003). However, there might be another potential source for the deviation: it could be the slightly too high initial temperature in the current study, which is stated as 285.5 K (Tab. 1 of the manuscript) but 285.2 K in Kreidenweis et al. (2003, their Tab. 3).

Thank you for catching this error. This is just a typo in the manuscript and the actual initial temperature used in the parcel model simulations was 285.2K: see https://github.com/igfwu/parcel/blob/master/chem_conditions.py#L13.

We have tested the dependence of the parcel model on different pressure $p$ and dry air density $\rho_d$ calculations. We repeated the parcel model simulations using three different ways to calculate $p$ and $\rho_d$. In all simulations $p$ is calculated by integrating the hydrostatic equation:

$$ \frac{dp}{dz} = -pg, $$

where $p$ represents pressure, $z$ is the vertical displacement, $\rho$ is the density of air and $g$ is the gravitational acceleration.

- In the first simulation, to integrate Eq. (1) we assume that $p$ is constant and equal $p_0 = 1.15 \text{ kg m}^{-3}$. The assumed density value is taken from Tab 3 in Kreidenweis et al. (2003). At each time level $n$ the pressure is given by

$$ p^n = p^0 - \rho_0gz^n. $$

The dry air density is given by:

$$ \rho_d(p, \theta, r_v) = \frac{p - p_v(p, r_v)}{R_d\theta(p_{1000})\frac{v_g}{R_d}}, $$

where $p_v(p, r_v)$ represents partial pressure of water vapor, $p_{1000}$ stands for pressure equal 1000 $\text{hPa}$ that comes from the definition of potential temperature, $R_d$ is the gas constant for dry air and $c_{pd}$ is the specific heat at constant pressure for dry air. The dry air density at a given time level $n$ is calculated as $\rho_d(p^n, \theta^n, r_v^n)$. This simulation is labeled const_rho

- For the second simulation, we assume that the density is constant at each height level (i.e. piecewise constant profile). As a result, at a given time level $n$ the pressure used to predict $\rho_d$ is defined as

$$ p^n = p^{n-1} - \rho^{n-1}gz^n. $$

Again the dry air density at a given time level $n$ is calculated using Eq.(3) as $\rho_d(p^n, \theta^n, r_v^n)$. This simulation is labeled piecewise_const_rho and was used in the submitted manuscript.

- In the third simulation, we assume that the potential temperature, $\theta = \theta^0$, and water vapor mixing ratio, $r_v = r_v^0$ are constant. At a given time level $n$ $p$ is then defined as

$$ p^n = p_{1000}\left(\frac{p^0_{1000}}{p_{1000}}\right)^{\frac{n}{v_g}} - \frac{R_d\theta(z^n - z^0)}{c_{pd}\theta^0R(r_v^n)}, $$

where $R(r_v^0)$ is the gas constant for moist air. The dry air density at a given time level $n$ is calculated using Eq.(3) as $\rho_d(p^n, \theta^0, r_v^0)$. This method follows the procedure used in the kinematic model, described
in Arabas et al. (2015) and used in the manuscript in the 2-dimensional simulations. This simulation is labeled const_th_rv.

The results of the tests are presented in Fig. 1 and can be compared with Fig. 1 in Kreidenweis et al. (2003). Additionally, the values of parcel acidity, sulfate formation, number of droplets at maximum saturation level and the maximum saturation are reported in Tab. 1. The simulation results are similar for the SO\textsubscript{2} concentration and pH calculations. The const_th_rv simulation is the biggest outlier. The biggest differences can be seen for the liquid water content profile (first panel) of Fig. 1. In general, the const_rho simulation agrees best with the results shown in Fig. 1 in Kreidenweis et al. (2003). In order to better match the results presented in Kreidenweis et al. (2003) the parcel model simulations shown in the manuscript are changed to use the constant density assumption (const_rho). The relevant descriptions and figures are changed.

Figure 1: The parcel model simulations using three different ways to calculate the pressure and dry air density profiles: assuming constant density profile (blue), assuming piecewise constant density profile (red) and assuming constant potential temperature and water vapor mixing ratio profile (green). The panels show the liquid water mixing ratio, the SO\textsubscript{2} concentration and the water volume weighted average pH, respectively.

<table>
<thead>
<tr>
<th></th>
<th>const_rho</th>
<th>piecewise_const_rho</th>
<th>const_th_rv</th>
<th>Kreidenweis et al. (2003)</th>
</tr>
</thead>
<tbody>
<tr>
<td>water volume weighted pH</td>
<td>4.86</td>
<td>4.83</td>
<td>4.79</td>
<td>4.82 - 4.85</td>
</tr>
<tr>
<td>sulfate formation (ppt)</td>
<td>171</td>
<td>168</td>
<td>161</td>
<td>170 - 180</td>
</tr>
<tr>
<td>sulfate formation by H\textsubscript{2}O\textsubscript{2} (ppt)</td>
<td>99</td>
<td>100</td>
<td>104</td>
<td>85 - 105</td>
</tr>
<tr>
<td>sulfate formation by O\textsubscript{3} (ppt)</td>
<td>72</td>
<td>68</td>
<td>57</td>
<td>70 - 85</td>
</tr>
<tr>
<td>number of droplets in cm\textsuperscript{-3}</td>
<td>269</td>
<td>272</td>
<td>265</td>
<td>275 - 358</td>
</tr>
<tr>
<td>maximum supersaturation</td>
<td>0.27</td>
<td>0.27</td>
<td>0.26</td>
<td>0.23 - 0.37</td>
</tr>
</tbody>
</table>

Moreover, it might be helpful to clarify the calculation of the pressure (and density) profile by using equations (the description on p. 8, ll. 21-29 is hard to follow).

The corresponding description was changed to reflect the changed way of calculating pressure and density profile.

Finally, why does the axis stating the height above cloud base start with 100 m but the time above cloud base with 0 s (Fig. 1)?

Thank you, we have corrected the axis labels.

**Introduction:** Many previous models (so-called moving bin or moving grid models) used a Lagrangian framework for the representation of aqueous sulfur chemistry as done in this study. Accordingly, they are based on similar (sometimes identical) model equations as the described...
in the current manuscript. As the LCM, these moving bin models do not suffer from numerical diffusion (e.g., p. 4, ll. 7-8). However, these models are only applicable in simplified parcel or trajectory simulations (e.g., Feingold et al., 1998; Kreidenweis et al., 2003). The authors should make clear that their main advancement is the possibility to use a Lagrangian representation of aqueous sulfur chemistry in multidimensional simulations with a fully coupled dynamics model, which has not been possible before.

Thank you for pointing this out! Using the term "Lagrangian" can be indeed confusing as it might refer to different microphysics schemes: (i) "Lagrangian-in-droplet-radius" - as in the "moving bin" models from Kreidenweis et al. (2003) or (ii) "Lagrangian-in-droplet-radius-and-space" - as it is done in the super-droplet method (Shima et al., 2009) used in this study. To avoid confusion we added a clarification in the introduction. We also decided to stop using the term Lagrangian when referring to our scheme. Instead, following the notation used in the original super-droplet paper by Shima et al. (2009), we will refer to this method as particle-based.

- P. 2, ll. 8-24: Please rearrange/rewrite these two paragraphs. Especially the first paragraph gives no real reasons why aqueous phase oxidation of sulfur is an important chemical reaction, it only states it is a very common reaction. Moreover, there are many interesting points in these two paragraphs but it would be nice to connect them more closely with the presented study/model.

We rearranged the paragraphs to hopefully put more emphasis on the importance of the aqueous phase oxidation of sulfur. Overall, it is worth underlining that the main reason for adding the aqueous phase oxidation of sulfur to the particle-based microphysics scheme was to study its effects on the aerosol size distribution and its potential further effects on cloud droplet size distribution.

Any suggestions what other aqueous-phase chemical reactions would be interesting from the atmospheric chemistry standpoint are greatly appreciated.

- P. 2, ll. 28-29: What about the collection of submicron aerosol by droplets (e.g., Ardon-Dryer et al. (2015), Ladino et al. (2011))? These processes are not considered in the current model but also affect the size of aerosols.

The particle-based method is very well suited for studying such process. The deliquescent aerosol particles and droplets are represented in the same way in the super-droplet method - there is no artificial category of "aerosol particle", "cloud droplet" or "rain drop". The only change needed to resolve the collection of submicron aerosol particles by droplets in the current setup is adding a collision efficiency look-up table based on the works of Ladino et al. (2011) and Ardon-Dryer et al. (2015) that would correctly represent the probability of collision between submicron aerosols and bigger drops.

We have added this process to the list of cloud-related processes that affect the aerosol sizes. We also added information on how it should be included in the particle-based microphysics scheme in Sec. 2.

- P. 3, l. 2: What is meant by an atmospheric aerosol particle in this context? There must certainly be a size restriction to the cited statement since very small aerosols are not cloud-active.

The mentioned study serves as an estimate of the global average recycling by clouds. From our understanding of the derivation of those estimates, the relevant aerosol sizes are between 0.1\(\mu m\) and 1 \(\mu m\).

- P. 3, ll. 5-6: The effect of collision and coalescence on the amount of aerosol inside a cloud droplet (in absence of aqueous sulfur chemistry!) has been considered in an LCM framework before (see Hoffmann, 2017).

We are mentioning this work in Sec. 2.

- P. 4, ll. 11-12: This is a vague statement since it depends on the equations considered and the applied numerical solvers.

We crossed out this sentence.

- P. 5, l. 2: How does the wind speed affect the chemical scheme?

The scheme requires the velocity field (or wind field) to be provided, in order to advect the super-droplets in the computational domain. Because the scheme can be used in models of different dimensionalities, this
velocity field could either be a constant number (as in parcel model simulations in Sec. 4) 2-dimensional prescribed flow field (as in the 2-dimensional simulations in Sec. 5) or a 3-dimensional LES velocity field (hopefully in the future). To clarify the sentence we changed the "wind speed" to "wind field".

- P. 5, ll. 8-13: Does this restriction affect the results? And if yes, how? And what happens to giant and ultra-giant aerosols, which remain less diluted throughout their lifecycle?

The assumption that the solution droplets are diluted only affects the chemical calculations of the presented scheme. The impact of GCCN on rain formation could be studied from the collisions perspective only. No changes to the current version of the scheme would be necessary for such simulations. To study the impact of GCCN on the aqueous phase chemistry another aerosol types should be added first. For example, the sea-salt aerosol particles and chemical processes relevant to them should be included, see the discussion starting in line 30 on page 13.

In the simulations discussed in the manuscript, the assumption that solution droplets are diluted excludes from the aqueous phase simulations deliquescent aerosol particles and rain drops at the last stages of their evaporation. In both cases, excluding aqueous phase oxidation reaction is very beneficial for the condensation/evaporation scheme. Especially during activation of cloud droplets the wet radius of droplets changes very fast. It would be difficult for the condensation scheme to reach convergence during cloud droplet activation if the dry and wet radii of the super-droplets were allowed to change simultaneously.

We tested the dependence of the parcel model simulations for the threshold values of 0.01 moles/liter and 0.1 moles/liter. In both cases we found no significant difference from the 0.02 moles/liter threshold that is used by default. Any further tests should be done in a full 3-dimensional LES setup that would allow for a longer simulation times with a more realistic representation of rain formation and aerosol recycling by clouds. However such simulations are out of the scope of the current paper.

- P. 5, ll. 20-22: In what aspect is the LCM beneficial? Using a low number of super-droplets, the LCM will suffer from the same problems as a moving bin scheme.

We crossed out this statement.

- P. 5, Eq. (1): For a reaction \( ab \rightarrow a + b \) dissociation constant is defined as \( K = [a][b]/[ab] \). Accordingly, Eq. (1) is just a relationship to adjust a measured dissociation constant \( K \) to a given temperature. Denote the meaning of \( T \) and \( R \).

Done!

- P. 6, Eq. (2): I understand that this equation is based on the neutral charge condition. For clarity, it might be necessary to comment on how the expressions for the individual terms are derived. (Or state an appropriate reference.)

We added a reference to the relevant chapter in Seinfeld and Pandis (2016).

- P. 6, ll. 15-16; p. 7, l. 15: How is this small time-step determined? Is there a criterion? How does the user know that an appropriate time-step is chosen?

Unfortunately we don’t have a criterion to determine the necessary time-step. The system is non-linear and the exact criterion depends on the changes in droplet radius due to condensation and collisions as well as the actual chemical composition of the droplets. The processes limiting the time-step length in the simulations presented in the manuscript are the condensation/evaporation and the aqueous phase chemistry. The parcel model simulations from Sec. 4 and the 2-dimensional simulations from Sec. 5 use model time-step equal to 1 second. For the simulations from Sec. 5 additional 10 sub-steps were used per each model time-step for both condensation and aqueous phase chemistry schemes.

The libcloudph++ can be compiled and run in the debugging mode. In this mode all the asserts (for example checks if the trace-gas fields are non-negative) are performed. When using the library in a new modeling setup it is worthwhile to first use it in a debugging mode to check that the simulation is stable and if necessary adjust the model time-step.

- P. 7, ll. 26-28: Why is H2SO4 used instead of SIV to denote sulfuric acid? Please clarify.

We changed it to SVI.
• P. 8, ll. 4-5: What is meant by initial conditions above supersaturation?
  To clarify, we changed it to "initial conditions where supersaturation is present in the environment". In general, such initial conditions are used in the kinematic model setups.

• P. 8, ll. 15-20; p. 12, Fig. 3: How is the aerosol distribution transformed to super-droplets? Is the distribution binned and a super-droplet represents each bin or is a random generator used that follows the given distribution? This is important information to understand the slight fluctuations in the initial spectrum (Fig. 3).
  The dry radius of super-droplets is chosen randomly with a uniform distribution in the logarithm of the radius. The minimal and maximal values of the dry radius are chosen automatically by evaluating the initial size distribution. The criterion is that the super-droplet multiplicity has to be greater or equal to 1. We added the reference to Sec. 5.1.6 in Arabas et al. (2015) where the procedure is discussed in detail.

• P. 10, ll. 9-10: How is it possible to have activated particles at a relative humidity of 100 %?
  To activate a droplet, a supersaturation ($RH > 1$) is necessary.
  Agreed. We crossed out this statement and only report the number of cloud droplets after the maximum super-saturation is reached.

• P. 13, ll. 3-4: The Hall (1980) collection efficiencies are not very suitable for big droplets since they neglect the effect of a coalescence efficiency of less than unity. How are big droplets defined?
  The collision efficiency values used in the simulations presented in the manuscript are shown in Fig. 2. Left panel shows the collision efficiencies based on Hall (1980). They are used for droplets with wet radius greater than 20 $\mu$m. Right panel shows the collision efficiencies based on Pinsky et al. (2008), which are used for collisions between droplets with radius smaller than 20 $\mu$m. For easier comparison, the middle panel of Fig. 2 zooms in on the Hall (1980) collision efficiencies for the small droplets.
  In general, the 2-dimensional simulations presented in Sec. 5 of the manuscript do not produce very big rain drops. At the end of the simulation the concentration of rain drops with radius greater that 100 $\mu$m does not exceed 0.01 per mg of dry air. The rain drop sizes are limited by the short simulation time and 2-dimensional kinematic setup. Different collision kernels and efficiencies should be tested for longer 3-dimensional simulations.

Figure 2: The collision efficiencies used in the presented simulations. The left panel shows the Hall (1980) collision efficiency. The middle panel zooms-in to show the Hall (1980) collision efficiency for small droplets. The right panel shows the collision efficiency for small droplets based on Pinsky et al. (2008).

• P. 16, Fig. 4b: What is meant by spec. in the panels title?
  It was supposed to mean specific concentration - i.e. concentration defined per unit mass of dry air. We have updated the caption of Fig. 4b.

• P. 18, l. 5: Eq. (5) shows that the oxidation of H2O2 is not independent of the acidity ([H+] in the denominator). (Although the dependence might be very weak.)
  Corrected.

• P. 18, l. 19-20: This is because ... aerosol particles. Is this a direct result of Eq. (2)?
  Yes. But it also depends on the speed of the dissolution of trace gases. We added the reference to Eq. (2).
• P. 1, ll. 20-21: clearly defined user interface is a slightly subjective statement.
P. 1, l. 21: Make clear that separation of concerns is a software design principle.
P. 2, l. 17: Define IPCC.
P. 2, ll. 33-34: All these references should be in one pair of parentheses.
P. 3, l. 20-22: As a service to the reader, it would be nice to state how many new attributes are added for the consideration of aqueous sulfur chemistry.
P. 4, l. 15: The citation is Lee et al. (2014) not Junghwa et al. (2014).
P. 4, ll. 22-23: The pH is a scale used to quantify the acidity of a solution. It is (linguistically) awkward to keep track of a scale and not the underlying physical meaning.
P. 6, Eq. (2): The underbrace needs to include all terms on the right-hand side of the equation.
P. 10, l. 4: These values are stated in Tab. 2 of Kreidenweis et al. (2003).
P. 12, ll. 7-8, 15: Semicolons in an enumeration of references are unusual.
P. 13, l. 21: the the
P. 15, l. 27, 29: Is Fig. 4f addressed of Fig. 4c?

All lines have been corrected. Thank you!

Interactive comment on libcloudph++ 1.1: aqueous phase chemistry extension of the Lagrangian cloud microphysics scheme by Anna Jaruga and Hanna Pawlowska Anonymous Referee #2

• p. 7, l. 14: The sentence is not clear. Should it read This approach does ensure. . . the total dissolved mass. . .does not exceed?

No. The dissolution of trace gases is calculated individually for each super-droplet. The trace gas mixing ratio is not updated after each super droplet dissolves the required trace gas mass. Instead, at the end of the dissolution calculation, the sinks from all super-droplets are summed and then applied to the ambient trace gas mixing ratio. In principle, when used with long model time-steps, this could result in dissolving more than the total trace gas mass available in the environment. To prevent that, short model time-steps or even sub-steps are required. The time-step length required by the dissolution scheme is similar to the time-step required by the condensation scheme.

The according sentence in the manuscript was corrected.

• p. 4, l. 8: A more detailed discussion of the errors would be useful to better appreciate the usefulness of your framework. How large are the statistical errors as compared to the numerical diffusion errors in previous schemes?

Unfortunately, it is difficult for us to provide an estimate of the total error due to the numerical diffusion in the bin schemes referenced in this manuscript. It depends on the number of bins used in the scheme, as well as the actual numerical algorithm used to integrate the equations.

The study by Li et al. (2017) provides a good comparison between the bin and particle-based schemes. As mentioned in the manuscript, the study by Unterstrasser et al. (2017) tests the behavior of different implementations of collisions in the super-droplet scheme. Additionally, Dzielen and Pawlowska (2017) tested the accuracy of the collision scheme in the super-droplet method when compared to the solutions of master equation, direct numerical simulations and the Smoluchowski equation. They found that for high super-droplet concentrations, the super droplet method resolves the collisions between droplets very accurately.

Interestingly, the work by Grabowski and Abade (2017) suggests that the fluctuations of the supersaturation should be included in super-droplet method and will result in an additional broadening of the droplet size distribution. In contrast to the broadening due to numerical diffusion in the bin schemes, such process can be accurately modeled in the super-droplet method when the turbulence dissipation rate in the model grid-box is known.

• p. 2, l. 1-6: This paragraph should be moved to the end of the introduction.

p. 2, l. 2: It should be also mentioned here that uptake processes between the gas and aqueous phases are included. As it is, it reads that only aqueous phase chemistry is considered but the full multiphase system is implemented.

p. 2, l. 10: SO2 can be oxidized within minutes or a few hours within clouds.
p. 5, l. 29, and throughout manuscript: $\Delta H/R$ are not correction coefficients. They should be called temperature dependence or enthalpy of ionization (= only $\Delta H$).

p. 7, l. 5, and throughout manuscript: The equilibrium dissolution constants are commonly referred to as Henry’s law constants. Their temperature dependencies are enthalpy of solution.

p. 7, l. 25: setup misspelled

p. 7, l. 29: Add unit M to [H+]

p. 7, l. 30: Why do you refer here to Table C2?

p. 9, l. 8, and throughout manuscript: I suggest using S(VI) instead of H2SO4.

p. 15, l. 31: new aerosol particles is misleading. It implies new particle formation. Please clarify.

p. 17, l. 7: minimize

p. 19, l. 7/8, and 14/15: These sentences sound awkward as SVI molecules are not oxidized. Try 92% of S(VI) originates from S(IV) oxidation by H2O2 (or similar).

All lines have been corrected. Thank you!

Additional changes to the manuscript:

- We changed the library version number from 1.1 to 2.0. The change reflects the new library programming interface related to the aqueous phase chemistry. The final version of the manuscript will be accompanied by the corresponding code release on GitHub
- To simplify we have removed from the abstract "(...) which allow reusing the created scheme from models implemented in other programming languages".
- p. 2, l.33: We added a list of names used to describe this process.
- p.14: We have deleted the last sentence.
- We have changed the Appendix A section name from User Interface to Programming Interface
- We have updated the list of references.

References


libcloudph++ 1.12.0: aqueous phase chemistry extension of the Lagrangian particle-based cloud microphysics scheme

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Abstract. This paper introduces a new scheme available in the library of algorithms for representing cloud microphysics in numerical models named libcloudph++. The scheme extends the Lagrangian microphysics scheme particle-based microphysics scheme with Monte-Carlo coalescence available in libcloudph++ to the aqueous phase chemical processes occurring within cloud droplets. The representation of chemical processes focuses on the aqueous phase oxidation of the dissolved SO\(_2\) by O\(_3\) and H\(_2\)O\(_2\). The Lagrangian Microphysics and Chemistry (LMC)-particle-based microphysics and chemistry scheme allows tracking the changes in the cloud condensation nuclei (CCN) distribution caused by both collisions between cloud droplets and aqueous phase oxidation.

The scheme is implemented in C++ and equipped with bindings to Python which allow reusing the created scheme from models implemented in other programming languages. The scheme can be used on either CPU or GPU, and is distributed under the GPL3-GPLv3 license. Here, the LMC particle-based microphysics and chemistry scheme is tested in a simple 0-dimensional adiabatic parcel model and then used in a 2-dimensional prescribed flow framework. The results are discussed with the focus on changes to the CCN sizes and compared with other model simulations discussed in the literature.

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1 Introduction

libcloudph++ is an open-source library of schemes for representing cloud microphysics in numerical models. It was first introduced in Arabas et al. (2015) where the authors present the different microphysics schemes available in the library, show its user-programming interface, and discuss its performance. The flagship component of libcloudph++ is the Lagrangian particle-based (i.e. particle tracking or particle-based"Lagrangian-in-droplet-radius-and-space") microphysics scheme. The scheme resolves the evolution of aerosol, cloud droplet, and rain drop\(^1\) size spectrum. It allows representing from the first principles cloud microphysical processes and is especially well suited to track changes in the CCN size distribution that are caused by clouds (i.e. cloud-aerosol processing). The scheme can be used in models of any dimensionality or dynamical core, and can be run on both CPU and GPU. The open-source availability of the code, its clearly defined user interface, and the

\(^1\) For convenience cloud droplets and rain drops will be often labeled together as water drops
Aqueous phase oxidation of sulfur is an important chemical reaction. Sulfur is emitted to the atmosphere by phytoplankton in the oceans as dimethyl sulfide and is then oxidized in the gas phase to sulfur dioxide trace gas. It is also emitted by anthropogenic and volcanic activity. The gas phase is oxidized in a matter of days either by the gas phase or aqueous phase reactions, see the review by Falooa (2009). The aqueous phase oxidation is reported as a dominant mechanism of production of sulfate. A numerical study by Barth et al. (2000) reports that for the in-cloud conditions, aqueous phase reactions accounts for 81% of sulfate production rate. According to their study total of ~50% – 60% of sulfate burden in the troposphere is produced by aqueous phase chemistry. Noteworthy, sulfate is a common component of aerosol particles (10%-67% of the submicron particle mass is made of sulfate, 32% on average; see Zhang et al., 2007). Based on analysis of 20 modeling studies, the review by Falooa (2009) marks wet deposition of aerosol sulfate, dry deposition of SO2 and heterogeneous (aqueous phase) oxidation of SO2 in aerosol particles and clouds as the most challenging to quantify in models. For an overview of the representation of sulfur oxidation in regional and global models see Ervens (2015).

The aqueous phase oxidation is reported as a dominant mechanism of production of sulfate: A numerical study by Barth et al. (2000) reports that for the in-cloud conditions, aqueous phase reactions accounts for 81% of sulfate production rate. According to their study total of ~50% – 60% of sulfate burden in the troposphere is produced by aqueous phase chemistry. The gas phase SO2 is oxidized in a matter of days by the gas phase reactions or within minutes or a few hours within clouds by aqueous phase reactions, see the review by Falooa (2009). From the cloud microphysics stand point, aqueous phase oxidation of sulfur is interesting because it affects the CCN within water drops. The reaction

\[
2 \text{SO}_2 + 4 \text{H}_2\text{O} \rightarrow 2 \text{H}_2\text{SO}_4 + 4 \text{H}^+ + 4 \text{e}^-
\]

is produced by aqueous phase chemistry. The gas phase SO2 is oxidized in a matter of days by the gas phase reactions or within minutes or a few hours within clouds by aqueous phase reactions, see the review by Falooa (2009).

From the cloud microphysics stand point, aqueous phase oxidation of sulfur is interesting because it affects the CCN within water drops. The reaction

\[
2 \text{SO}_2 + 4 \text{H}_2\text{O} \rightarrow 2 \text{H}_2\text{SO}_4 + 4 \text{H}^+ + 4 \text{e}^-
\]

is produced by aqueous phase chemistry. The gas phase SO2 is oxidized in a matter of days by the gas phase reactions or within minutes or a few hours within clouds by aqueous phase reactions, see the review by Falooa (2009).
increases the dissolved CCN mass. Collisions and the subsequent coalescence of water drops in as well as collisions between the aerosol particles and water drops are another in-cloud irreversible process that affects the CCN-sized aerosol particles. As the water drops collide and coalesce, the created new water drop carries within the combined CCN mass of all of its colliding predecessors. Efficient collisions between cloud droplets may quickly lead to the onset of precipitation which can in turn effectively cleanse the atmosphere from aerosol particles and water soluble trace gases. In non-precipitating clouds, aerosol particles that served as CCN are altered by cloud microphysical and chemical processes and then returned to the atmosphere after water drops evaporate (the process is referred to as the CCN deactivation, aerosol regeneration, aerosol recycling or aerosol resuspension). The cloud-processed aerosol particles can be observed in measurements (Hoppel et al., 1986, 1994), (Werner et al., 2014), (Hudson et al., 2015). Hoppel et al. 1986, 1994, Werner et al. 2014, Hudson et al. 2015). The case without precipitation reaching the surface is especially interesting as it allows for aerosol-cloud interactions to loop for several cloud life-cycles without removing the altered aerosol particles. The cloud-processed aerosol particles may again serve as CCN and influence microphysical properties of the next generation of clouds. The study of Pruppacher and Jaenicke (1995) estimates that on global average an atmospheric aerosol particle has been cycled 3 times by cloud systems. The LMC-Particle-Based Microphysics and Chemistry (PBMC) scheme introduced here offers a chance to represent the effects of such cloud-processing on CCN sizes stemming from both collisions between water drops and aqueous phase oxidation reaction within water drops. The PBMC can be used in multidimensional simulations with a fully coupled dynamics model, which has not been possible before. To the authors knowledge, the presented scheme is the first to represent the impact of both collisions and aqueous phase chemistry on the aerosol size spectrum in the Lagrangian-particle-based microphysics framework.

The This paper documents the extension of the particle-based microphysics scheme with a numerical scheme that represents aqueous phase chemical reactions inside cloud droplets and the uptake of the trace gases into cloud droplets. The representation of chemical reactions includes only the aqueous phase processes (i.e., no gas phase chemical reactions) and revolves around oxidation of sulfur dissolved in water drops to sulfate. Two reaction paths are considered – oxidation by ozone and by hydrogen peroxide. In total, six trace gases are included in the chemistry description: sulfur dioxide (SO2), ozone (O3), hydrogen peroxide (H2O2), carbon dioxide (CO2), nitric acid (HNO3), and ammonia (NH3). Their dissolution, and if applicable, dissociation is resolved. The structure of the presented work is as follows: Section 2 presents briefly the Lagrangian-particle-based scheme available in libcloudph++. Section 3 discusses the design and user-programming interface of the new aqueous chemistry scheme. Section 4 presents the validation tests of the new scheme. Section 5 discusses the results from simulations where the LMC-PBMC scheme is incorporated into a simple model of a stratuscumulus cloud. The effects of both collisions between water drops and aqueous phase oxidation of sulfur on the aerosol particle size distribution are presented.

2 Lagrangian-Particle-based microphysics scheme

The Lagrangian-particle-based scheme used in this work is described in detail in Arabas et al. (2015) and this section only briefly summarizes its major concepts. In the Lagrangian-particle-based approach to modeling cloud microphysics the compu-
tional domain is filled with “numerical point particles” representing a multiplicity of real particles (aerosol particles, cloud droplets or rain drops) of the same properties. Following the nomenclature introduced by Shima et al. (2009), the “numerical particles” are labeled here as super-droplets (SD). Each SD has a set of attributes describing the properties of the aerosol particles or water drops it represents. As discussed in Arabas et al. (2015), for microphysical purposes, the required attributes are: multiplicity ($N$), the position of SD in the computational domain, the wet radius ($r_w$), the dry radius ($r_d$) and the hygroscopicity parameter ($\kappa$). The aqueous phase chemistry scheme extends the list of required attributes by masses of chemical compounds dissolved in droplets. Eight new attributes are needed, see Sec. 3 for details.

The key attribute of the Lagrangian-particle-based microphysics scheme is the SD multiplicity. Multiplicity defines the number of aerosol particles or water drops represented by a given SD. All particles represented by one SD are assumed to be identical. Using multiplicites allows to reduce the complexity of the problem and enables efficient numerical computations.

The Lagrangian-particle-based scheme used here requires no division into artificial categories of aerosol particles, cloud and rain water, as it is often done in the bulk schemes, for example Kessler (1995), Seifert and Beheng (2001), Morrison and Grabowski (2007). All the modeled microphysical processes are represented by calculating the changes to the SD attributes. The equation of condensational growth is solved iteratively for each SDs wet radius (see Sec. 5.1.3 in Arabas et al. (2015) for details). The process of condensational growth from deliquescent aerosol particles to cloud droplets is thus resolved and no additional parametrisation of cloud droplet activation is required as it is again often done in the bulk microphysics schemes, see for example Morrison and Grabowski (2007).

Following Arabas et al. (2009) the collisions between SDs are represented using Monte-Carlo scheme (see Sec. 5.1.4 in Arabas et al. (2015) for details).

Information about SD attributes is retained within the model throughout the whole simulation. This means that the size distribution of both water drops and aerosol particles in each computational grid-cell can be easily obtained by taking into account the SD attributes of $r_w$, $r_d$ and $N$. As a result, the Lagrangian-particle-based scheme is capable of resolving the changes to both aerosol and water drop size distributions. The same functionality is offered by the 2-dimensional bin schemes, for example Ovchinnikov and Easter (2010) or Lebo and Seinfeld (2011). However, the Lagrangian-particle-based approach greatly reduces the numerical diffusion errors. As discussed in Unterstrasser et al. (2017), it does introduce statistical errors, i.e. fluctuations between different realizations of the same collision/coalescence scenario. These errors are easier to minimize than diffusion numerical errors, for example by increasing the number of SDs in the computational domain or by averaging over an ensemble of simulation runs. Lagrangian schemes solve ordinary differential equations instead of the partial differential equations encountered in the bin schemes, which is computationally more efficient. Dziekan and Pawlowska (2017) showed that for high SD concentrations the super-droplet method accurately represents collisions between the drops (with regard to

---

3 It is a volume equivalent radius for solute in the water drop.
4 It is a volume equivalent radius for solute in the water drop.
5 Following Ghan et al. (2001) and Petters and Kreidenweis (2007) it is a single parameter representing the hygroscopicity of the solvent. In this work we use the notation from Petters and Kreidenweis (2007).
6 Following Petters and Kreidenweis (2007) it is a single parameter representing the hygroscopicity of the solvent.
the expected value and the standard deviation of the autoconversion time. An interesting comparison between the bin and the particle-based schemes is provided by Li et al. (2017). On a side note, Grabowski and Abade (2017) show that an additional scheme modeling the broadening of droplet spectra due to supersaturation fluctuations might be necessary.

The collision efficiency used in this study is based on Hall (1980) and Pinsky et al. (2008). It is well suited for representing the collisions between water drops. An additional collision efficiency look-up table based for example on Ladino et al. (2011) or Ardon-Dryer et al. (2015) should be used to study the collection of submicron aerosol particles by droplets. Similarly, additional collision efficiency corrections based for example on Chen et al. (2018) should be applied to study the effects of turbulence on the aerosol size distribution.

The Lagrangian Particle-based methods are becoming a well known tool for studying cloud microphysics in both warm-clouds (Shima et al. (2009), Arabas and Shima (2013), Arabas and Shima (2017), Andrejczuk et al. (2010) Andrejczuk et al. (2010), Riechelmann et al. (2012), Arabas and Shima (2013), Andrejczuk et al. (2014), Lee et al. (2014), Naumann and Seifert (2015), Sato et al. (2017), Hoffmann (2017), Grabowski et al. (2018), Sardina et al. (2018) ); warm rain clouds (Riechelmann et al. (2012), Lee et al. (2014), Naumann and Seifert (2015)); and ice-phase clouds (Sölch and Kärcher (2010), Unterstrasser and Sölch (2014)). None of the above however, included description of the aqueous phase chemical reactions happening within cloud droplets.

3 Aqueous phase chemistry scheme

In order to represent the chemical composition of water drops the aqueous-phase chemistry scheme extends the list of SD attributes. The additional attributes are defined as the total mass of each of the chemical compounds in a given SD (including both the dissolved and, if applicable, dissociated fraction). An additional variable – the mass of the H⁺ ions – is also added, in order to keep track of SD’s pH acidity. This results in eight new SD attributes needed for simulations with aqueous phase chemistry:

- the total mass of dissolved O₃,
- the total mass of dissolved H₂O₂,
- the total mass of dissolved SO₂ (including: SO₂ + H₂O, HSO₃⁻ and SO₄²⁻),
- the total mass of dissolved CO₂ (including: CO₂ + H₂O, HCO₃⁻ and CO₃²⁻),
- the total mass of dissolved NH₃ (including: NH₃ + H₂O and NH₄⁺),
- the total mass of dissolved HNO₃ (including: HNO₃(aq) and NO₃⁻),
- the total mass of created H₂SO₄ (including: HSO₄²⁻ and SO₄²⁻)
- the total mass of H⁺ ions.
The scheme needs to be coupled to the driver model (i.e., dynamical core) that provides information about the environment in which SD are immersed (i.e., temperature, humidity, trace gas mixing ratio, and wind speed). The representation of aqueous phase chemistry more than doubles the number of required SD attributes and significantly increases the computational time. On the other hand, thanks to the added attributes, the mass of any ion for any SD can be easily diagnosed using just a dissociation constant. This, in turn, allows for a very straightforward representation of the aqueous chemical processes and does not call for any additional parametrisation.

All aqueous-phase chemistry included in the scheme is formulated under the assumption that solution droplets are diluted. Therefore, in the LMC-PBMC scheme, chemical processes are only performed for SDs with ionic strength smaller than 0.02 moles/liter (the same criterion is used for example in Ovchinnikov and Easter, 2010). In practice, this condition results in excluding from aqueous chemistry calculations SDs with small wet radii (i.e., SDs representing haze particles and very small cloud droplets). Exclusion of SDs with small wet radii also prevents numerical issues during condensation procedure when changes in dry radius caused by oxidation could prevent convergence of the condensation scheme during the initial rapid growth of cloud droplets during activation.

Combining the Lagrangian-particle-based microphysics scheme with aqueous phase chemistry is straightforward. Condensation/evaporation does not affect the chemical attributes of SDs. During collisions the mass of chemical compounds is summed when recalculating SD attributes (it is an extensive parameter). In principle the \( \kappa \) attribute should be recalculated in every time-step based on the new chemical composition of each SD. However, the \( \kappa \) values relevant for this study are very similar - the \( \kappa \) value of ammonium bisulfate is 0.61 (Petters and Kreidenweis, 2007) and of sulfuric acid is 0.64 (Kim et al., 2016). Therefore, the hygroscopicity parameter is assumed to be constant.

As discussed in Kreidenweis et al. (2003), one of the major sources of errors in modeling the in-cloud aqueous phase oxidation are the uncertainties when resolving the cloud droplet size distribution. It follows that combining the aqueous phase chemistry model with the very detailed Lagrangian microphysics scheme is beneficial. This will be discussed further in Sec. 4.

3.1 Dissociation

Dissociation is a reversible process of splitting of the molecules dissolved in water drops into ions. It is treated as an equilibrium process and is described using the dissociation constants. The dissociation constant of chemical compound \( A \) is denoted here by \( K_A \). The dissociation constants are corrected for temperature using the formula of van 't Hoff (1885):

\[
K_A(T) = K_A(T_0) \exp \left( \frac{-\Delta H_D}{R} \left( \frac{1}{T} - \frac{1}{T_0} \right) \right),
\]

where \( \Delta H_D \) denotes the reaction enthalpy of dissociation at constant temperature and pressure, \( T \) is the temperature of air and \( R \) is the gas constant. The list of considered dissociation constants as well as their temperature correction coefficients is available in Tab. C2. The dissociation of water, although very small, is also taken into account\(^5\). No temperature

\(^{5}\)The concentration of undissociated water molecules is so big that it is usually assumed constant and it traditionally multiplies the dissociation constant of water. This leads to a different definition of the dissociation constant for water: \( K_{\text{H2O}} = [\text{H}^+] [\text{OH}^-] \)
It is assumed that there is no electric charge of water drops and therefore the concentrations of positive and negative ions created during dissociation should balance each other. Using dissociation constants (see Tab. C2), all ion concentrations can be expressed as a function of the total concentration of the dissolved chemical compounds and the concentration of H\(^+\) ions. For example derivation see Chap. 7.6.2 in Seinfeld and Pandis (2016). The neutral charge condition can be expressed as

\[
\begin{align*}
[H^+] + \frac{[N^{III}][H^+]}{K_{H2O} + K_{NH3}[H^+]} &= \frac{[N^{IV}][K_{HNO3}]}{[H^+] + K_{HNO3}} + \frac{[S^{IV}][K_{SO2}([H^+] + 2K_{HSO3})]}{([H^+]^2 + [H^+][K_{SO2} + K_{SO2}K_{HSO3}])} + \\
&\frac{[K_{H2O}][K_{H2SO4}]}{[H^+] + [K_{H2SO4}]} + \frac{[C^{IV}][K_{CO2}([H^+] + 2K_{HCO3})]}{([H^+]^2 + [H^+][K_{CO2} + K_{CO2}K_{HCO3}])}.
\end{align*}
\]

The \([\ ]\) brackets denote the concentration of each of the chemical compounds, (traditionally defined in units of moles per liter), capital letters denote chemical compound and roman numbers mark its oxidation state. In Eq. (2) the dissociation constants of SO\(_2^2^-\), CO\(_2^3^-\), and SO\(_2^4^-\) ions (i.e. \(K_{HSO3}\), \(K_{HCO3}\), and \(K_{H2SO4}\)) are multiplied by a factor of two, to take into account the respective electric charge number of those ions.

Equation (2) has only one unknown variable – the new equilibrium concentration of the H\(^+\) ions. The new concentration is obtained iteratively using a numerical root-finding algorithm\(^6\). The algorithm searches for solution between pH = -1 and pH = 9. The lower bound for the pH scale is unrealistically acidic and is only necessary at the start of the simulation when the initial SDs have very small volume and thus highly acidic pH. The upper bound is set arbitrarily, but is sufficient for the expected pH of the modeled droplets. At the end of the dissociation procedure the mass of H\(^+\) ions is updated based on the new equilibrium concentration.

When the SD wet radius is quickly changing, for example during the initial condensational growth of cloud droplet or rain drop evaporation, the dissociation procedure requires small time-steps to reach convergence. The time-step used in dissociation procedure can be divided into user-specified number of sub-steps in order to prevent limiting the overall simulation time-step by dissociation.

### 3.2 Dissolution

The amount of the chemical compound that can dissolve into water drop from the gas phase is proportional to its partial pressure above the surface of the drop. Due to the longer timescale of the process, in contrast to dissociation, the transfer between the gas and liquid phase is not treated as an instantaneous process. Assuming that the water drop is internally mixed, the gas-liquid transfer is limited by the diffusion of gas phase particles to the drop surface (gas-phase limitation) and the probability that the molecule will enter the drop after collision (interfacial limitation). Following chapter 8.4.2 in Warneck (1999), for a chemical

\(^6\) TOMS 748 algorithm from Boost library. See www.boost.org/doc for documentation and Alefeld et al. (1995) for derivation
compound “A” the rate of transfer from the gas phase to the aqueous phase equals
\[
\frac{d[A]}{dt} = \left( \frac{4r_w}{3\sqrt{\langle v \rangle} M_A \alpha} + \frac{r_w^2}{3D_A} \right) \left( c_{\text{avg}} - \frac{[A]}{H_{A}^{\text{eff}} RT} \right) \tag{3}
\]
where \(D_A\) and \(\alpha M_A\) are the diffusion and mass accommodation coefficients for the chemical compound “A”, \(\langle v \rangle = \sqrt{\frac{8RT}{\pi M_A}}\) is the average velocity of the molecules calculated from the Maxwell-Boltzmann distribution function, \(M_A\) is the molar mass of the chemical compound “A”, \(c_{\text{avg}}\) is the ambient concentration of the trace gas “A” and \(H_{A}^{\text{eff}}\) is the effective Henry’s law constant of the chemical compound “A”. The dissolution (i.e., the equilibrium dissolution constant). The Henry’s law constants depend on the temperature following similar relation as for dissociation Eq. (1). Table C3 shows the equilibrium dissolution-Henry’s law constants and their temperature correction dependencies and Tab. C4 presents the diffusion and mass accommodation coefficients. The term “effective” marks that the dissolution constants take into account the increase of the efficiency due to dissociation (see Seinfeld and Pandis (2016) chap. 7.4–7.3 for the exact equations). Equation (3) is solved for each SD and for each of the considered trace gases. It is solved implicitly with respect to aqueous-phase concentration and explicitly with respect to the gas-phase concentration. The input ambient trace gas concentration is calculated from the trace gas mixing ratio provided by the driver model to which the LMC-PBMC scheme is coupled. Obtained aqueous phase concentration is recalculated to the mass of dissolved chemical compounds and the corresponding SD attribute is updated. The changes in the ambient trace gas mixing ratios are calculated by LMC-PBMC scheme by summing the changes in chemical composition in all SDs in a given grid-cell and then subtracting them from the trace gas mixing ratio of the driver model. This approach does not ensure that per each time-step the total dissolved mass of a given trace gas. To ensure that the sum of sinks from each SD does not exceed the available ambient mixing ratio. To prevent that, trace gas mixing ratio, a relatively short time-steps should be applied. If necessary the user can divide the model time-step into sub-steps.

3.3 Oxidation

The reaction rates of oxidation by ozone and hydrogen peroxide can be described as (Hoffmann and Calvert, 1985):

\[
R_{O_3} = \left. \frac{d[S^{VI}]}{dt} \right|_{O_3} = \left( k_0 + k_1 \frac{[SO_2]}{[H^+]} + k_2 \frac{[SO_2][HSO_3]}{[H^+]^2} \right) [O_3][SO_2 + H_2O] \tag{4}
\]

\[
R_{H_2O_2} = \left. \frac{d[S^{VI}]}{dt} \right|_{H_2O_2} = \frac{k_3 [SO_2]}{1 + k_4 [H^+]} [H_2O_2][SO_2 + H_2O] \tag{5}
\]

where \(R_A\) is the reaction rate of the chemical compound “A” and \(k_{0,...,4}\) are the reaction rate coefficients. \(k_{0,...,4}\) depend on the temperature following similar relation as for dissociation Eq. (1). Table C5 shows the values of reaction rate coefficients and their temperature correction dependence coefficients.

Equations (4) and (5) return the new concentration of \(S^{VI}\) created in each SD in each time-step. Based on the new concentration, the new mass of \(S^{VI}\) and the new dry radius are calculated and the corresponding SD attributes are updated. The dry particle density of 1.8 g/cm\(^3\) is assumed while evaluating the dry radius from the \(S^{VI}\) mass.
For the typical atmospheric conditions, say pH between 3 and 6 (i.e. [H⁺] between 10⁻³ M and 10⁻⁶ M), it can be said that the rate of oxidation by H₂O₂ does not depend on pH (see Tab. C2 for the dissociation constant values). In contrast, oxidation by ozone depends strongly on pH of the solution and can become very fast if pH is high. For example, increasing pH by 1 point results in approximately 100 increase in O₃ reaction rate.

3.4 Initialization

The initial aerosol is assumed to be ammonium bisulfate (NH₄HSO₄), with dry particle density of 1.8 g cm⁻³. Using dry particle density and dry radius of each SD, the initial mass of H⁺, NH⁺, and SO₄²⁻ ions is calculated. The initial mass of other molecules and ions is equal set to zero and is therefore not in equilibrium with the initial ambient trace gas conditions. For the initial conditions above supersaturation where supersaturation is present in the environment, it is advisable to allow for a spin-up period with only condensation/evaporation and the equilibrium chemical processes enabled, to allow the model to reach equilibrium. Such initial conditions are mostly relevant for the kinematic models.

4 Validation

The LMC–PBMC scheme is set to reproduce results from model intercomparison study by Kreidenweis et al. (2003), where several bulk and bin moving-bin schemes representing cloud microphysics and aqueous-phase chemistry were tested in an adiabatic parcel model setup. A parcel model. The parcel model used here is a 0-dimensional model that represents an idealized scenario of a finite volume of air rising adiabatically with a constant vertical velocity. As the parcel of air raises, its temperature decreases leading to supersaturation. This results in activation and further condensational growth of cloud droplets. For the studied oxidation reaction, the presence of liquid water enables aqueous-phase chemical reactions and leads to creation of sulfuric acid within cloud droplets. The collisions between cloud droplets are not included in the parcel simulations to allow an easy comparison with Kreidenweis et al. (2003).

The initial conditions are the same as in Kreidenweis et al. (2003) and are provided for convenience in Tab. 1. The simulation starts below cloud base (i.e. with subsaturation). The initial aerosol is ammonium bisulfate and the initial aerosol particle size distribution is assumed to be lognormal with one mode

$$n(r_d) = \frac{n_{tot}}{r_d \sqrt{2\pi \ln(\sigma_g)}} \exp \left( -\frac{(\ln(r_d) - \ln(r\bar{d}))^2}{2\ln^2(\sigma_g)} \right)$$

where $n(r_d)$ is the spectral density function of aerosol particle sizes, $n_{tot}$ is the total aerosol concentration, $r\bar{d}$ is median radius and $\sigma_g$ is the geometric standard deviation. See Sec. 5.1.6 in Arabas et al. (2015) for the details on how the super-droplet dry and wet radii are initialized.

The parcel model employed in this study uses dry air density $\rho_d$, dry air potential temperature $\theta$, water vapor mixing ratio $r_v$, and mixing ratios of ambient trace gases as model variables. In order to calculate $\rho_d$ at each time-level (or each height-level of the parcel ascent) the model assumes needs to assume a vertical profile of pressure. In the presented simulations the assumed pressure profile is obtained by integrating the hydrostatic equation and assuming that density of air is constant.
Table 1. Initial conditions for the adiabatic parcel test.

<table>
<thead>
<tr>
<th>factor</th>
<th>value</th>
<th>units</th>
</tr>
</thead>
<tbody>
<tr>
<td>number of super-droplets</td>
<td>1024</td>
<td>-</td>
</tr>
<tr>
<td>time-step</td>
<td>0.1</td>
<td>s</td>
</tr>
<tr>
<td>temperature at (t = 0)</td>
<td>285.5 ± 0.2</td>
<td>K</td>
</tr>
<tr>
<td>pressure at (t = 0)</td>
<td>950</td>
<td>hPa</td>
</tr>
<tr>
<td>relative humidity at (t = 0)</td>
<td>95</td>
<td>%</td>
</tr>
<tr>
<td>updraft velocity</td>
<td>0.5</td>
<td>m/s</td>
</tr>
<tr>
<td>median radius</td>
<td>0.4</td>
<td>µm</td>
</tr>
<tr>
<td>geometric standard deviation</td>
<td>2</td>
<td>-</td>
</tr>
<tr>
<td>total aerosol number concentration</td>
<td>566</td>
<td>cm⁻³</td>
</tr>
<tr>
<td>dry particle density</td>
<td>1.8</td>
<td>g/cm³</td>
</tr>
<tr>
<td>hygroscopicity</td>
<td>0.61</td>
<td>-</td>
</tr>
<tr>
<td>concentration of (\text{SO}_2) at (t = 0)</td>
<td>0.2</td>
<td>ppb-v</td>
</tr>
<tr>
<td>concentration of (\text{O}_3) at (t = 0)</td>
<td>50</td>
<td>ppb-v</td>
</tr>
<tr>
<td>concentration of (\text{H}_2\text{O}_2) at (t = 0)</td>
<td>0.5</td>
<td>ppb-v</td>
</tr>
<tr>
<td>concentration of (\text{CO}_2) at (t = 0)</td>
<td>360</td>
<td>ppm-v</td>
</tr>
<tr>
<td>concentration of (\text{HNO}_3) at (t = 0)</td>
<td>0.1</td>
<td>ppb-v</td>
</tr>
<tr>
<td>concentration of (\text{NH}_3) at (t = 0)</td>
<td>0.1</td>
<td>ppb-v</td>
</tr>
</tbody>
</table>

at each height level (piecewise constant) and equal to 1.15 kg m⁻³. The assumed density is based on the density provided in Tab. 3 in Kreidenweis et al. (2003). Then, at each level, \(\rho_d\) is calculated from the ideal gas law taking into account the current \(r_v\) and \(\theta_c\):

\[
\rho_d(p, \theta, r_v) = \frac{p - p_v(p, r_v)}{R_d(\frac{p}{p_{1000}})^{\frac{R_g}{c_{pd}}}},
\]

(7)

where: \(p_v\) represents partial pressure of water vapor, \(p_{1000}\) stands for pressure equal 1000 hPa that comes from the definition of potential temperature, \(R_d\) is the gas constant for dry air and \(c_{pd}\) is the specific heat at constant pressure for dry air. Because the simulated air parcel is assumed to be adiabatic, only processes resolved by the Lagrangian-particle-based scheme can change \(\theta\), \(r_v\) and other trace gas mixing ratios. In each model time-step, the Lagrangian-particle-based microphysics scheme changes \(\theta\) and \(r_v\) according to Eq. 25 and 26 from Arabas et al. (2015). The changes in the trace gas mixing ratios are resolved following procedure discussed in Sec. 3.2. It is assumed that the initial mass of dry air within the parcel is 1 kg.
Figure 1 shows the general physical and chemical conditions from the cloud base up to the end of the test run 1.2 km above the cloud base. Two vertical axes are used, representing either the time or the height above the cloud base. Figure 1a shows the liquid water mixing ratio (LWC). The increase in LWC is linear and the LWC reaches above 2 g/kg at height 1.2 km above the cloud base. Figure 1b shows the total \( \text{SO}_2 \) concentration (both in gas phase and dissolved in water) in ppb units. The concentration of \( \text{SO}_2 \) is decreasing due to oxidation taking place in cloud droplets. Figure 1c shows the water volume weighted average pH of cloud droplets. The pH near the cloud base is very low due to acidic nature of the assumed initial aerosol and small size of activated cloud droplets. As the drops grow bigger and become more diluted, the average pH increases. Figure 1 compares well with Fig. 1 from Kreidenweis et al. (2003), the overall differences between the two figures are small. The biggest discrepancy is in the LWC profile. This is arguably caused by different pressure profiles in the two parcel models (the pressure profile in Kreidenweis et al. (2003) is not specified).

At the end of the test simulation, \( 84.85\% \) of \( \text{SO}_2 \) is converted into \( S^{VI} \) and the final water volume weighted average pH is equal to 4.834.86. The total sulfate production is 168 ppt with 100–171 ppt with 99 ppt produced by the \( \text{H}_2\text{O}_2 \) reaction path and 68–72 ppt produced by the \( \text{O}_3 \) reaction path. Based on Fig. 2 in Kreidenweis et al. (2003), the range of average pH values reported by different size resolving (bim-moving-bin) schemes was between 4.82 and 4.85, and the range of total sulfate production values was between 170 – 180 ppt. Based on Fig. 3 in Kreidenweis et al. (2003) the production by \( \text{H}_2\text{O}_2 \) ranged between 85 and 105 ppt, and by \( \text{O}_3 \) between 70 and 85 ppt for the size resolving schemes. In short, the results from the Lagrangian particle-based scheme are close to the range of values reported by bim-moving-bin schemes.
Microphysics schemes taking part in the Kreidenweis et al. (2003) intercomparison study reported significant differences between the number of activated cloud droplets. Based on Tab. 2 in Kreidenweis et al. (2003) the droplet number concentration at the cloud base varied between 275 and 358 cm$^{-3}$. One of the differences between the bin-moving-bin schemes responsible for causing this discrepancy is the different water vapor mass accommodation coefficient $\alpha_M$ leading to different predicted maximum supersaturation. Figure 6 in Kreidenweis et al. (2003) shows that the observed maximum supersaturations were lower (between 0.23% – 0.26%) for schemes using high values of $\alpha_M$ (either 0.5 or 1). In contrast, a scheme using $\alpha_M = 0.042$ predicted maximum supersaturation equal to 0.37%. The Lagrangian-particle-based scheme used in this study reports concentration $264 - 269$ cm$^{-3}$ at RH $= 1$ and $272$ cm$^{-3}$ at the level of maximum supersaturation. The maximum supersaturation is equal to 0.27%. The Lagrangian-particle-based scheme assumes $\alpha_M$ equal to unity and therefore fits with the trend of high $\alpha_M$ causing lower supersaturation presented in Fig. 6 in Kreidenweis et al. (2003).

Another cause for the discrepancy between the bin schemes listed in the intercomparison study are the different sizes and locations of bins in different models, see also the discussion in Arabas and Pawlowska (2011). Along those lines, here it is tested how sensitive the Lagrangian-particle-based scheme is to the number of SDs. The results of this test are summarized in Fig. 2 showing the cloud droplet concentration at the cloud base (a), the maximum supersaturation (b), the average pH (c) and the total sulfate production (d). The results are plotted against the logarithm of base two of the number of SDs in the computational domain (meaning that “0” represents one SD and “10” represents 1024 SDs). All values seem to converge for SD numbers greater than 64. The average pH, maximum supersaturation and total sulfate production do not change for those four test-runs. The concentration of droplets at the cloud base varies little (between 265 cm$^{-3}$ and 274 cm$^{-3}$269 cm$^{-3}$ and 281 cm$^{-3}$). The cloud droplet concentration from simulation with 128 SDs is the largest outlier. The concentrations from simulations with SD number between 256-512 and 1024 vary between 274 cm$^{-3}$ and 272 cm$^{-3}$269 cm$^{-3}$. For SD numbers between 32 and 128-64 there are insignificant changes in the maximum supersaturation. The values of pH vary by 0.01 and the total sulfate production increases by 1 ppt. There are, however, large differences between the number of droplets at the cloud base (between 265 cm$^{-3}$ and 281 cm$^{-3}$ and 332 cm$^{-3}$). This confirms the observations from Kreidenweis et al. (2003) that the predicted cloud droplet number concentration strongly depends on the representation of the size distribution of modeled aerosol particles and cloud droplets and that this may become a major source of uncertainties in the microphysics representation. Decrease in the SD number below 32 leads to a big variance in the cloud droplet concentration as well as other parameters.

Figure 3 shows the simulated modification of the aerosol size distribution. Red line depicts the initial distribution and the green line shows model state at the end of adiabatic parcel test. For convenience, Fig. 3 uses both logarithmic (left panel) and linear (right panel) scale on ordinate. The change in aerosol size distribution is caused by oxidation. Aerosol particles that are too small to become cloud droplets are not affected by aqueous phase oxidation and they do not grow in size. Large aerosol particles grow in size due to $S^VI$ production during oxidation, but the increase in size is small compared to their initial size. The smallest activated aerosol particles are those that are affected most by oxidation. The increase in their size due to the produced $S^VI$ is the largest compared to their initial size. In short, oxidation produces a “gap”, often labeled the “Hoppel minimum”, between the CCN processed by the cloud and the smaller unactivated aerosol particles.
The effect of in-cloud sulfate production on aerosol particle size distribution presented in Fig. 3, combined with other tests presented in this chapter, document the correctness of the implementation of aqueous chemistry in the Lagrangian particle-based scheme. The formation of the “Hoppel minimum” was reported by many observational studies, see Hoppel et al. (1994); Bower et al. (1997) and Hudson et al. (2015). Figure 3 compares well with aerosol size distribution plots from the intercomparison study shown in Fig. 9 in Kreidenweis et al. (2003). Other numerical schemes also reported the formation of “Hoppel minimum”, see for example Flossmann (1994); Feingold and Kreidenweis (2000, 2002); Ovchinnikov and Easter (2010). Flossmann (1994), Feingold and Kreidenweis (2000, 2002) and Ovchinnikov and Easter (2010). The work by Cantrell et al. (1999) shows that the maximum supersaturation and the cloud droplet concentration in the clouds which processed the aerosol particles can be inferred based on the location of the aerosol size distribution minimum.
Figure 3. Modification of the dry aerosol sulfate mass. Red line shows the initial condition and green line shows the final model state. Left panel uses logarithmic ordinate and right panel uses linear ordinate.

5 Example simulations

5.1 2D kinematic model

The kinematic model mimics a single 2D eddy spanning a stratocumulus cloud deck and a boundary layer below it. The model is based on a test scenario from the 8th International Cloud Modeling Workshop (ICMW; Muhlbauer et al., 2013, case 1). The velocity field is prescribed as in Szumowski et al. (1998); Morrison and Grabowski (2007); Rasinski et al. (2011). The same model was used when presenting the initial release of libcloudph++, see Sec. 2 in Arabas et al. (2015) for the details of the model formulation. The model operates on the Eulerian grid. At each model time-step, the kinematic model is based on the open-source library of parallel MPDATA-based solvers for systems of generalised transport equations, see Jaruga et al. (2015).

The temperature, moisture, and trace gas fields are discretized on the Eulerian grid and are advected using the prescribed velocity field. Then, the model variables are passed to the LMC-PBMC scheme, where the microphysical and chemical processes are resolved. Finally, the source and sink terms due to microphysics and chemistry are calculated and applied in each model grid-cell as described in Sec. 2 and 3.

The collisions between water drops are represented using the geometric kernel with collision efficiency for big drops (i.e., radius greater than 20 μm) from Hall (1980) and for small droplets (i.e., radius smaller than 20 μm) from Pinsky et al. (2008). For big drops, the collision efficiencies were obtained from the fit to measurements, see Hall (1980). For small droplets, the collision efficiencies were based on numerical simulations taking into account turbulence typical for stratocumulus clouds, see Pinsky et al. (2008). The collision efficiencies are provided via a look-up tables for different drop sizes.

The initial conditions are summarized in Tab. 2. The computational domain size is 1.5 km in both directions and the computational grid is composed of 75×75 cells of equal size (the grid lengths are 20 m) and is periodic in the horizontal direction.
The initial air density profile corresponds to the hydrostatic equilibrium with the pressure of 1015 hPa at the bottom of the domain. At the beginning of the simulation it is assumed that there is no condensed water, and the initial profiles of $\theta$ and $r_V$ are constant with altitude. To keep the simulation setup simple and due to a relatively low vertical extent of the computational domain, the initial trace gas volume fractions are also assumed to be constant with altitude. This unrealistic initial condition results in very high initial supersaturation in the upper part of the domain. As a consequence a $10^5$ second ($\sim 2h 45\text{min}$) spin-up period is necessary to allow for the simulated water drops to reach equilibrium with their environment. During the spin-up only the reversible processes (condensation and evaporation, dissolving of trace gases and dissociation into ions) are allowed and the supersaturation is limited to 5% (RH=1.05). After spin-up the simulations are run for 30 minutes. The chosen simulation time is enough to deplete the SO$_2$ available in the cloudy part of the domain as well as to create precipitation.

Similarly to the adiabatic parcel test, the initial aerosol is ammonium bisulfate and the aerosol particle size distribution is lognormal with one mode. The initial condition for trace gases is defined in terms of volume fractions and then translated to mixing ratios that serve as the model variables. The initial SO$_2$, O$_3$ and H$_2$O$_2$ volume fractions are taken from the simulation setup used in Ovchinnikov and Easter (2010). The values for SO$_2$ and O$_3$ are based on the measurements from MASE campaign (Wang et al., 2008) and the value for H$_2$O$_2$ is based on the representative values for the Eastern Pacific Ocean (Genfa et al., 1999). The NH$_3$, HNO$_3$ and CO$_2$ volume fractions are the same as in the parcel test from Sec. 4.

The setup detailed in Tab. 2 corresponds to “very clean conditions” (i.e. low aerosol particle concentrations). Initial aerosol particle sizes are also relatively small. Three additional simulation cases are studied to check the sensitivity of the model to different conditions. In case1 the reversible chemical processes are allowed, but oxidation is prohibited. In case2 the initial volume fraction of NH$_3$ is increased and in case3 the initial aerosol size distribution is changed. The conditions for all the sensitivity simulation cases are summarized in Tab. 3.

As discussed in Flossmann (1994), the initial chemical scenario is idealized. For instance, although the initial conditions represent clean maritime environment, the setup lacks sea salt aerosol particles. As discussed by Twohy et al. (1989), sea salt aerosol particles are alkaline, which may in turn increase the pH of water drops and thus affect the oxidation rate. On the other hand, a study by von Glasow and Sander (2001) indicates that alkaline sea salt particles are quickly converted to acidic due to the uptake of HCl vapor. More importantly, including sea salt would result in aerosol particles with very different hygroscopicity values ($\kappa$ of ammonium bisulfate is 0.61, whereas $\kappa$ of NaCl is 1.28; Petters and Kreidenweis, 2007). Including sea salt would also result in the initial bi-modal size distribution with one mode representing smaller ammonium bisulfate aerosol particles and the second mode representing larger sea salt particles. In general, including sea salt should result in a very different condensational growth of aerosol particles. The setup used in this study also lacks other particles containing sulfate, such as ammonium sulfate or sulfuric acid aerosol particles. The reason behind the chosen setup is that this is the first attempt to include aqueous chemistry into the Lagrangian cloud microphysics and therefore the decision was made to start with the simplified setup.

The initial aerosol size distribution parameters are based on the test cases studied in Feingold and Kreidenweis (2002). The discussion presented in their study introduced two regimes for oxidation with regard to the mean aerosol size $r_d$ and precipitation: (i) for relatively small initial $r_d$ production of sulfate enhances precipitation, (ii) for relatively big initial $r_d$
Table 2. Initial conditions for base case of 2-dimensional kinematic model.

<table>
<thead>
<tr>
<th>factor</th>
<th>value</th>
<th>units</th>
</tr>
</thead>
<tbody>
<tr>
<td>number of super-droplets</td>
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<td>#/grid-cell</td>
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<tr>
<td>model time-step</td>
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<td>s</td>
</tr>
<tr>
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<td>s</td>
</tr>
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<td>K</td>
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<tr>
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<td>g/kg</td>
</tr>
<tr>
<td>pressure at z = 0</td>
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<td>hPa</td>
</tr>
<tr>
<td>median radius</td>
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<td>µm</td>
</tr>
<tr>
<td>geometric standard deviation</td>
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<td>-</td>
</tr>
<tr>
<td>total aerosol number concentration</td>
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<td>cm⁻³</td>
</tr>
<tr>
<td>dry particle density</td>
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<td>g/cm³</td>
</tr>
<tr>
<td>hygroscopicity</td>
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<td>-</td>
</tr>
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<td>concentration of SO₂ at t=0</td>
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<td>ppb-v</td>
</tr>
<tr>
<td>concentration of O₃ at t=0</td>
<td>25</td>
<td>ppb-v</td>
</tr>
<tr>
<td>concentration of H₂O₂ at t=0</td>
<td>0.4</td>
<td>ppb-v</td>
</tr>
<tr>
<td>concentration of CO₂ at t=0</td>
<td>360</td>
<td>ppm-v</td>
</tr>
<tr>
<td>concentration of HNO₃ at t=0</td>
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<td>ppb-v</td>
</tr>
<tr>
<td>concentration of NH₃ at t=0</td>
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<td>ppb-v</td>
</tr>
</tbody>
</table>

Table 3. Initial conditions for sensitivity test cases of 2-dimensional kinematic model. Specified are: aqueous phase chemistry choice, initial volume fraction of NH₃, mean radius of the assumed lognormal aerosol particle size distribution $r_d$, total aerosol concentration $n_{tot}$ and geometric standard deviation $\sigma_g$. Other parameters for each case are the same as in base case (Tab. 2). The parameters that distinguish each sensitivity test case are marked in bold.

<table>
<thead>
<tr>
<th>case</th>
<th>oxidation reaction</th>
<th>NH₃ [ppb-v]</th>
<th>$r_d$ [µm]</th>
<th>$n_{tot}$ [cm⁻³]</th>
<th>$\sigma_g$</th>
</tr>
</thead>
<tbody>
<tr>
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<td>0.05</td>
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<td>1.8</td>
</tr>
<tr>
<td>case2</td>
<td>on</td>
<td>0.4</td>
<td>0.05</td>
<td>50</td>
<td>1.8</td>
</tr>
<tr>
<td>case3</td>
<td>on</td>
<td>0.1</td>
<td>0.05</td>
<td>150</td>
<td>1.8</td>
</tr>
</tbody>
</table>

Production of sulfate suppresses precipitation. The overall impact depends strongly on the initial concentration of aerosol particles, see Feingold and Kreidenweis (2002) for the discussion. The short simulation time used in this study hinders analysis of the impact of oxidation on the overall precipitation. The work presented here focuses on the evolution of aerosol particle sizes and pH values of cloud and drizzle droplets. Future LES simulations should focus on the impacts of aqueous chemistry on precipitation, cloud lifetime, and cloud dynamics.
Figure 4. Base case setup (see Tab. 2). All panels depict model state after 30 minutes simulation time (excluding the spin-up) and show: unactivated aerosol concentration (a), cloud droplet specific concentration (b), rain water mixing ratio (c), mean dry radius (d), cloud droplet effective radius (e) and concentration of S\textsuperscript{VI} molecules (f). The thresholds for particle radii are: unactivated aerosol < 1 \(\mu\)m; 1 \(\mu\)m < cloud < 25 \(\mu\)m; rain > 25 \(\mu\)m. Note the logarithmic scale for the rain water mixing ratio plot.

The kinematic setup precludes any links between cloud microphysical processes and dynamics of the air motion. The setup limits the study to the smooth velocity and therefore smooth saturation fields and prevents mixing between air parcels with different trajectories and properties. On the other hand, the kinematic setup has low computational cost and allows easy testing and sensitivity analysis. Prescribing the velocity ensures that all changes to the aerosol particle and water drop size distributions are caused by the cloud microphysics and aqueous-phase chemistry alone. Moreover, the kinematic setup allows for a straightforward selection of the updraft and downdraft regions, further simplifying the analysis of the microphysical processes.
5.2 Results

Figure 4 shows the model state after 30 minutes of simulation from base case (see Tab. 2). Figure 4a shows the concentration of unactivated aerosol particles (defined as the SDs with wet radius smaller than 1 \( \mu m \)). The lower part of the plot (below 900 m) shows cloud-free conditions and corresponds to the initial concentration of aerosol particles. The upper part of the plot shows the interstitial aerosol particles, i.e. those in-cloud aerosol particles that did not activate. The difference between the upper and lower parts of Fig. 4a shows the impact of nucleation scavenging on aerosol population. The regions with slightly higher concentration of the in-cloud aerosol particles near the cloud base correspond to regions with low vertical velocities, lower supersaturations and thus lower concentrations of cloud droplets. Figure 4b shows the concentration of cloud droplets (defined as the SDs with wet radii between 1 and 25 \( \mu m \)). It is nearly constant with height, that agrees with the observations in stratocumulus clouds (e.g. Pawlowska et al., 2000). The regions with lower cloud droplet concentrations correspond to the regions with drizzle (see Fig. 4c). Figure 4c shows the rain water mixing ratio (water drops with wet radius greater than 25 \( \mu m \)) using a logarithmic color scale. Rain forms quickly in the simulation due to the relatively high values of cloud droplet radii after the spin-up caused by the low initial aerosol particle concentration. The footprint of precipitation can be seen in Fig. 4a and f where the cloud droplet concentration is depleted in regions of drizzle. Figure 4d shows the mean dry radius of all particles (both aerosol particles and water drops). The mean dry radius is increasing due to oxidation. In the updraft (left-hand side of panel d) new aerosol particles

...are advected into the cloudy region. Once the cloud droplets are formed, the aqueous phase oxidation starts to produce sulfate and changes the CCN size distribution. In the downdraft (right-hand side of panel d) cloud droplets are advected out of the cloud and they evaporate. The cloud-processed CCN are returned to the environment and change the ambient air aerosol particle size distribution. Figure 4e depicts the cloud droplet effective radius. As expected, the effective radius increases with height. At the top of the cloud the effective radius reaches 20 \( \mu m \), which is linked to the small cloud droplet concentration. High effective radii imply efficient drizzle production after the spin-up (usually water drop radius \( \sim \) 12 \( \mu m \) is reported as the threshold value for efficient collisions between water drops and the production of precipitation, for example Rosenfeld and Gutman (1994); Pawlowska and Brenguier (2003)). Figure 4f shows the concentration of S\textsuperscript{VI} molecules (all molecules containing sulfur at +6 oxidation state) and represents molecules from the initial ammonium bisulfate (\( \text{NH}_4\text{HSO}_4 \)) aerosol and the molecules created during oxidation. It corresponds to the mean dry radius plotted in Fig. 4d. Additionally, some effects of collisions and precipitation can be seen when comparing the irregular features from Fig. 4f with rain water mixing ratio in Fig. 4c. Precipitation displaces the largest water drops, which causes the irregular distribution of S\textsuperscript{VI} molecules in cloudy grid-cells. Figure 4f also shows that the Lagrangian particle-based scheme can track the dissolved chemical compounds in the evaporating rain drops below the cloud base.

Figure 4b, d, e and f show a layer of very clean air above the cloud which is caused by sedimentation of cloud droplets. In the downdraft region, the prescribed velocity field advects the clean layer into the domain. This feature is not present in the aerosol concentration plot (Fig. 4a) because the clean layer contains small aerosol particles with small sedimentation velocity. The depicted clean layer is an artifact caused by the prescribed velocity field and the absence of aerosol sources in the com-
putational domain. The relatively short simulation time is chosen to minimize the impact of the clean layer on the simulation.

**Figure 5.** Liquid water volume weighted average pH from base case (a), case1 (b), case2 (c), case3 (d). See Tab. 2 and Tab. 3 for a definition of simulation setups.

Figure 5 shows the liquid water volume weighted average pH in each computational grid-cell from base case (a) and sensitivity test cases (b-d). In order to better adjust the color scale to the in-cloud pH variability, pH values below 3 that correspond to very acidic aerosol particles below the cloud base have been clipped. Figure 5 captures the pH of cloud droplets as well as the pH of some evaporating rain drops below the cloud base. The droplets in the downdraft of Fig. 5a have higher acidity that is caused by $S^{VI}$ created during aqueous phase oxidation. For base case, Fig. 5a, pH increases with height above the cloud base. Initially water drops are very acidic, but they grow in size and become more diluted. Even though $S^{VI}$ is created during oxidation, the average pH still increases with height due to dilution. The same behavior is shown in the adiabatic parcel tests discussed in Sec. 4 and shown in Fig. 1. The increase of pH with height is also observed in 1-dimensional model representing processing of sulfur in small cumuli in marine environment (Alfonso and Raga, 2002). Due to the pH variability shown in Fig. 5a oxidation by $O_3$ happens mostly near the cloud top in base case. As discussed in Sec. 3.3, the rate of oxidation by $O_3$ increases significantly with increasing pH (see Eq. 4) whereas oxidation by $H_2O_2$ does not depend very weakly on
the acidity (Eq. 5). The study by Walcek and Taylor (1986) also reported that the pH of droplets increased with height due to dilution despite the production of sulfuric acid. In turn, increased pH promotes oxidation by O$_3$ in the upper parts of the cloud, whereas oxidation by H$_2$O$_2$ dominates in lower parts of the cloud, according to their study.

Case 1 shown in Fig. 5b represents a hypothetical “no oxidation” scenario where all physical and chemical conditions are the same as in base case, the reversible chemical processes are allowed, and oxidation is prohibited. The scenario without oxidation is overall less acidic than base case (Fig. 5a). Additionally, without oxidation there is no difference between the pH values in the updraft and downdraft in Fig. 5b. Without oxidation, all the chemical processes are reversible and the dissolved chemical compounds are outgassed to the atmosphere as the cloud droplets evaporate in the downdraft.

Case 2 differs from base case by increasing the initial NH$_3$ volume fraction from 0.1 ppb-v to 0.4 ppb-v (see Tab. 2 and Tab. 3). Because the initial aerosol particle size distribution is the same as in base case, the mean aerosol and droplet sizes and concentrations at the end of the simulation are not different from base case (not shown). Figure 5c shows the liquid water volume weighted average pH for case 2. The average pH in case 2 (Fig. 5c) is higher than in base case (Fig. 5a), that is, both cloud droplets and rain drops are less acidic in case 2 than in base case. In contrast to base case, in the updraft (left-hand side of the plots), the pH in case 2 actually decreases with height above the cloud base. This is because the higher initial NH$_3$ volume fraction increases its uptake and counters the low pH values caused by initial acidic aerosol particles (see Eq. (2)). Then, as the water drops are advected upwards, oxidation produces sulfuric acid and the average pH decreases. Near the cloud top, the NH$_3$ is degassed back to the environment. Case 2 results are in agreement with the trajectory ensemble model simulations by Zhang et al. (1999). In their study, the initial aerosol size distribution is the same as in base case and case 2. However, their initial trace gas volume fractions are much higher and aim to represent a “moderately polluted marine environment” (their base case NH$_3$ volume fraction is ten times larger than base case value assumed here). As in case 2 presented here, the high initial NH$_3$ volume fractions in Zhang et al. (1999) increase the pH near the cloud base and promote oxidation by O$_3$ during the first minutes after the simulated parcels entered the cloud. Because the sulfuric acid was produced, the pH dropped and oxidation by H$_2$O$_2$ becomes dominant in the higher regions of the cloud, as reported in their study.

Case 3 increases the initial aerosol concentration to 150 cm$^{-3}$, while keeping all other initial conditions the same as in base case (see Tab. 2 and Tab. 3). In general, higher initial aerosol particle concentration results in higher cloud droplet concentrations. This in turn creates smaller cloud droplet effective radii that virtually prohibits the onset of precipitation during the 30 minutes simulation time (not shown). Figure 5d shows the liquid water volume weighted average pH for case 3. Similar to base case (Fig. 5a), the pH increases with height due to the dilution and the downdraft droplets are more acidic due to the ongoing oxidation. However, case 3 is more acidic than base case because the overall droplet sizes are smaller and they are therefore less diluted.

At the end of base case simulation, 18% of the total available S$^{IV}$ is oxidized. As a result, 0.14 µg/m$^3$ of dry particulate matter are created during oxidation (an average value for the whole computational domain reported in relation to the dry air volume). In total, 40% of the final dry particulate matter is created due to oxidation and 60% originates from the initial aerosol mass. The oxidation is a significant source of dry particulate matter because the initial aerosol mass is very low (only 0.21
µg/m³ dry air). Oxidation by H₂O₂ is the dominant path: 92% of the S^{VI} molecules created during oxidation are oxidized by S^{VI} originating from S^{IV} oxidation by H₂O₂. More alkaline conditions of case2 enhance the efficiency of oxidation. At the end of case2, 21% of available S^{IV} is oxidized. As a result, oxidation produces 0.16 µg of dry particulate matter per m³ of dry air (average over the whole computational domain). For case2, 44% of the final dry particulate matter is created due to oxidation and 56% originates from the initial ammonium bisulfate aerosol. Similarly to base case, the significance of oxidation as a source of dry particulate matter is caused by a very low initial aerosol mass. Due to more alkaline conditions, oxidation by O₃ becomes more important than in base case. At the end of case3 simulation, 39% of the S^{VI} molecules that are created during oxidation are produced by the originates from S^{VI} oxidation by O₃ path and 61% by the H₂O₂ path. In contrast, more acidic conditions of case3 hinder the O₃ reaction path. Virtually all molecules of sulfate that are created during oxidation are oxidized by H₂O₂. As a result, the conversion of sulfur to sulfate is slightly less effective in case3. At the end of case3 simulation, 17% of available S^{IV} is oxidized. As a result 0.13 µg of dry particulate matter are created per m³ of dry air. At the end of case3 simulation, 17% of the dry particulate matter is created by oxidation and 83% originates from the initial aerosol. The initial aerosol mass is larger in case3 than in base case due to the higher initial aerosol concentration (case3 contains initially 0.61 µg/m³ of dry particulate matter). Because of this, even though the produced sulfate mass is only slightly lower than in base case, the relative importance of oxidation decreases by more than 20 percentage points in case3. Due to the simple kinematic setup chosen in this study the values reported here cannot be treated as representative for the atmospheric conditions. They are shown to allow comparison between base case and the sensitivity test cases.

Figure 6. Size distributions of dry radii for base case (a) and case3 (b). The initial dry radius size distribution is marked in black, final dry radius size distribution from grid-cells with \( r_c > 0.01 \) g/kg in green and from grid-cells with \( r_r > 0.01 \) g/kg in red. See Tab. 2 and Tab. 3 for a definition of simulation setups.
Finally, the impact of collisions and aqueous phase oxidation of sulfur on the aerosol and water drop size distributions is examined. For this purpose, the aerosol particle size distributions from base case (Fig. 6a) and case3 (Fig. 6b) are compared. The black line represents the initial aerosol size distribution and the green and red lines represent the final aerosol size distribution for in-cloud \( r_c > 0.01 \text{ g/kg} \) and precipitating \( r_r > 0.01 \text{ g/kg} \) grid-cells. The two cases are chosen because they have different initial aerosol size distributions. In both cases the cloud-processed aerosol size distributions (green and red lines) have a bimodal shape. This is a footprint of oxidation that creates the Hoppel minimum in the dry radius size distribution. The same effect is obtained in the adiabatic parcel tests discussed in Sec. 4 and shown in Fig. 3. Moreover, the efficient collisions between water drops in base case create a tail of bigger aerosol sizes in Fig. 6a. The effect is stronger for the precipitating grid-cells (red line). In case3 fewer collisions between water drops occur than in base case and therefore no precipitation and no tail of big aerosol particles is created. Also, in case3, the change in size distribution of aerosol particles caused by oxidation is smaller because the produced sulfate is divided among bigger number of aerosol particles.

6 Summary and outlook

The work presented here describes an extension of the *libcloudph++* library that allows including aqueous phase chemical reactions within water drops in the Lagrangian particle-based microphysics scheme. The extension covers the aqueous phase oxidation of sulfur to sulfate. The modular way in which the library is implemented along with the provided documentation should allow, if needed, further development to cover more chemical compounds and reactions. The 0-dimensional and 2-dimensional tests described in this work as well as comparison with other numerical studies using bin microphysics schemes along with aqueous chemistry representation document the correctness of the design and the implementation of the LMC-validate the the particle-based microphysics and chemistry scheme. Additionally, the changes in the user-programming interface due to aqueous chemistry extension are described in Sec. A in the Appendix. Section C in the Appendix completes the description with a list of chemical constants used in the library and chemical reactions included.

The models used in this study to test the chemistry scheme provide a simplified view of the macrophysical cloud properties. They enable validation and testing of the Lagrangian particle-based scheme but do not provide a good balance between the representation of cloud microphysics and dynamics. As a next step, the Lagrangian particle-based scheme needs to be coupled to an eddy-resolving model. This would allow quantifying how microphysical and chemical processes affect precipitation in the model and how they affect the cloud lifetimes simulated by the model.

*Code availability.* The *libcloudph++* library along with the aqueous-phase chemistry extension, the parcel model and the 2D kinematic model are released under GNU General Public License v3.0. The version of *libcloudph++* accompanying this publication is tagged as “2.0.0” at the project repository and is also available as an electronic supplement to this paper. *libcloudph++* and the 2D slice model are available at: https://github.com/igfuw/libcloudphxx and the parcel model is available at: https://github.com/igfuw/parcel.

The supported platforms are: Linux with GNU g++, Linux with LLVM clang++ and Apple OSX with the Apple clang++. The code requires c++14 support. The compilation is tested using Travis continuous integration framework.
The user interface of the Lagrangian microphysics scheme in liycloudph++ is used. The user interface of the Lagrangian microphysics scheme is presented in Appendix A. Here, additional information related to the new aqueous phase chemistry scheme is provided. The liycloudph++ is implemented in C++ and therefore some nomenclature related to this programming language is used. For a thorough introduction to C++ programming language see Stroustrup (2013).

The aqueous chemistry module is implemented as an optional extension to the Lagrangian particle-based microphysics scheme in liycloudph++. It uses the same liycloudphxx::lgrngn namespace as the original scheme. Again the template parameter real_t selects between floating point formats of simulations. The Lagrangian particle-based microphysics scheme options are grouped into a structure named lgrngn::opts_t. Chemistry module adds three Boolean fields to this structure: chem_dsl chem_dsc and chem_rct, see code listing in Fig. A1. When set to true by the user, they switch on dissolving of trace gases into water drops, dissociation of chemical compounds in water drops and oxidation reaction, respectively. The parameters in lgrngn::opts_t can be changed during simulation. For example during the 2-dimensional kinematic simulations from Sec. 5.3, oxidation is enabled by setting the chem_rct parameter to true at the end of spin-up. Other parameters that cannot be changed during simulation are encapsulated in lgrngn:: opts_init_t structure. Chemistry module adds three fields to this structure: (i) A Boolean chem_switch field that enables memory allocation for additional variables needed for chemistry representation. (ii) An integer sstp_chem field that defines the number of sub-steps to be carried out in aqueous chemistry calculations. (iii) A real_t chem_rho field that defines the dry aerosol density, see code listing in Fig. A2.

The names of chemical compounds available in the aqueous phase chemistry module are stored in a chem_species_t enumerator enum, see code listing in Fig. A3. The state of all variables used by the Lagrangian particle-based scheme is stored in an instance of the lgrngn::particles_t structure shown in code listing in Fig. A4. The second template parameter of that structure selects between CPU and GPU calculations (see the discussion in Sec. 5.2 in Arabas et al., 2015, for details). The initialization, time-stepping and output from the Lagrangian particle-based scheme are done using the methods of lgrngn::particles_t structure. Their signatures are provided in code listing in Fig. A4.

```
template<typename real_t>
struct opts_t
{
    // process toggling for chemistry
    bool chem_dsl, chem_dsc, chem_rct;
    // ...

Figure A1. lgrngn::opts_t definition

Appendix A: User-Programming Interface

The user-programming interface of the Lagrangian particle-based microphysics scheme of liycloudph++ is presented in Sec 5.2. in Arabas et al. (2015). Here, additional information related to the new aqueous phase chemistry scheme is provided. The liycloudph++ is implemented in C++ and therefore some nomenclature related to this programming language is used. For a thorough introduction to C++ programming language see Stroustrup (2013).

The aqueous chemistry module is implemented as an optional extension to the Lagrangian particle-based microphysics scheme in liycloudph++. It uses the same liycloudphxx::lgrngn namespace as the original scheme. Again the template parameter real_t selects between floating point formats of simulations. The Lagrangian particle-based microphysics scheme options are grouped into a structure named lgrngn::opts_t. Chemistry module adds three Boolean fields to this structure: chem_dsl chem_dsc and chem_rct, see code listing in Fig. A1. When set to true by the user, they switch on dissolving of trace gases into water drops, dissociation of chemical compounds in water drops and oxidation reaction, respectively. The parameters in lgrngn::opts_t can be changed during simulation. For example during the 2-dimensional kinematic simulations from Sec. 5.3, oxidation is enabled by setting the chem_rct parameter to true at the end of spin-up. Other parameters that cannot be changed during simulation are encapsulated in lgrngn:: opts_init_t structure. Chemistry module adds three fields to this structure: (i) A Boolean chem_switch field that enables memory allocation for additional variables needed for chemistry representation. (ii) An integer sstp_chem field that defines the number of sub-steps to be carried out in aqueous chemistry calculations. (iii) A real_t chem_rho field that defines the dry aerosol density, see code listing in Fig. A2.

The names of chemical compounds available in the aqueous phase chemistry module are stored in a chem_species_t enumerator enum, see code listing in Fig. A3. The state of all variables used by the Lagrangian particle-based scheme is stored in an instance of the lgrngn::particles_t structure shown in code listing in Fig. A4. The second template parameter of that structure selects between CPU and GPU calculations (see the discussion in Sec. 5.2 in Arabas et al., 2015, for details). The initialization, time-stepping and output from the Lagrangian particle-based scheme are done using the methods of lgrngn::particles_t structure. Their signatures are provided in code listing in Fig. A4.
or sedimentation). The list of arguments of
pointing to the corresponding trace gas mixing ratio fields from the driver model. This is an optional argument for simulations
describe the velocity field. The last argument of

During time-stepping the Lagrangian scheme overwrites the driver model fields during simulation. The

\begin{verbatim}
template<typename real_t>
struct opts_init_t
{
  // if false no chemical reactions are allowed
  // (no memory allocation)
  bool chem_switch;
  // substeps for chemistry
  int sstp_chem;
  // assumed dry aerosol density
  real_t chem_rho;
}
\end{verbatim}

\textbf{Figure A2. lgrngn::opts_init_t definition}

\begin{verbatim}
enum chem_species_t
{
  // both gas and total dissolved chem species
  HNO3, NH3, CO2, SO2, H2O2, O3,
  // created sulfate
  S_VI,
  // additional H+ for pH
  H
}
\end{verbatim}

\textbf{Figure A3. lgrngn::chem_species_t definition}

The \texttt{init()} method performs initialization and should be called first. As discussed in Arabas et al. (2015) the first three arguments are obligatory and should point to the dry air potential temperature, water vapor mixing ratio and dry air density fields of the driver model that uses the \texttt{libcloudph++}. The next three arguments should point to the Courant number field components. They are optional and depend on the dimensionality of the solved problem. For example, for the parcel model tests from Sec. 4 none are necessary, whereas for the 2-dimensional kinematic model from Sec. 5 two arguments are specified in order to describe the velocity field. The last argument of \texttt{init()} is a map with keys from \texttt{chem_species_t} and values pointing to the corresponding trace gas mixing ratio fields from the driver model. This is an optional argument for simulations with aqueous phase chemistry.

During time-stepping the \texttt{Lagrangian-particle-based} scheme computations are performed by \texttt{step_sync()} and \texttt{step_async()} methods. The first one gathers all the processes that affect the driver model fields (such as condensation/evaporation or aqueous phase chemistry) and the second one gathers all the processes that can be calculated asynchronously (for example collisions or sedimentation). The list of arguments of \texttt{step_sync()} method is extended by the chemistry module. Similar to the \texttt{init()} method, a map linking \texttt{chem_species_t} items with the driver model mixing ratio fields needs to be provided as the last optional argument. The \texttt{Lagrangian-particle-based} scheme overwrites the driver model fields during simulation. The
template <typename real_t, backend_t backend_t>
struct particles_t: particles_proto_t<real_t>
{
   // initialisation
   void init(
      const arrinfo_t<real_t> th,
      const arrinfo_t<real_t> rv,
      const arrinfo_t<real_t> rhod,
      const arrinfo_t<real_t> courant_x,
      const arrinfo_t<real_t> courant_y,
      const arrinfo_t<real_t> courant_z,
      const std::map<
         enum chem_species_t,
         arrinfo_t<real_t> > ambient_chem
   );

   // time-stepping methods
   void step_sync(
      const opts_t<real_t> &,
      arrinfo_t<real_t> th,
      arrinfo_t<real_t> rv,
      const arrinfo_t<real_t> rhod,
      const arrinfo_t<real_t> courant_x,
      const arrinfo_t<real_t> courant_y,
      const arrinfo_t<real_t> courant_z,
      std::map<
         enum chem_species_t,
         arrinfo_t<real_t> > ambient_chem,
   );

   void step_async(
      const opts_t<real_t> &
   );

   // diagnostic methods
   // ...
   void diag_chem(const enum chem_species_t&);
   // ...

Figure A4. lgrngn::particles_t definition
signature of `step_async` method is not changed by the new chemistry module.

As discussed in Arabas et al. (2015), the `lgrngn::particles_t` structure provides many methods for obtaining statistical information on the SD parameters (prefixed with `diag`). The chemistry model adds to them the `diag_chem` method that outputs the total mass of a chemical compound dissolved into droplets. The chemical compound is selected using the `chem_species_t` enumerator items. See the discussion in Sec. 5.2 in Arabas et al. (2015) for the details on how to select the size ranges of droplets specified for output or how to output other statistical parameters.
Appendix B: Glossary

\( \alpha_M \) mass accommodation coefficient of water vapor
\( \alpha_{MA} \) mass accommodation coefficient of the chemical compound “A”
\( c_{A \infty} \) ambient concentration of the trace gas “A”
\( c_{pd} \) specific heat at constant pressure for dry air
\( D_A \) diffusion coefficient of the chemical compound “A”
\( E \) reaction activation energy
\( H_A^{\text{eff}} \) effective Henry’s law constant of the chemical compound “A”
\( \Delta H_D \) reaction enthalpy of dissociation at constant temperature and pressure
\( \Delta H_H \) reaction enthalpy of dissolution at constant temperature and pressure
\( K_A \) dissociation constant
\( k_{0,...,4} \) reaction rate coefficients
\( \kappa \) hygrosopicity parameter
\( M_A \) molar mass of the chemical compound “A”
\( n(r_d) \) spectral density function of aerosol particle sizes
\( n_{\text{tot}} \) total aerosol concentration
\( N \) super-droplet multiplicity
\( \theta \) dry air potential temperature
\( p_v \) partial pressure of water vapor
\( p_{1000} \) pressure equal 1000 hPa
\( R_d \) gas constant for dry air
\( R_A \) reaction rate of the chemical compound “A”
\( \rho_d \) dry air density
\( r_d \) dry radius
\( \bar{r}_d \) mean radius of the assumed lognormal aerosol particle size distribution
\( r_w \) drop radius
\( r_c \) cloud water mixing ratio
\( r_r \) rain water mixing ratio
\( r_v \) water vapor mixing ratio
\( \sigma_g \) geometric standard deviation
\( <v> \) average velocity of the molecules

Appendix C: List of chemical compounds and constants
Table C1. Chemical compounds considered in this work.

<table>
<thead>
<tr>
<th>chemical compound</th>
<th>formula</th>
<th>molar mass (g moles(^{-1}))</th>
<th>source</th>
</tr>
</thead>
<tbody>
<tr>
<td>ammonia</td>
<td>NH(_3)</td>
<td>17</td>
<td>trace gas</td>
</tr>
<tr>
<td>carbon dioxide</td>
<td>CO(_2)</td>
<td>44</td>
<td>trace gas</td>
</tr>
<tr>
<td>hydrogen peroxide</td>
<td>H(_2)O(_2)</td>
<td>34</td>
<td>trace gas</td>
</tr>
<tr>
<td>nitric acid</td>
<td>HNO(_3)</td>
<td>63</td>
<td>trace gas</td>
</tr>
<tr>
<td>ozone</td>
<td>O(_3)</td>
<td>48</td>
<td>trace gas</td>
</tr>
<tr>
<td>sulfur dioxide</td>
<td>SO(_2)</td>
<td>64</td>
<td>trace gas</td>
</tr>
<tr>
<td>sulfuric acid</td>
<td>H(_2)SO(_4)</td>
<td>98</td>
<td>oxidation reaction product</td>
</tr>
<tr>
<td>ammonium bisulfate</td>
<td>NH(_4)HSO(_4)</td>
<td>115</td>
<td>initial aerosol</td>
</tr>
</tbody>
</table>

Table C2. Dissociation constants and their temperature correction dependence coefficients (taken from Kreidenweis et al., 2003). Dissociation of H\(_2\)SO\(_4\) is taken from Tab. 6.7A.1 in Seinfeld and Pandis (2016).

<table>
<thead>
<tr>
<th>equilibrium reaction</th>
<th>dissociation constant at 298K (moles liter(^{-1}))</th>
<th>temp. correction (\frac{\Delta H_d}{R}) (K)</th>
</tr>
</thead>
<tbody>
<tr>
<td>(K_{\text{HNO}})</td>
<td>HNO(_3)(aq) (\rightarrow) H(^+) + NO(_3^-)</td>
<td>15.4</td>
</tr>
<tr>
<td>(K_{\text{SO}})</td>
<td>SO(_2) + H(_2)O (\rightarrow) H(^+) + HSO(_3^-)</td>
<td>(1.3 \times 10^{-2})</td>
</tr>
<tr>
<td>(K_{\text{NH}})</td>
<td>NH(_3) + H(_2)O (\rightarrow) NH(_4^+) + OH(^-)</td>
<td>(1.7 \times 10^{-5})</td>
</tr>
<tr>
<td>(K_{\text{CO}})</td>
<td>CO(_2) + H(_2)O (\rightarrow) H(^+) + HCO(_3^-)</td>
<td>(4.3 \times 10^{-7})</td>
</tr>
<tr>
<td>(K_{\text{HSO}})</td>
<td>HSO(_3^-) (\rightarrow) H(^+) + SO(_3^{2-})</td>
<td>(6.6 \times 10^{-8})</td>
</tr>
<tr>
<td>(K_{\text{HCO}})</td>
<td>HCO(_3^-) (\rightarrow) H(^+) + CO(_3^{2-})</td>
<td>(4.68 \times 10^{-11})</td>
</tr>
<tr>
<td></td>
<td>H(_2)SO(_4)(aq) (\rightarrow) H(^+) + HSO(_4^{2-})</td>
<td>(\infty)</td>
</tr>
<tr>
<td>(K_{\text{H2SO}})</td>
<td>HSO(_4^{2-}) (\rightarrow) H(^+) + SO(_4^{2-})</td>
<td>(1.2 \times 10^{-2})</td>
</tr>
</tbody>
</table>
Table C3. Dissolution constants and their temperature correction coefficients (taken from Kreidenweis et al., 2003).

<table>
<thead>
<tr>
<th>equilibrium reaction</th>
<th>dissolution Henry's law constant at 298K (moles liter$^{-1}$ atm$^{-1}$)</th>
<th>temp. corredep. $\frac{-\Delta H_{R}}{R}$ (K)</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\text{H}_3\text{NO}_3$</td>
<td>$\text{HNO}_3(g) \rightarrow \text{HNO}_3(aq)$</td>
<td>$2.10 \times 10^5$</td>
</tr>
<tr>
<td>$\text{H}_2\text{O}_2$</td>
<td>$\text{H}_2\text{O}_2(g) \rightarrow \text{H}_2\text{O}_2(aq)$</td>
<td>$7.45 \times 10^4$</td>
</tr>
<tr>
<td>$\text{NH}_3$</td>
<td>$\text{NH}_3(g) \rightarrow \text{NH}_3 + \text{H}_2\text{O}$</td>
<td>62</td>
</tr>
<tr>
<td>$\text{SO}_2$</td>
<td>$\text{SO}_2(g) \rightarrow \text{SO}_2 + \text{H}_2\text{O}$</td>
<td>1.23</td>
</tr>
<tr>
<td>$\text{CO}_2$</td>
<td>$\text{CO}_2(g) \rightarrow \text{CO}_2 + \text{H}_2\text{O}$</td>
<td>$3.40 \times 10^{-2}$</td>
</tr>
<tr>
<td>$\text{O}_3$</td>
<td>$\text{O}_3(g) \rightarrow \text{O}_3(aq)$</td>
<td>$1.13 \times 10^{-2}$</td>
</tr>
</tbody>
</table>

Table C4. Diffusion constants (Massman, 1998; Tang et al., 2014) and accommodation coefficients (Kreidenweis et al., 2003) for relevant chemical compounds.

<table>
<thead>
<tr>
<th></th>
<th>diffusion coeff. $D_A$ (m$^2$/s)</th>
<th>mass accommodation coeff. $\sigma_{MA}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\text{HNO}_3$</td>
<td>$65.25 \times 10^{-6}$</td>
<td>0.05</td>
</tr>
<tr>
<td>$\text{H}_2\text{O}_2$</td>
<td>$87.00 \times 10^{-6}$</td>
<td>0.018</td>
</tr>
<tr>
<td>$\text{NH}_3$</td>
<td>$19.78 \times 10^{-6}$</td>
<td>0.05</td>
</tr>
<tr>
<td>$\text{SO}_2$</td>
<td>$10.89 \times 10^{-6}$</td>
<td>0.035</td>
</tr>
<tr>
<td>$\text{CO}_2$</td>
<td>$13.81 \times 10^{-6}$</td>
<td>0.05</td>
</tr>
<tr>
<td>$\text{O}_3$</td>
<td>$14.44 \times 10^{-6}$</td>
<td>0.00053</td>
</tr>
</tbody>
</table>

Table C5. Reaction rate coefficients and their temperature correction coefficients (taken from Kreidenweis et al., 2003).

<table>
<thead>
<tr>
<th>oxidation reaction path</th>
<th>reaction rate coefficient (liter moles$^{-1}$ s$^{-1}$) at 298K</th>
<th>temperature corredep. $\frac{-E}{R}$ (K)</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\text{O}_3(aq) + \text{SO}_2 + \text{H}_2\text{O} \rightarrow \text{S}^{VI}$</td>
<td>$k_0 = 2.4 \times 10^4$</td>
<td>0</td>
</tr>
<tr>
<td>$\text{O}_3(aq) + \text{HSO}_3^- \rightarrow \text{S}^{VI}$</td>
<td>$k_1 = 3.5 \times 10^5$</td>
<td>-5530</td>
</tr>
<tr>
<td>$\text{O}_3(aq) + \text{SO}_3^- \rightarrow \text{S}^{VI}$</td>
<td>$k_2 = 1.5 \times 10^9$</td>
<td>-5280</td>
</tr>
<tr>
<td>$\text{H}_2\text{O}_2(aq) + \text{HSO}_3^- \rightarrow \text{S}^{VI}$</td>
<td>$k_3 = 7.45 \times 10^7$</td>
<td>-4430</td>
</tr>
</tbody>
</table>
Competing interests. The authors declare that they have no competing interests

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References


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